

# Chiral two- and three-nucleon forces with explicit Delta degree of freedom

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in collaboration with

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# Outline

- Introduction&Motivation
- 2-N forces with explicit  $\Delta$
- 3-N forces with explicit  $\Delta$
- $\pi$ N scattering with explicit  $\Delta$
- Summary and Outlook

# EFT with explicit $\Delta(1232)$

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Hemmert, Holstein, Kambor '98

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Hemmert, Holstein, Kambor '98
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→ Explicit decoupling of  $\Delta$  makes comparison with  $\Delta$ -less theory more transparent




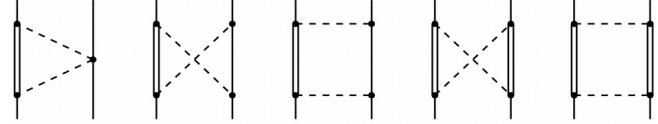

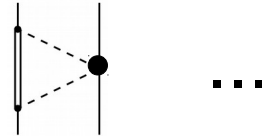

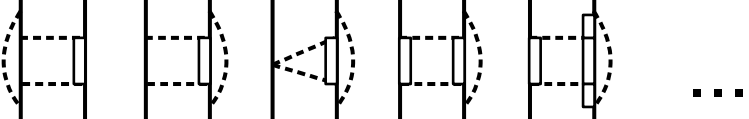


Bernard, Fearing, Hemmert, Meißner '98

finite parts of LECs can be always chosen such that

Appelquist, Carrazone '74 (Decoupling theorem)

$$\lim_{\Delta \rightarrow \infty} = \Delta\text{-less}$$

# Small scale expansion of 2NF

	$\Delta$ -less theory	$\Delta$ -full theory: additional graphs
LO		
NLO		 Keiser, Gerstendorfer, Weise '98
N <sup>2</sup> LO		 Krebs, Epelbaum, Meißner '07
N <sup>3</sup> LO		 ...
N <sup>4</sup> LO		 ...

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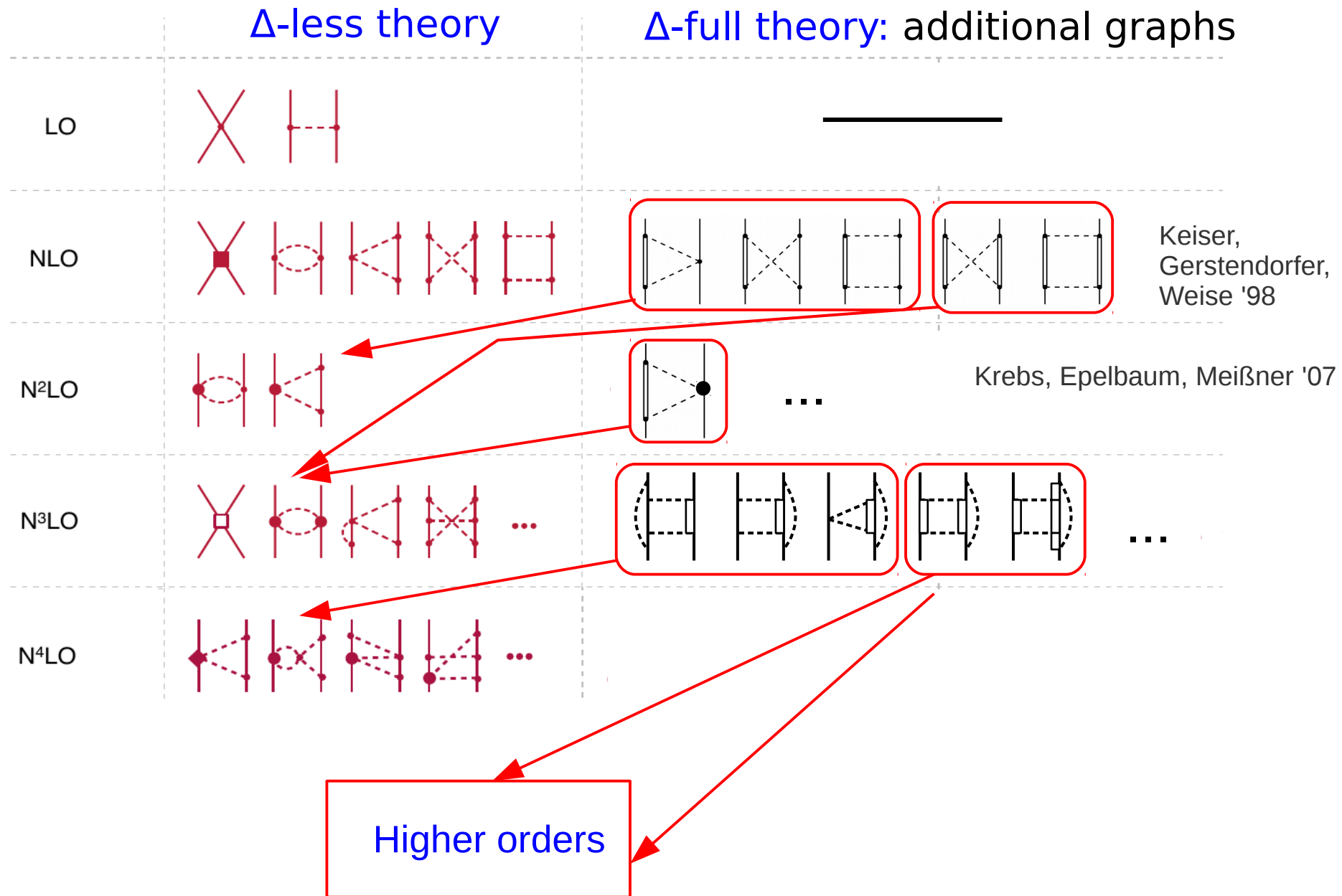
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# Small scale expansion of 2NF



# Preliminary results for N<sup>3</sup>LO 2N forces with explicit $\Delta$

- Only 2-pion-exchange contribution are considered (the long range part)
- $1/m_N$  corrections are not yet included
- Results for peripheral phases, no refitting of LEC's, no cut offs
- No additional parameters,  $h_A$  and  $g_1$  ( $\pi N\Delta$  and  $\pi\Delta\Delta$ ) are extracted from the fit to  $\pi N$  scattering

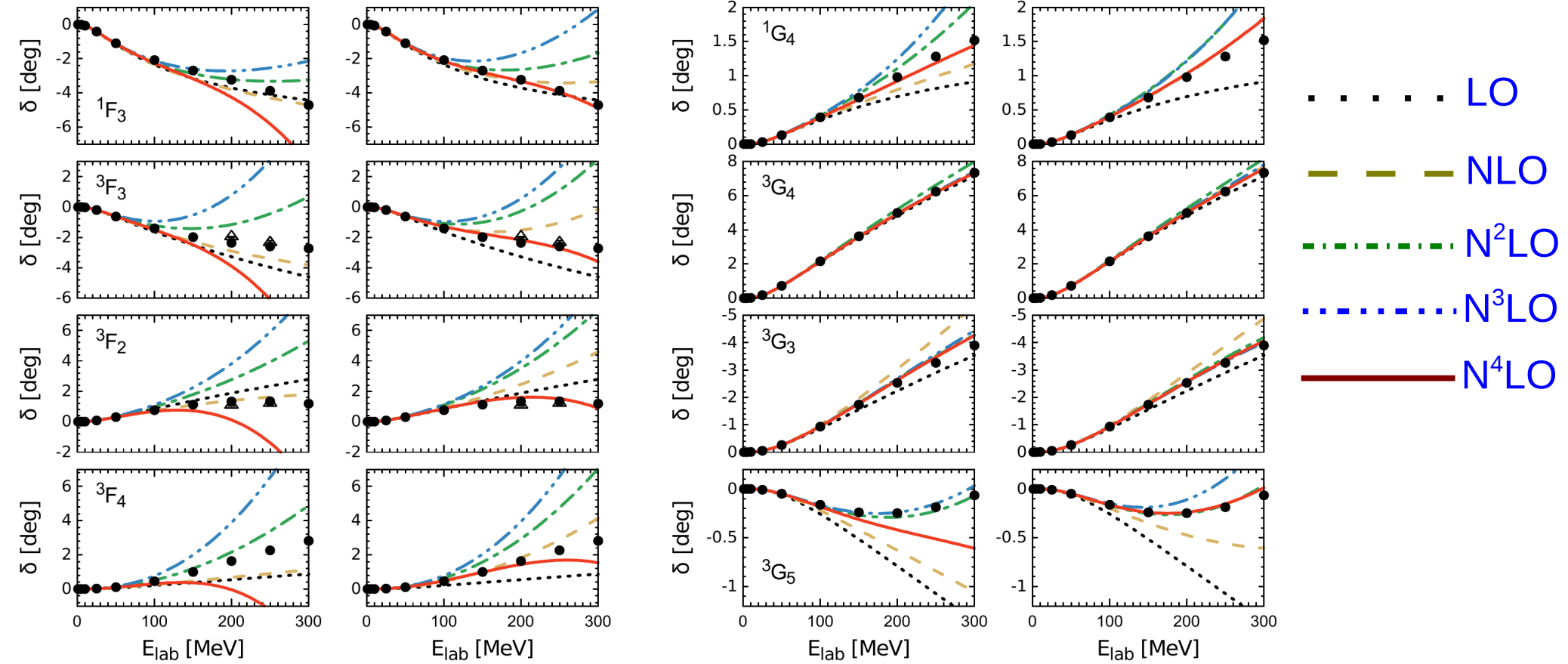
# F and G waves

$\Delta$ -less

$\Delta$ -full

$\Delta$ -less

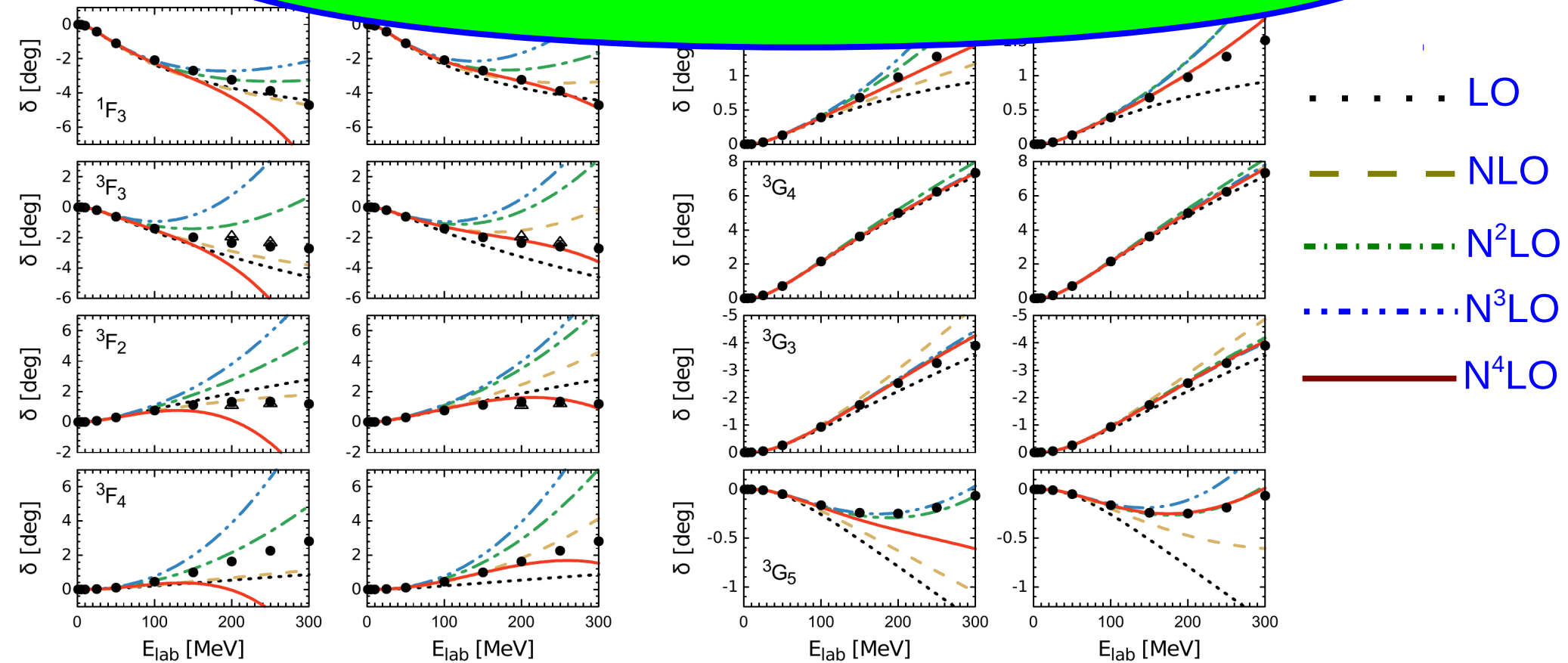
$\Delta$ -full



Data: Nijmegen PWA

# F and G waves

$\Delta$ - F-waves might be sensitive to the short-range physics

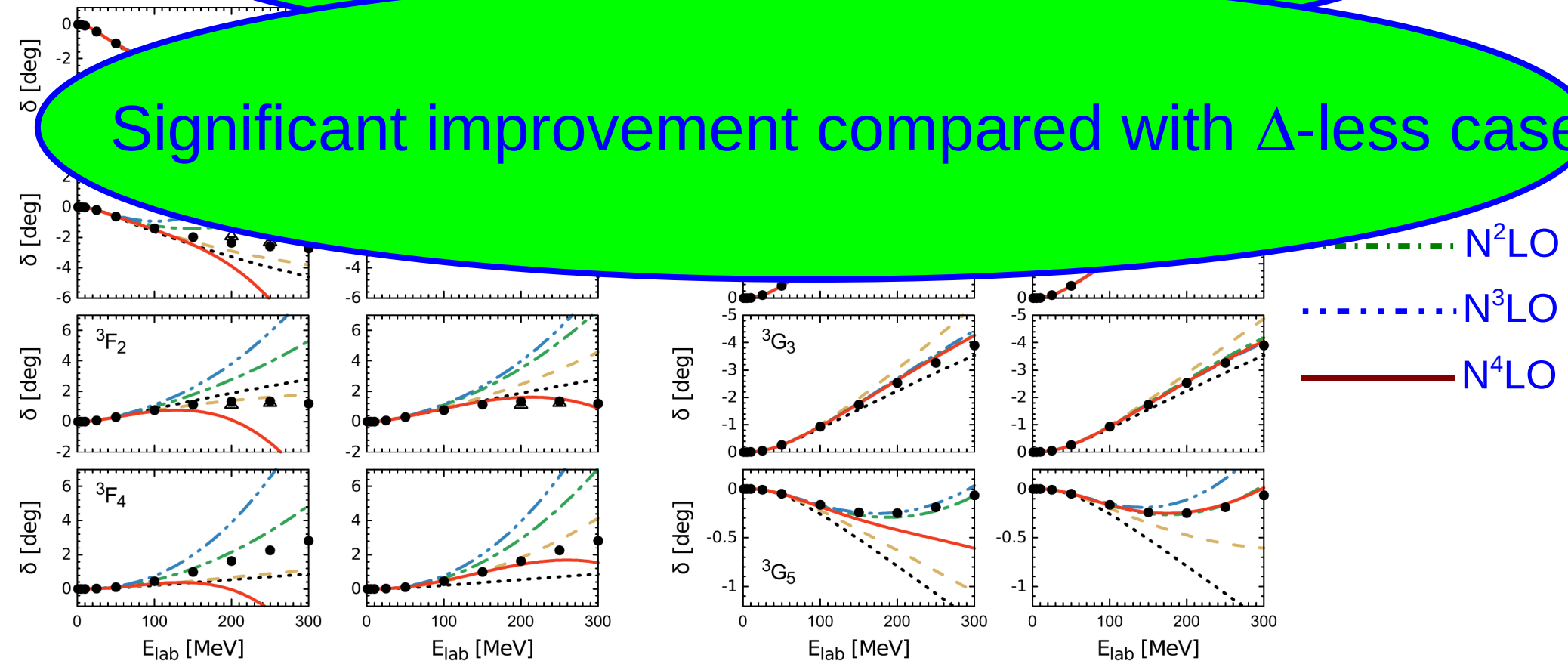


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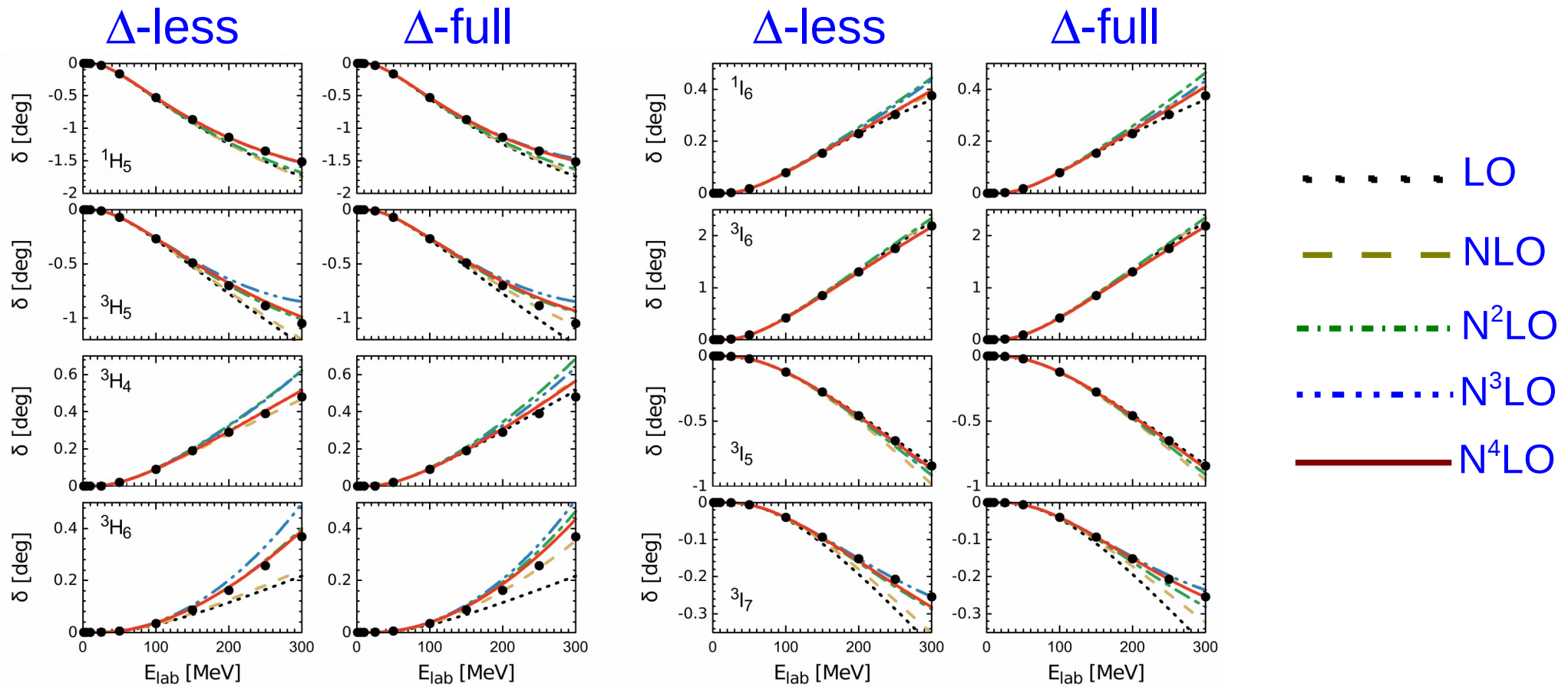
$\Delta$ - F-waves might be sensitive to the short-range physics

Significant improvement compared with  $\Delta$ -less case



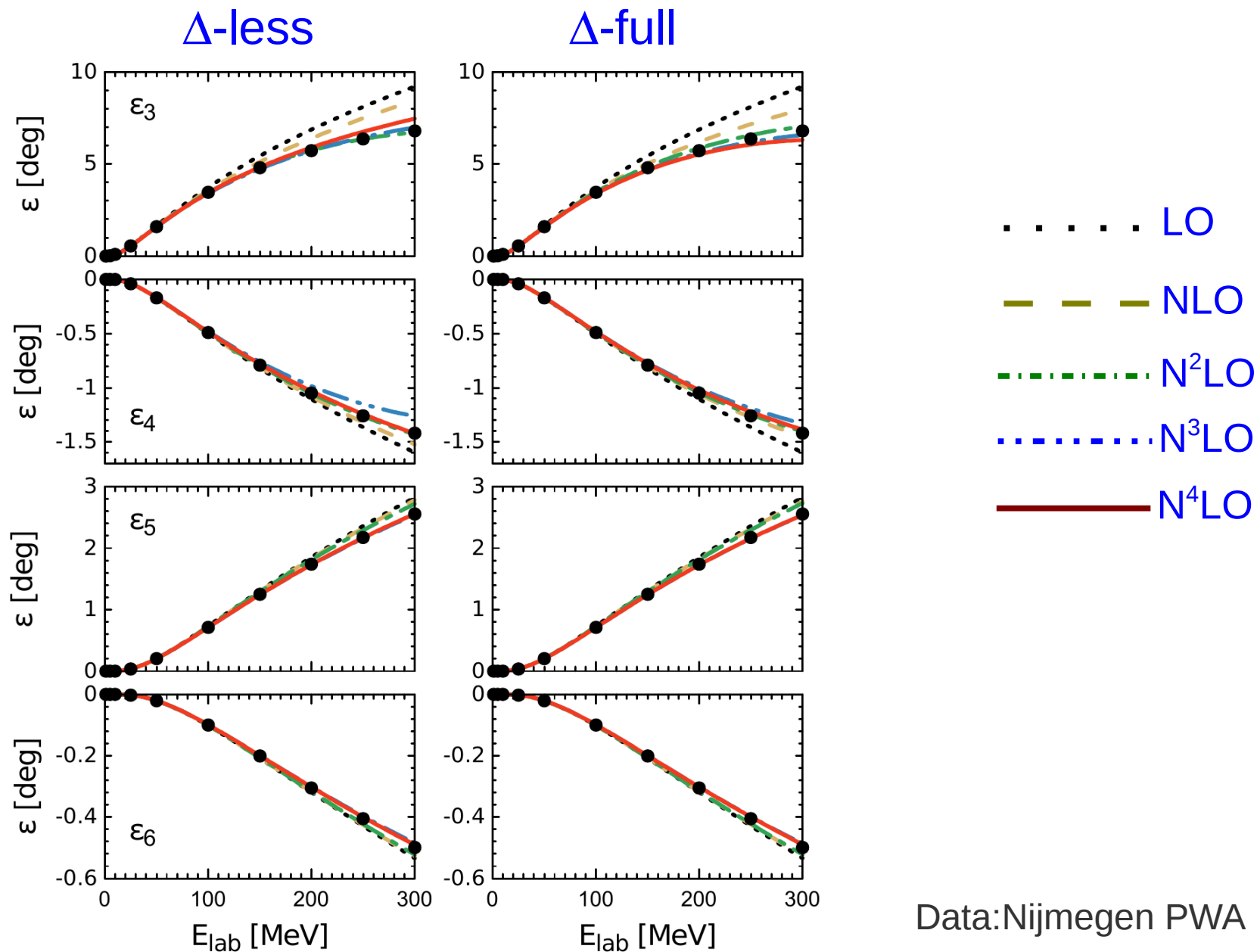
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# H and I waves




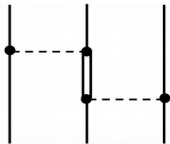
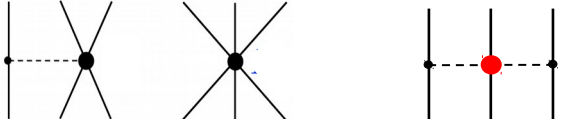

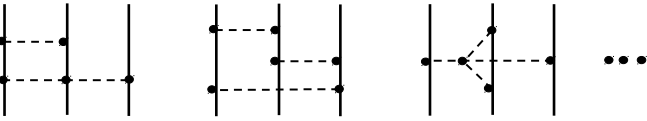
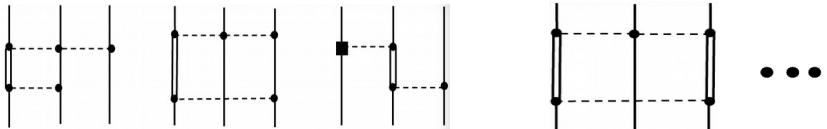
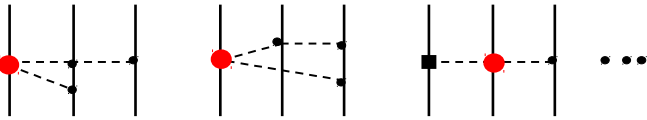

Data: Nijmegen PWA

# Mixing angles $\varepsilon_3, \varepsilon_4, \varepsilon_5, \varepsilon_6$





# Small scale expansion of 3NF

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N <sup>2</sup> LO		
N <sup>3</sup> LO		
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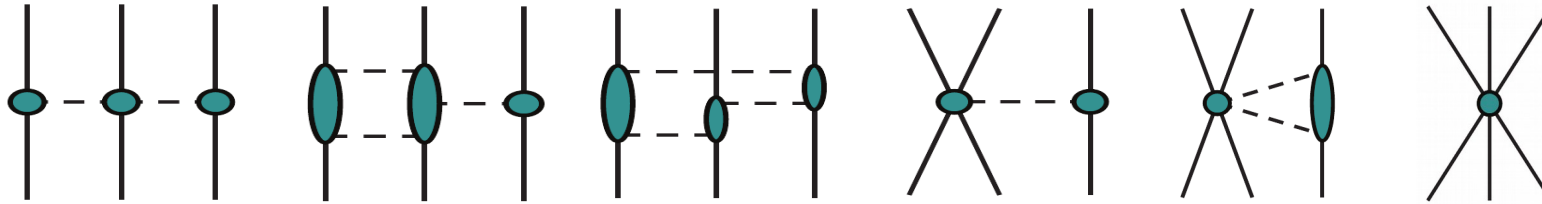
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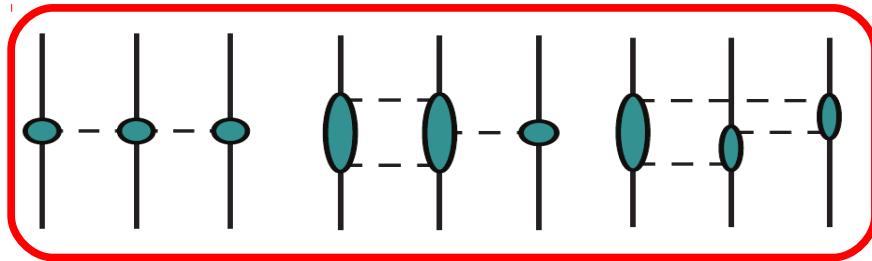
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# Long-range 3NF



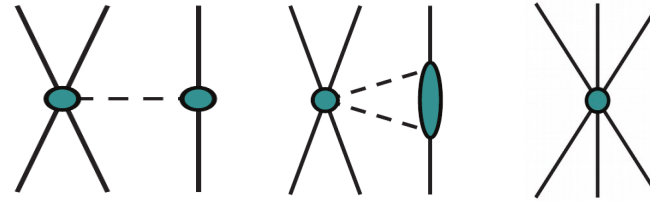
# Long-range 3NF



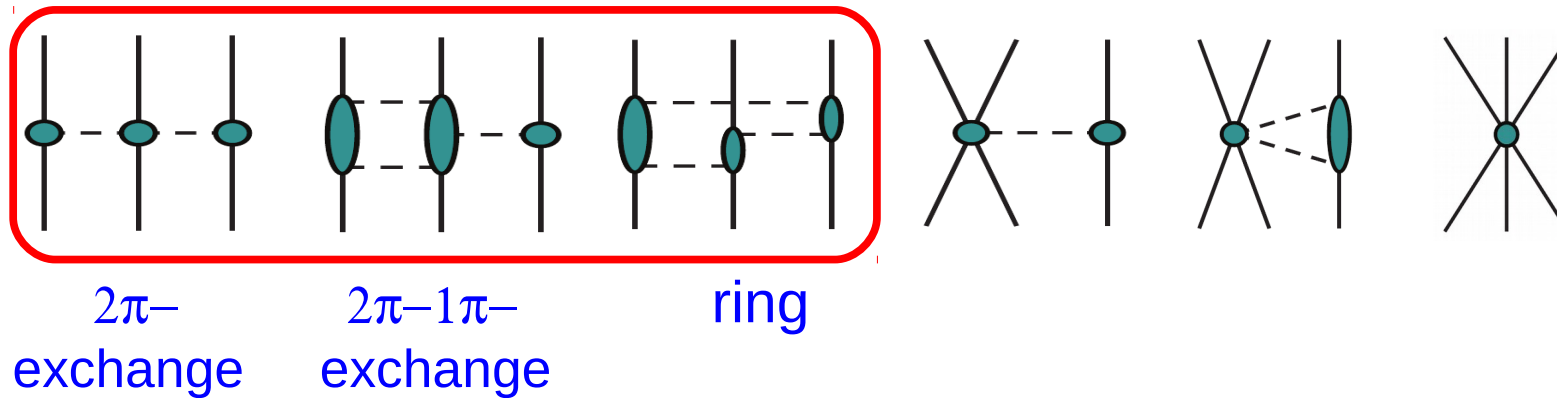
$2\pi$ -  
exchange

$2\pi-1\pi$ -  
exchange

ring



# Long-range 3NF



- Only the long range part considered (coordinate space)
- Scheme independent
- No unknown parameters

# Most general structure of a local 3NF

Krebs, Gasparyan, Epelbaum '13

Up to  $N^4$ LO all considered contribution are local

Constraints:

- Locality
- Isospin symmetry
- Parity and time-reversal invariance

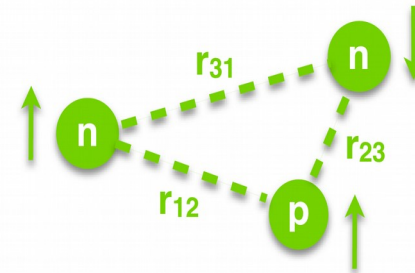
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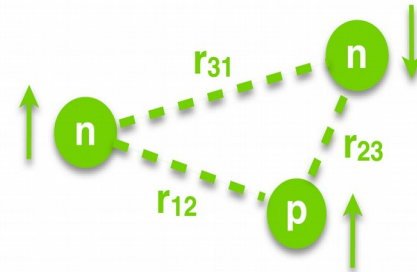
Krebs, Gasparyan, Epelbaum '13

Up to N<sup>4</sup>LO all considered contribution are local

$$\begin{aligned}
 \tilde{G}_1 &= 1, \\
 \tilde{G}_2 &= \boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_3, \\
 \tilde{G}_3 &= \vec{\sigma}_1 \cdot \vec{\sigma}_3, \\
 \tilde{G}_4 &= \boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_3 \vec{\sigma}_1 \cdot \vec{\sigma}_3, \\
 \tilde{G}_5 &= \boldsymbol{\tau}_2 \cdot \boldsymbol{\tau}_3 \vec{\sigma}_1 \cdot \vec{\sigma}_2, \\
 \tilde{G}_6 &= \boldsymbol{\tau}_1 \cdot (\boldsymbol{\tau}_2 \times \boldsymbol{\tau}_3) \vec{\sigma}_1 \cdot (\vec{\sigma}_2 \times \vec{\sigma}_3), \\
 \tilde{G}_7 &= \boldsymbol{\tau}_1 \cdot (\boldsymbol{\tau}_2 \times \boldsymbol{\tau}_3) \vec{\sigma}_2 \cdot (\hat{r}_{12} \times \hat{r}_{23}), \\
 \tilde{G}_8 &= \hat{r}_{23} \cdot \vec{\sigma}_1 \hat{r}_{23} \cdot \vec{\sigma}_3, \\
 \tilde{G}_9 &= \hat{r}_{23} \cdot \vec{\sigma}_3 \hat{r}_{12} \cdot \vec{\sigma}_1, \\
 \tilde{G}_{10} &= \hat{r}_{23} \cdot \vec{\sigma}_1 \hat{r}_{12} \cdot \vec{\sigma}_3, \\
 \tilde{G}_{11} &= \boldsymbol{\tau}_2 \cdot \boldsymbol{\tau}_3 \hat{r}_{23} \cdot \vec{\sigma}_1 \hat{r}_{23} \cdot \vec{\sigma}_2, \\
 \tilde{G}_{12} &= \boldsymbol{\tau}_2 \cdot \boldsymbol{\tau}_3 \hat{r}_{23} \cdot \vec{\sigma}_1 \hat{r}_{12} \cdot \vec{\sigma}_2, \\
 \tilde{G}_{13} &= \boldsymbol{\tau}_2 \cdot \boldsymbol{\tau}_3 \hat{r}_{12} \cdot \vec{\sigma}_1 \hat{r}_{23} \cdot \vec{\sigma}_2, \\
 \tilde{G}_{14} &= \boldsymbol{\tau}_2 \cdot \boldsymbol{\tau}_3 \hat{r}_{12} \cdot \vec{\sigma}_1 \hat{r}_{12} \cdot \vec{\sigma}_2, \\
 \tilde{G}_{15} &= \boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_3 \hat{r}_{13} \cdot \vec{\sigma}_1 \hat{r}_{13} \cdot \vec{\sigma}_3, \\
 \tilde{G}_{16} &= \boldsymbol{\tau}_2 \cdot \boldsymbol{\tau}_3 \hat{r}_{12} \cdot \vec{\sigma}_2 \hat{r}_{12} \cdot \vec{\sigma}_3, \\
 \tilde{G}_{17} &= \boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_3 \hat{r}_{23} \cdot \vec{\sigma}_1 \hat{r}_{12} \cdot \vec{\sigma}_3, \\
 \tilde{G}_{18} &= \boldsymbol{\tau}_1 \cdot (\boldsymbol{\tau}_2 \times \boldsymbol{\tau}_3) \vec{\sigma}_1 \cdot \vec{\sigma}_3 \vec{\sigma}_2 \cdot (\hat{r}_{12} \times \hat{r}_{23}), \\
 \tilde{G}_{19} &= \boldsymbol{\tau}_1 \cdot (\boldsymbol{\tau}_2 \times \boldsymbol{\tau}_3) \vec{\sigma}_3 \cdot \hat{r}_{23} \hat{r}_{23} \cdot (\vec{\sigma}_1 \times \vec{\sigma}_2), \\
 \tilde{G}_{20} &= \boldsymbol{\tau}_1 \cdot (\boldsymbol{\tau}_2 \times \boldsymbol{\tau}_3) \vec{\sigma}_1 \cdot \hat{r}_{23} \vec{\sigma}_2 \cdot \hat{r}_{23} \vec{\sigma}_3 \cdot (\hat{r}_{12} \times \hat{r}_{23}), \\
 \tilde{G}_{21} &= \boldsymbol{\tau}_1 \cdot (\boldsymbol{\tau}_2 \times \boldsymbol{\tau}_3) \vec{\sigma}_1 \cdot \hat{r}_{13} \vec{\sigma}_3 \cdot \hat{r}_{13} \vec{\sigma}_2 \cdot (\hat{r}_{12} \times \hat{r}_{23}), \\
 \tilde{G}_{22} &= \boldsymbol{\tau}_1 \cdot (\boldsymbol{\tau}_2 \times \boldsymbol{\tau}_3) \vec{\sigma}_1 \cdot \hat{r}_{23} \vec{\sigma}_3 \cdot \hat{r}_{12} \vec{\sigma}_2 \cdot (\hat{r}_{12} \times \hat{r}_{23}),
 \end{aligned}$$

Constraints:

- Locality
- Isospin symmetry
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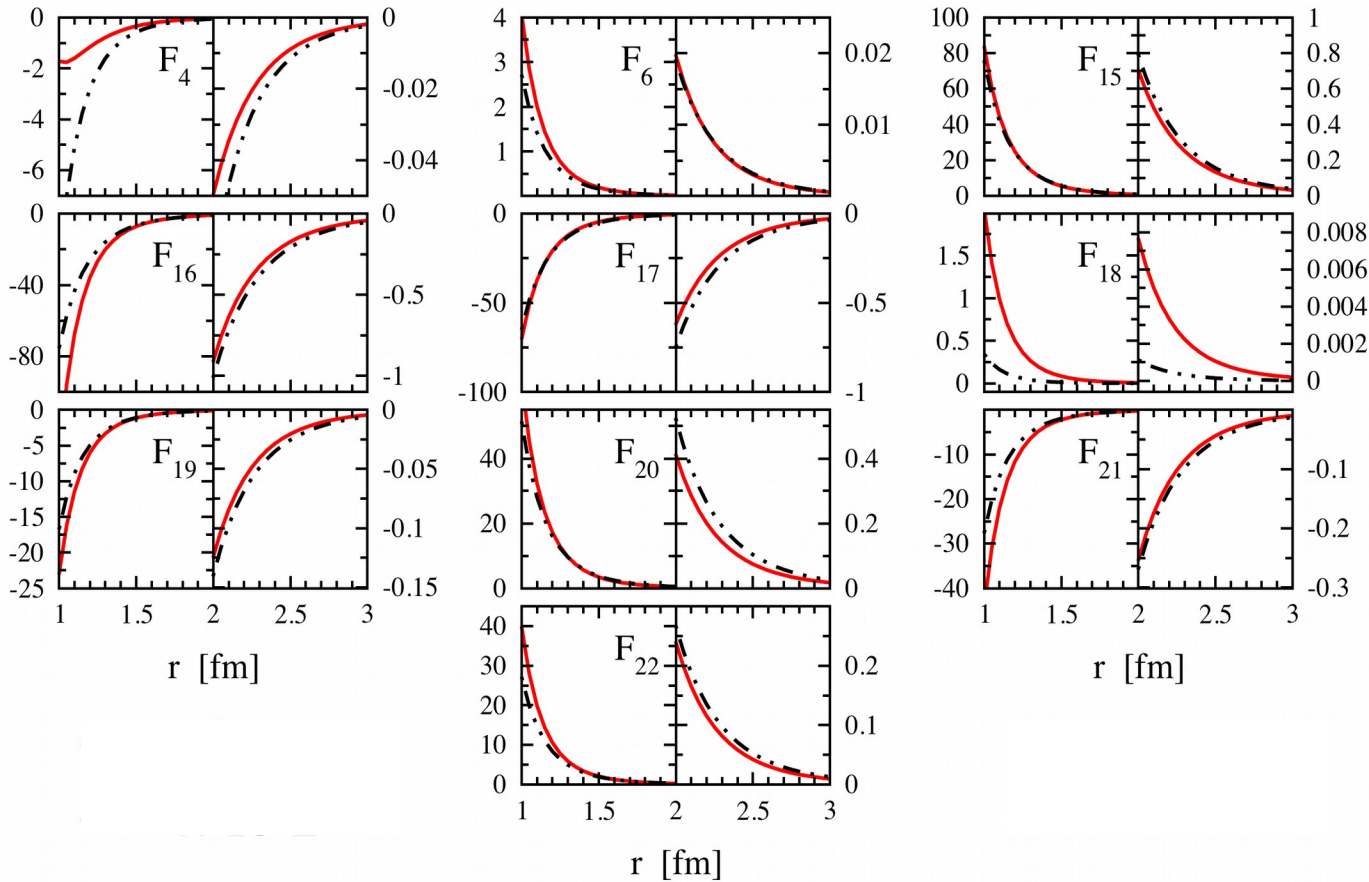
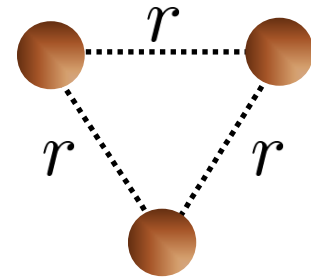


$$V_{3N} = \sum_{i=1}^{22} \tilde{G}_i F_i(r_{12}, r_{23}, r_{31}) + 5\text{perm.}$$

# Two-pion-exchange 3NF in $\Delta$ -full and $\Delta$ -less approach (preliminary)

Krebs, Gasparyan, Epelbaum, in preparation

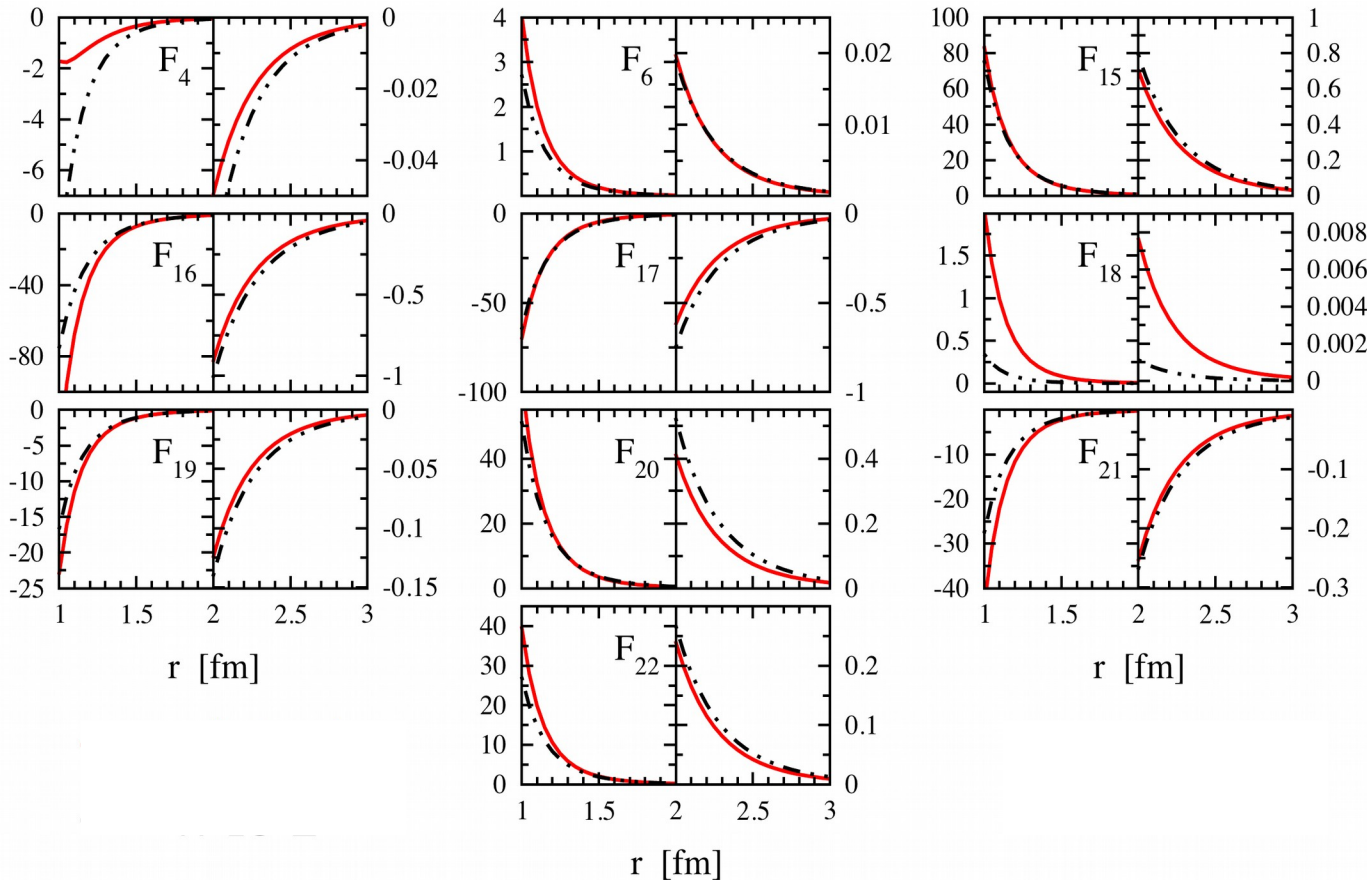
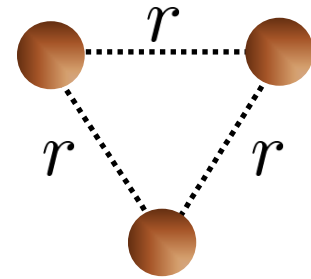
TPE “structure functions”  $F_i$  in MeV”  
in equilateral-triangle configuration



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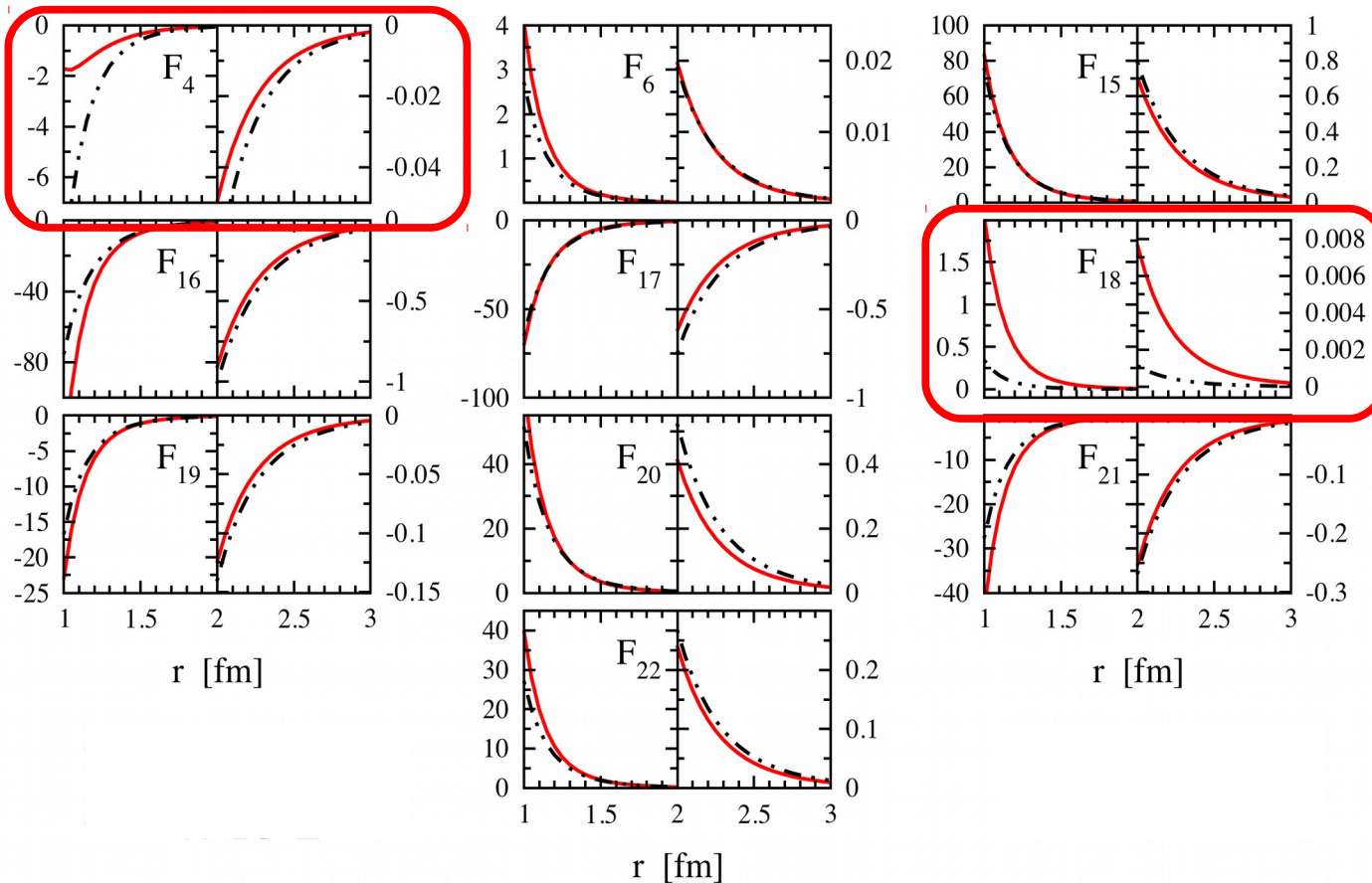
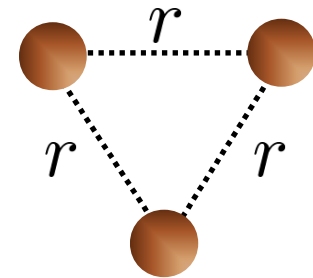
—  $N^4LO$   $\Delta$ -less  
- - -  $N^3LO$ - $\Delta$

→ similar results for large contributions

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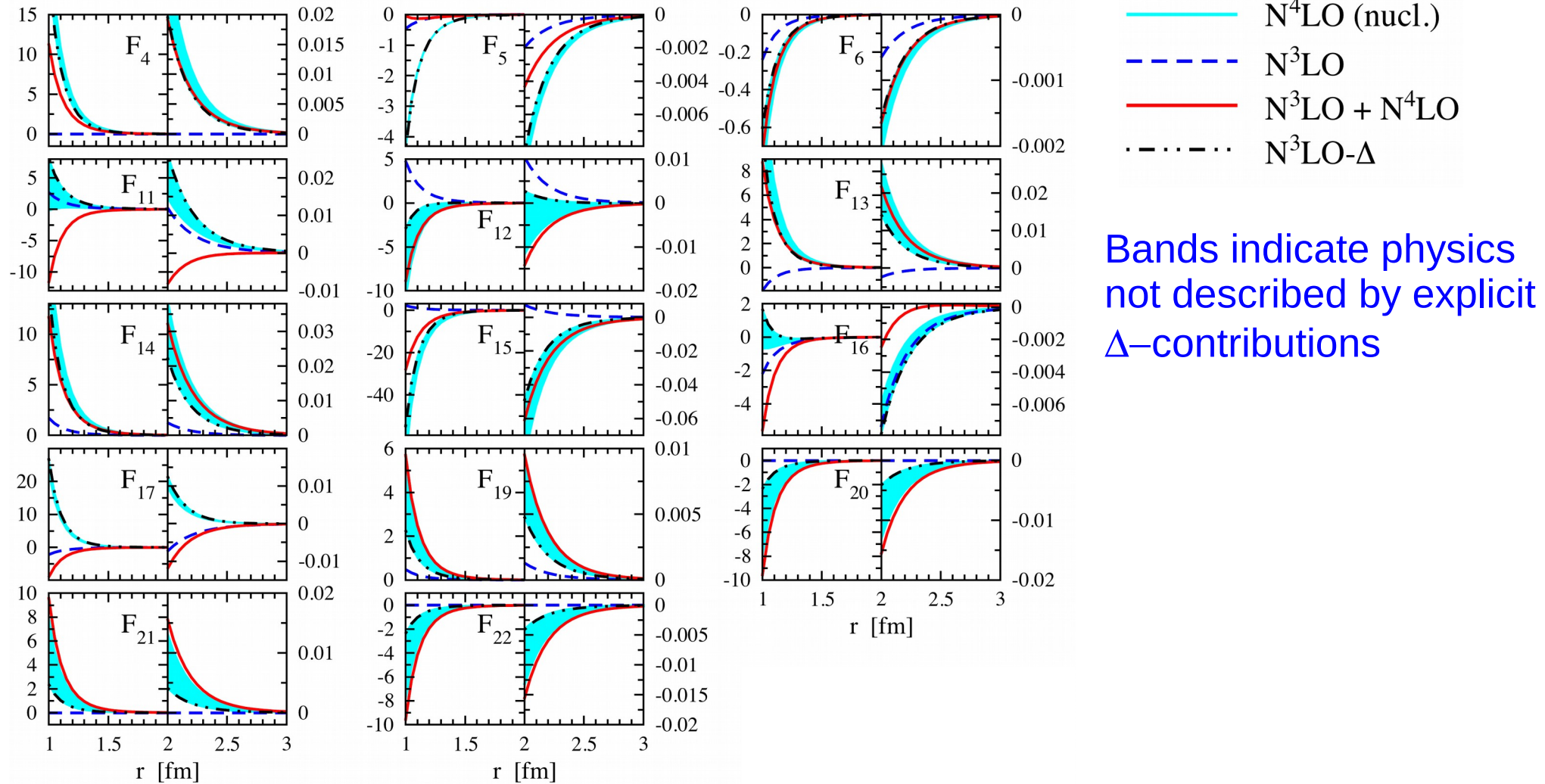


—  $N^4\text{LO } \Delta\text{-less}$   
- - -  $N^3\text{LO-}\Delta$

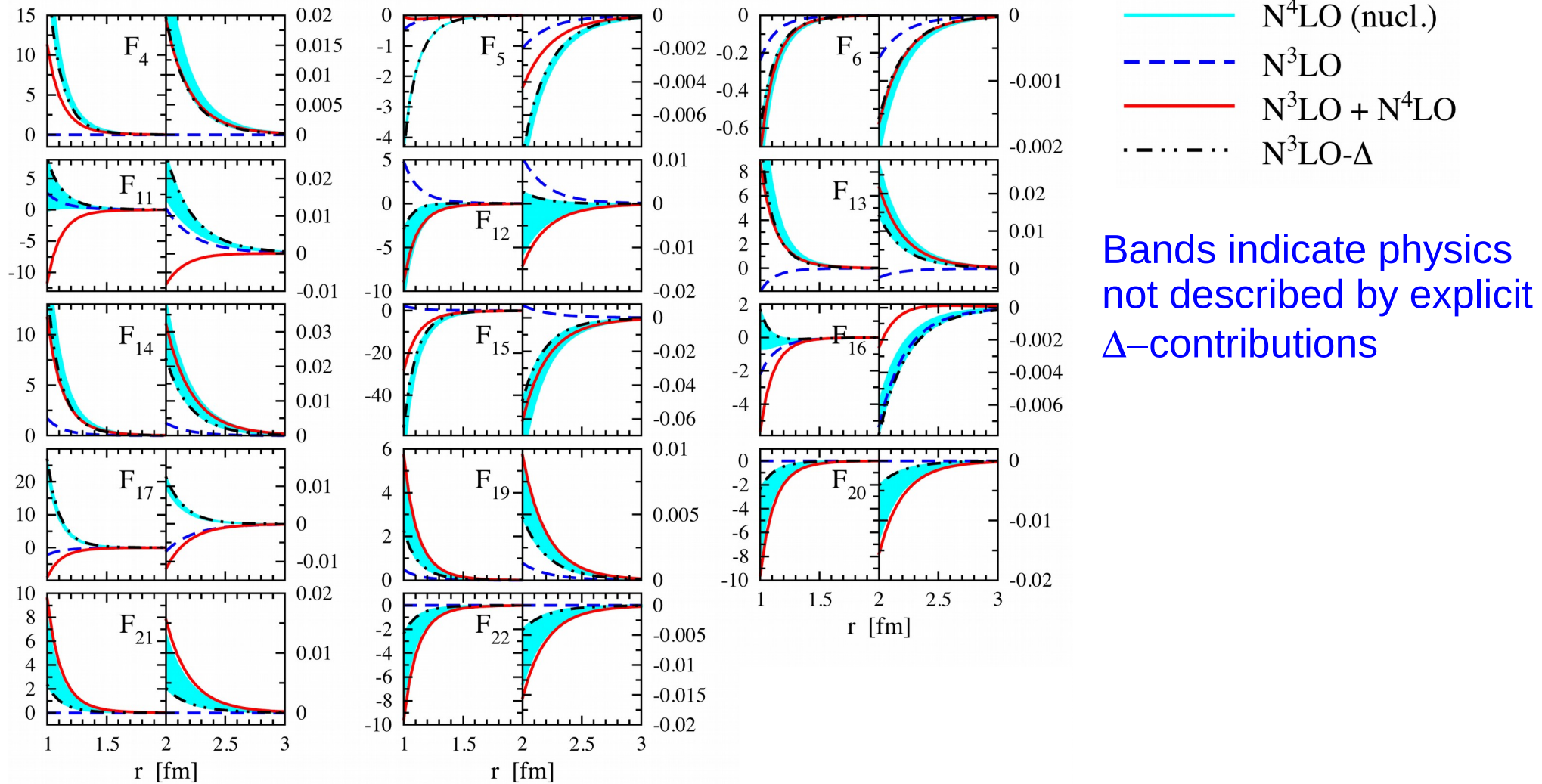
→ similar results for large contributions

→ slightly different for small contributions

# Two-pion-one-pion-exchange 3NF in $\Delta$ -full and $\Delta$ -less approach (preliminary)

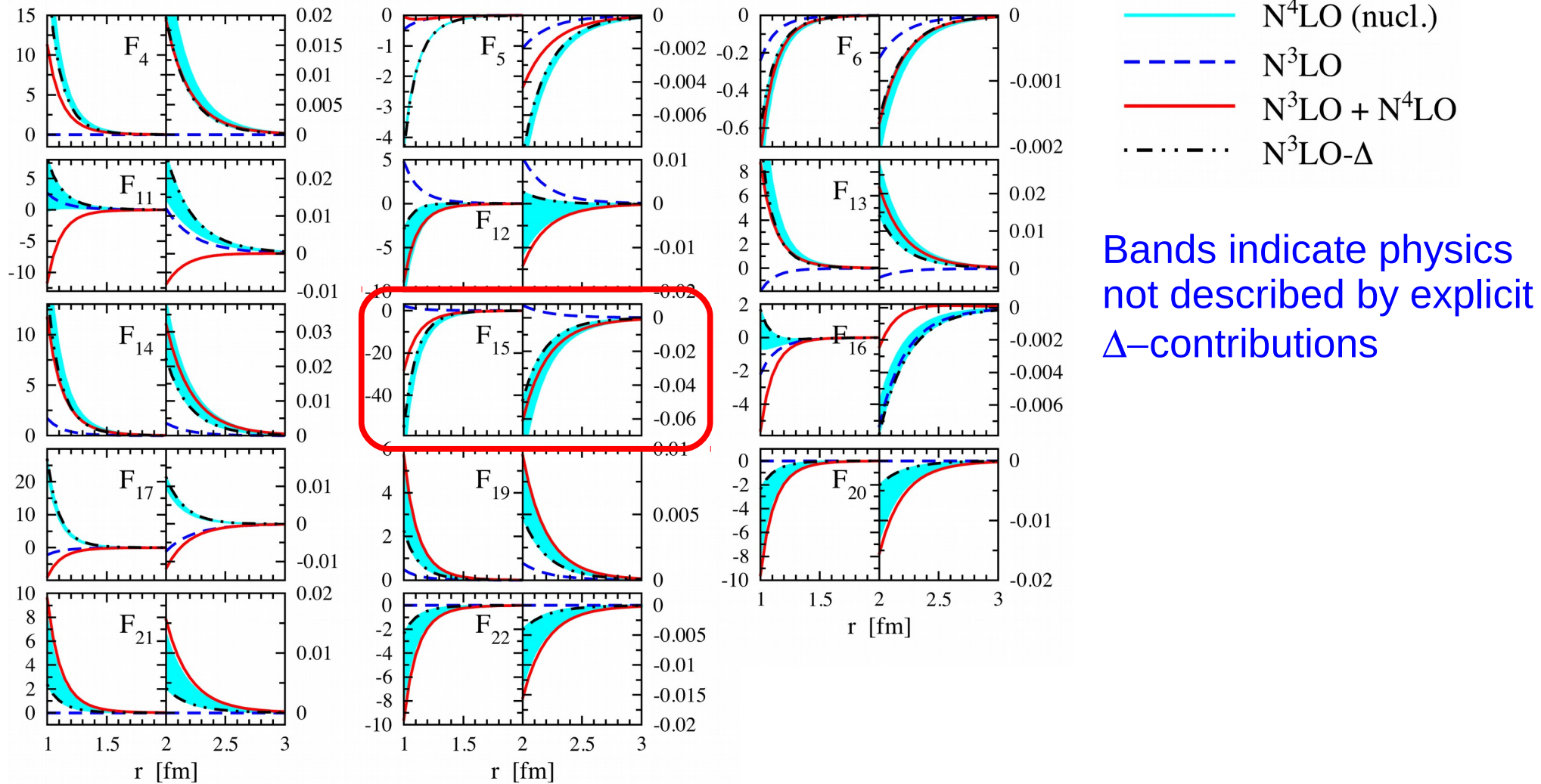


# Two-pion-one-pion-exchange 3NF in $\Delta$ -full and $\Delta$ -less approach (preliminary)



→ Dominant effects come from  $N^3\text{LO}-\Delta/N^4\text{LO}$

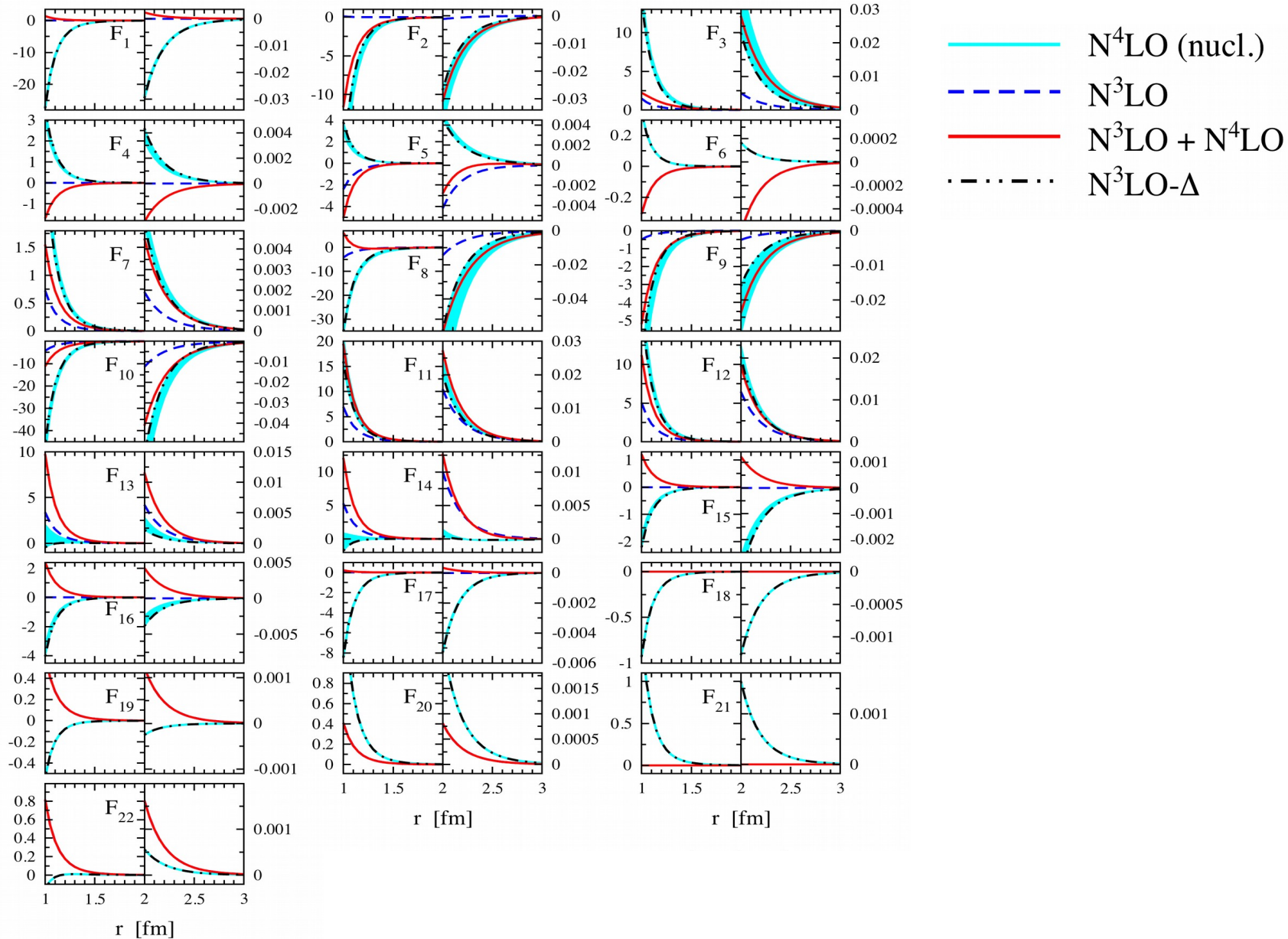
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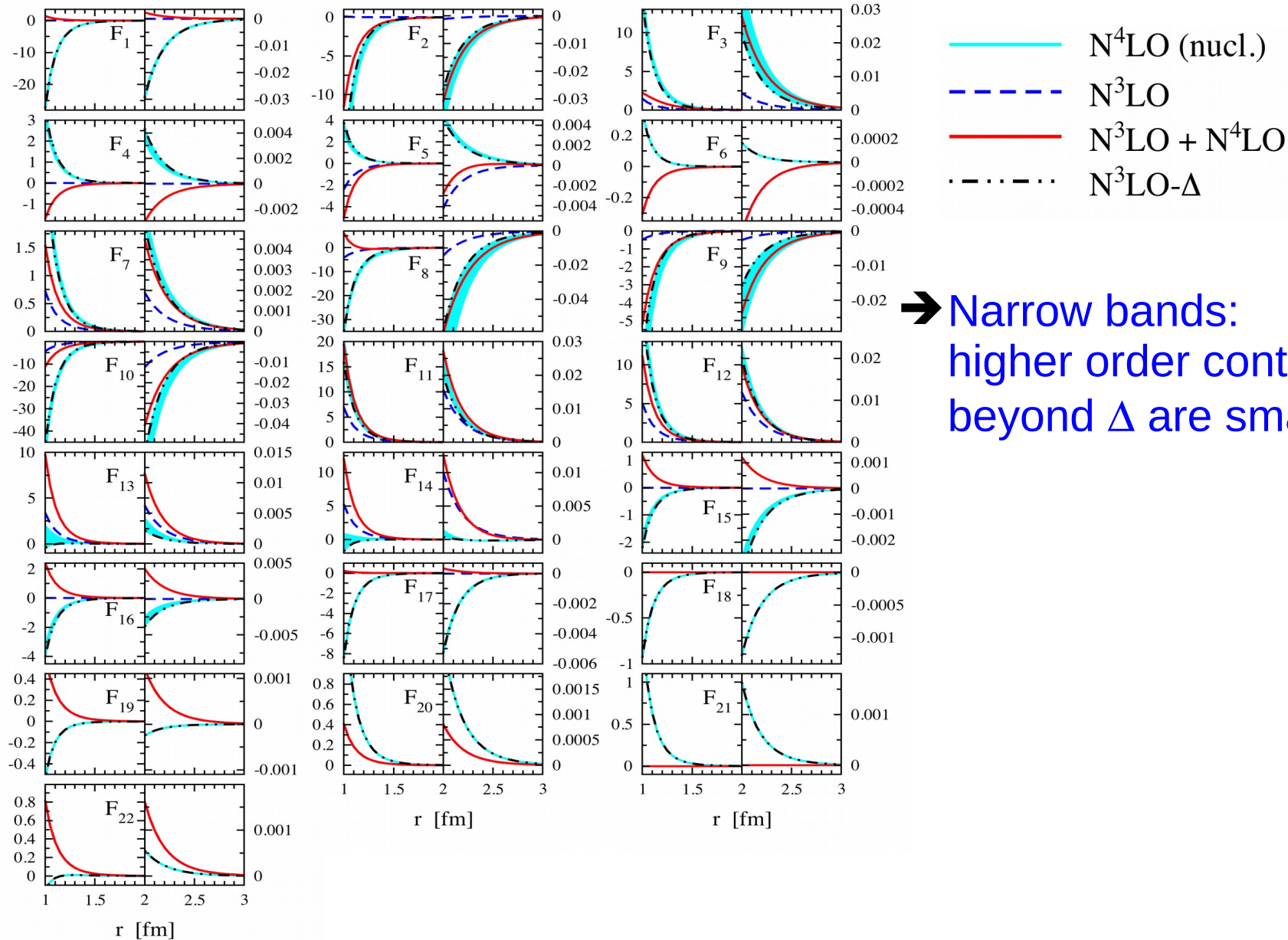
→ The largest  $N^4\text{LO}$  contribution is saturated by  $\Delta$

# Ring-topology 3NF in $\Delta$ -full and $\Delta$ -less approach (preliminary)

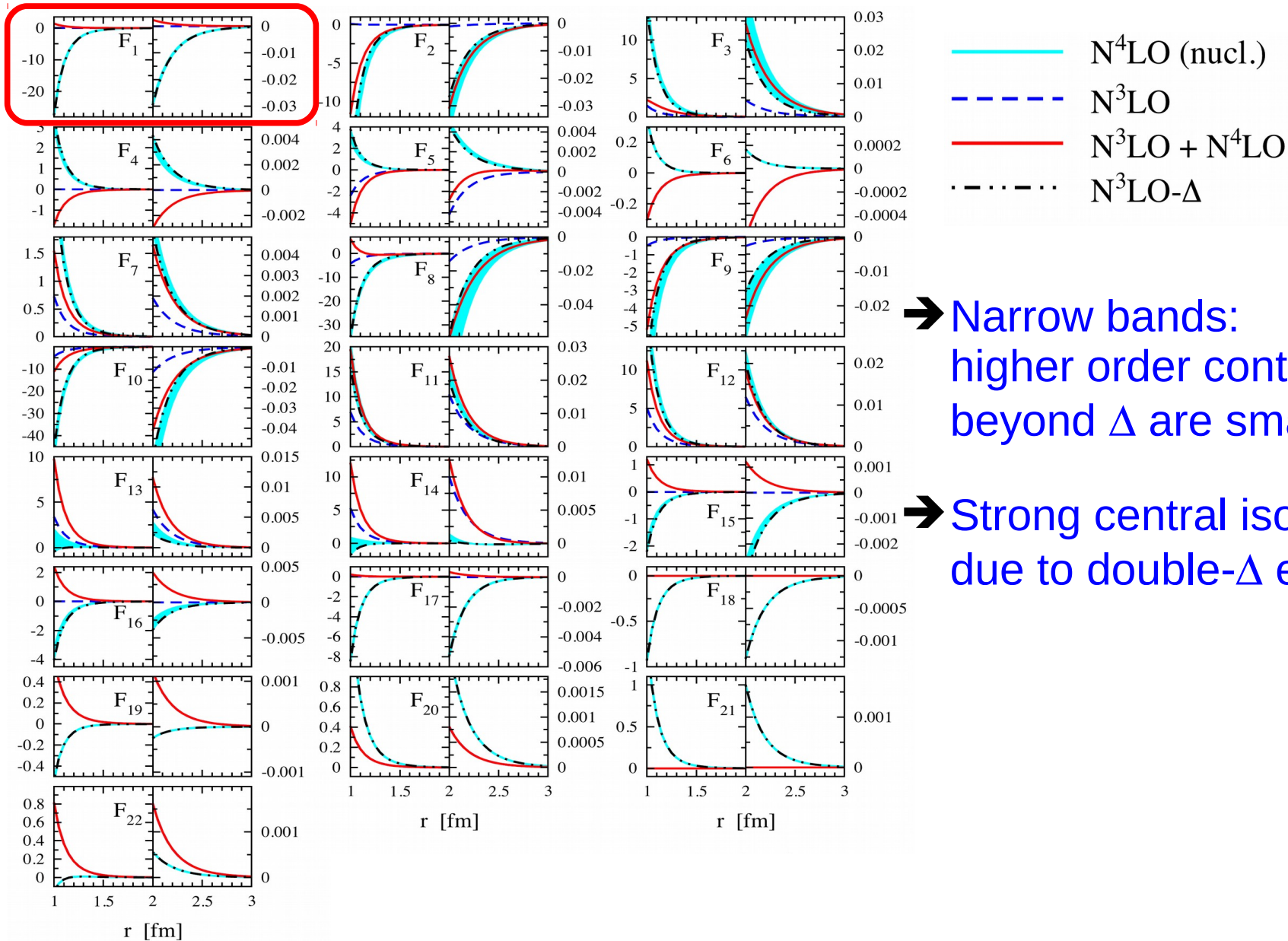




# Ring-topology 3NF in $\Delta$ -full and $\Delta$ -less approach (preliminary)



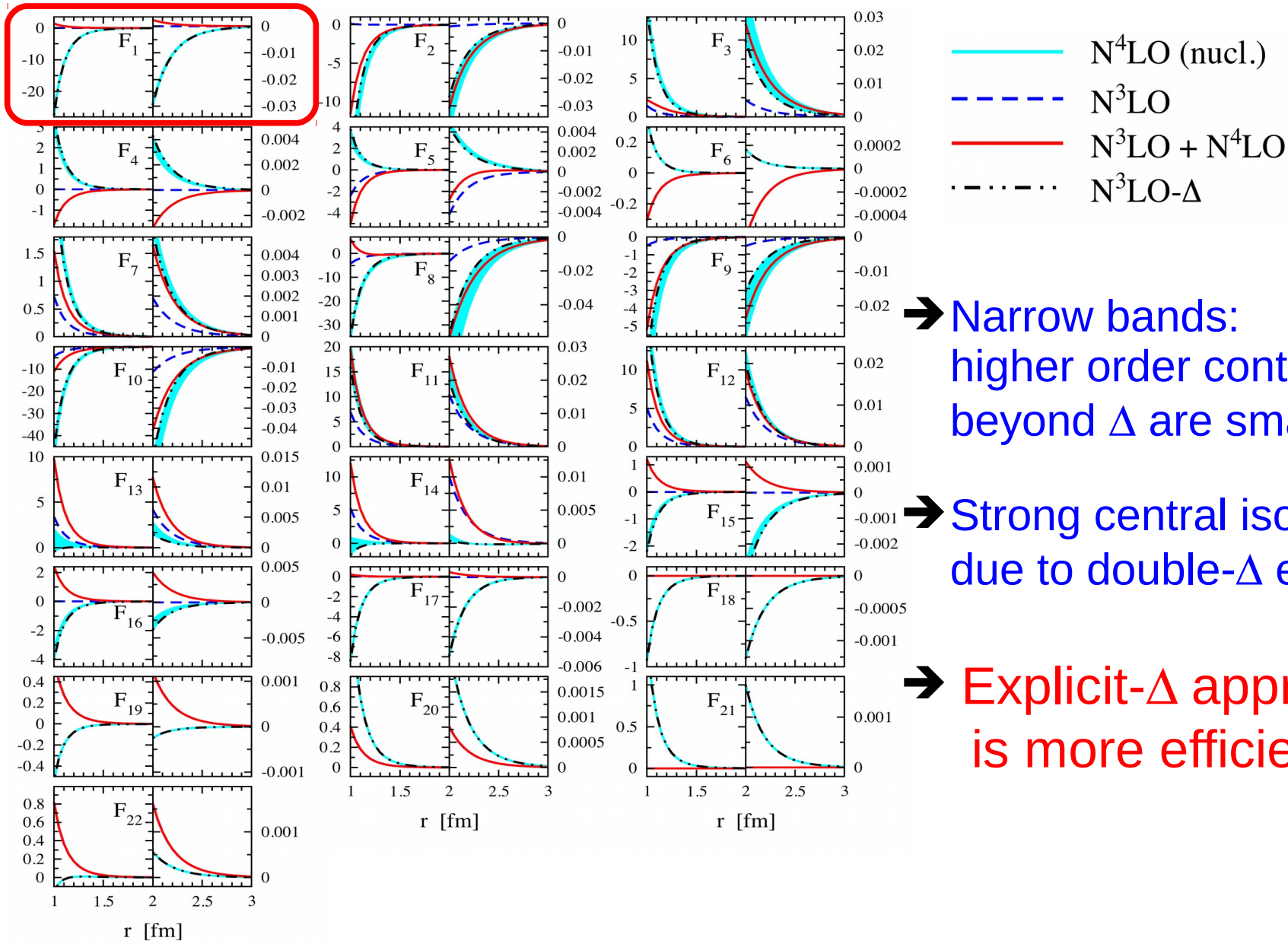
# Ring-topology 3NF in $\Delta$ -full and $\Delta$ -less approach (preliminary)



→ Narrow bands:  
higher order contributions  
beyond  $\Delta$  are small

→ Strong central isoscalar 3NF  
due to double- $\Delta$  excitation

# Ring-topology 3NF in $\Delta$ -full and $\Delta$ -less approach (preliminary)



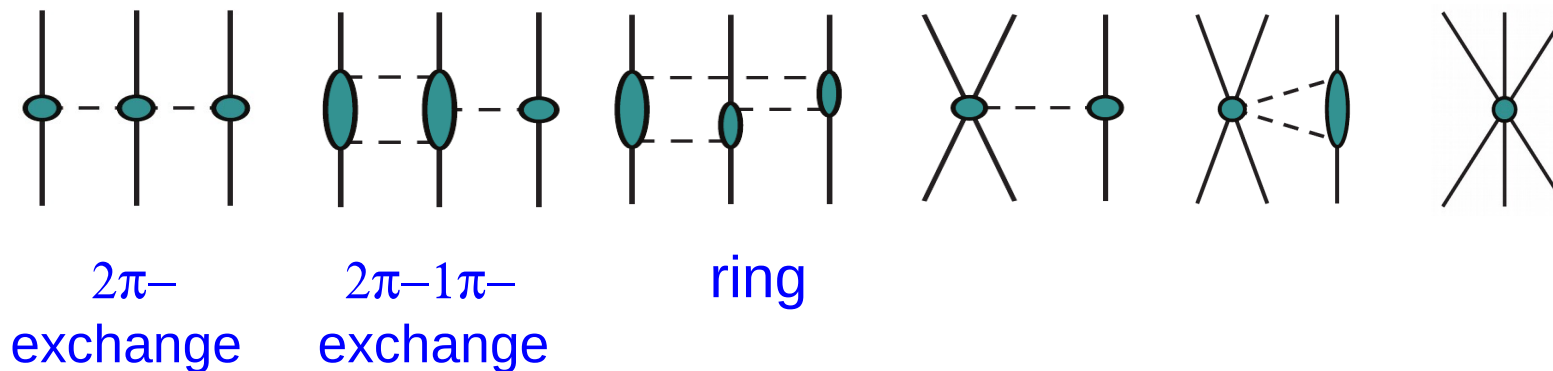
→ Narrow bands:  
higher order contributions  
beyond  $\Delta$  are small

→ Strong central isoscalar 3NF  
due to double- $\Delta$  excitation

→ Explicit- $\Delta$  approach  
is more efficient !

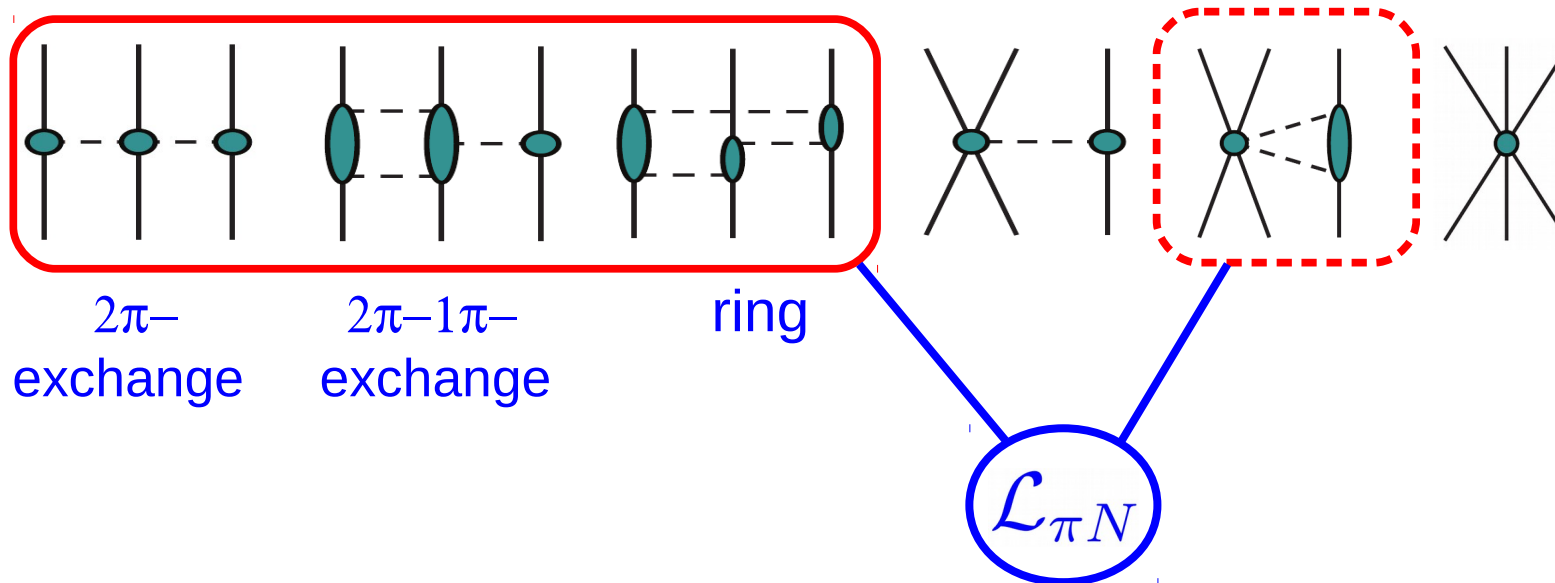
# $\pi$ N input for 3-Nucleon Forces

- Longest-range contributions
- Intermediate-range contributions
- Short-range contributions



# $\pi N$ input for 3-Nucleon Forces

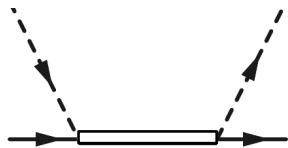
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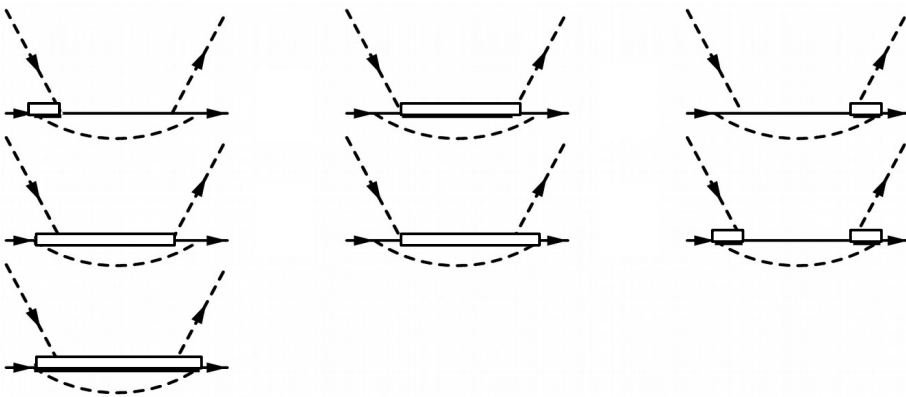
# $\pi N$ scattering up to $\varepsilon^4$

Siemens et al. In preparation

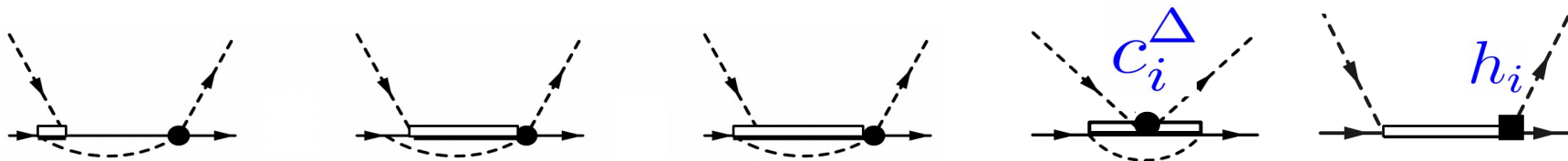
$\varepsilon^1$



$\varepsilon^3$



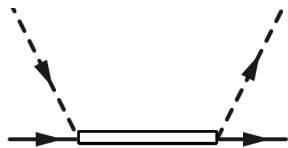
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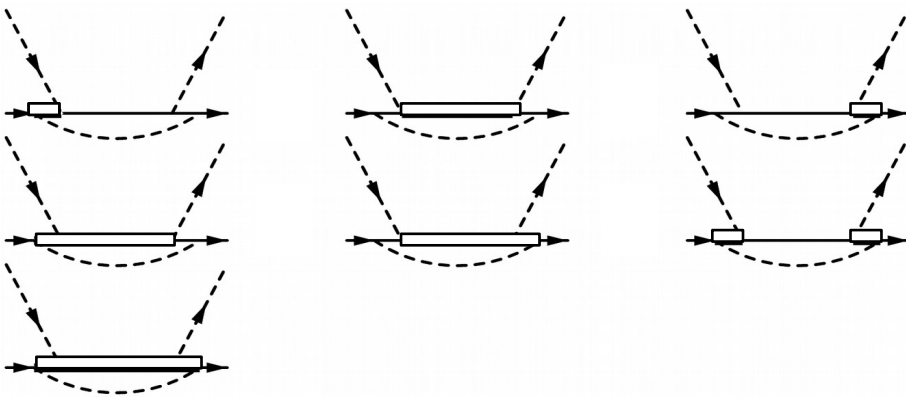
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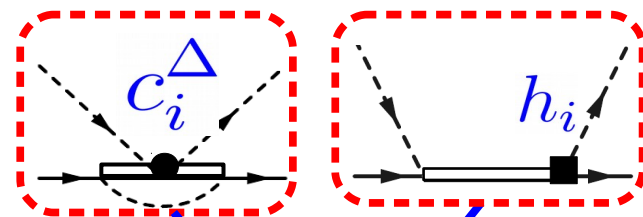
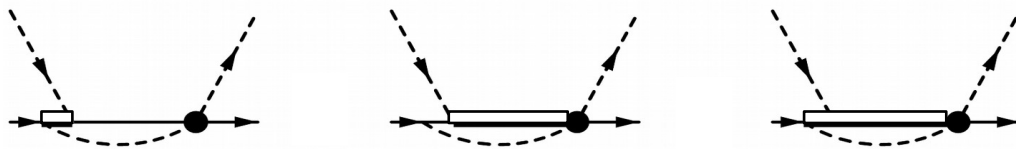
$\varepsilon^1$



$\varepsilon^3$



$\varepsilon^4$

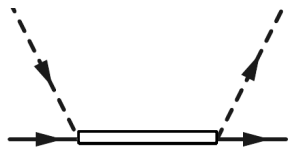


redundant, can be absorbed by redefining other LEC's

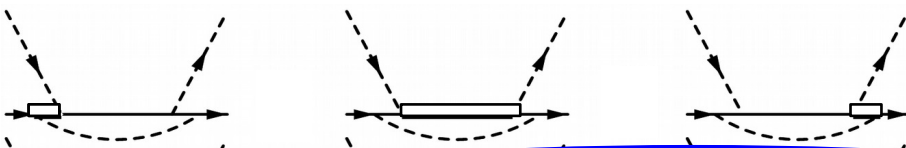
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Siemens et al. In preparation

$\epsilon^1$

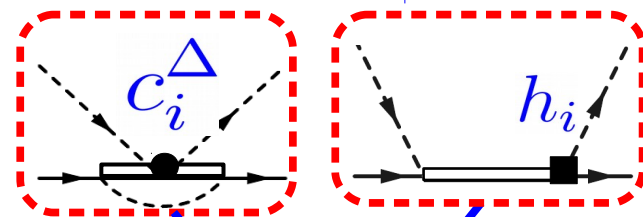
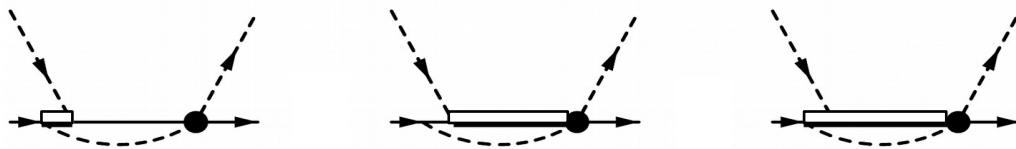


$\epsilon^3$



only two new parameters

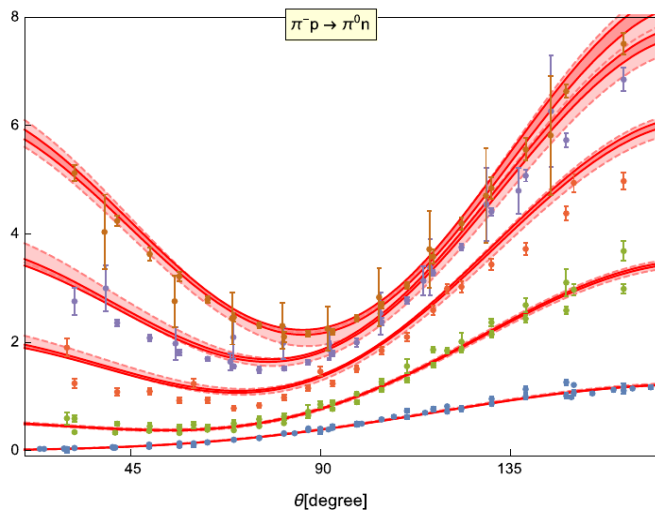
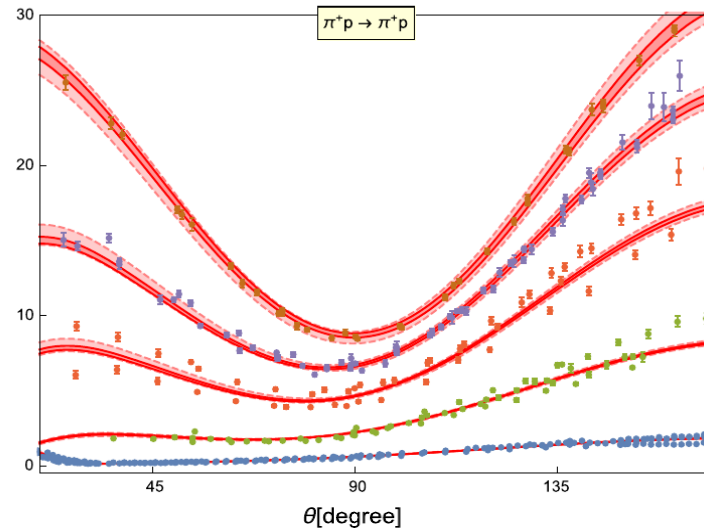
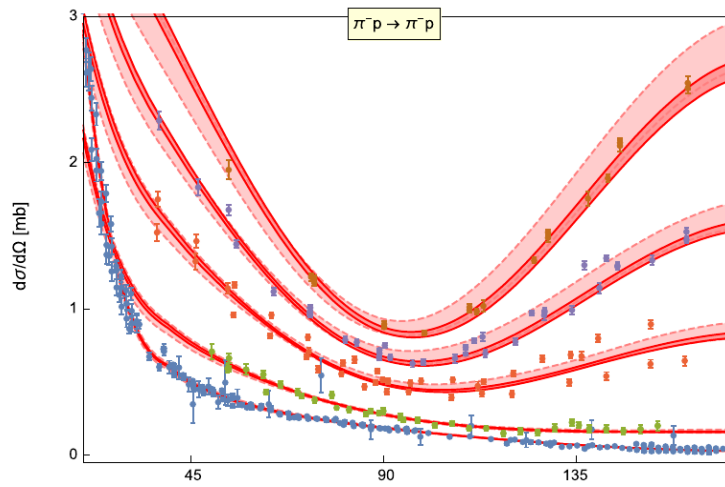
$\epsilon^4$



redundant, can be absorbed  
by redefining other LEC's



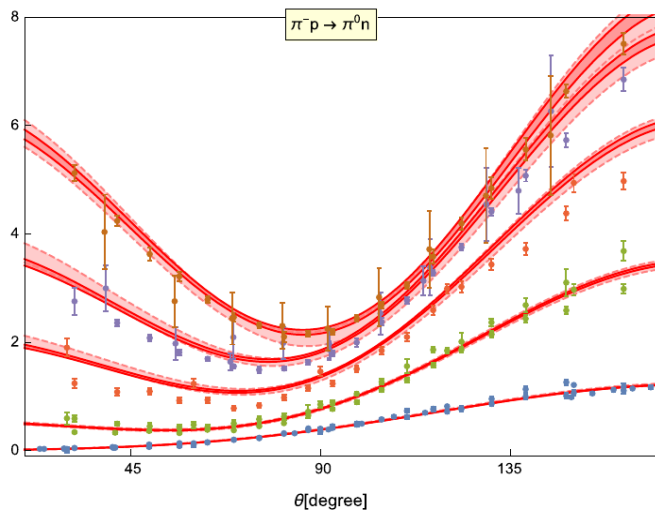
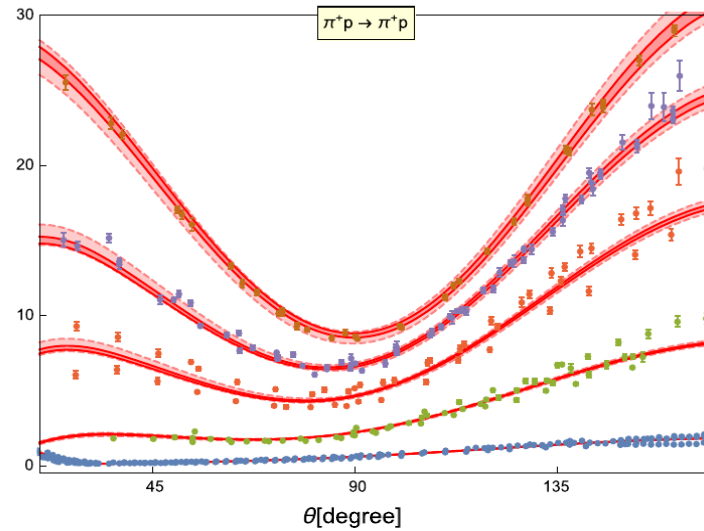
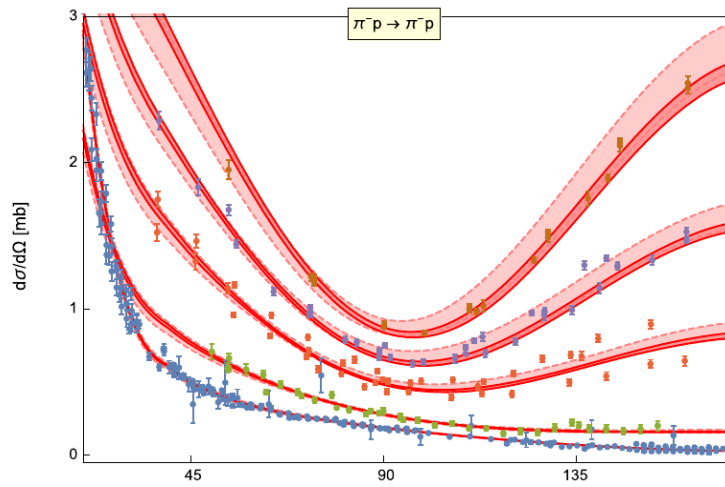
# $\pi N$ differential cross section



- $T_\pi = 167 \pm 5$  MeV
- $T_\pi = 140 \pm 5$  MeV
- $T_\pi = 121 \pm 5$  MeV
- $T_\pi = 90 \pm 5$  MeV
- $T_\pi = 42 \pm 5$  MeV

---  $\varepsilon^3$   
—  $\varepsilon^4$

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---  $\varepsilon^3$

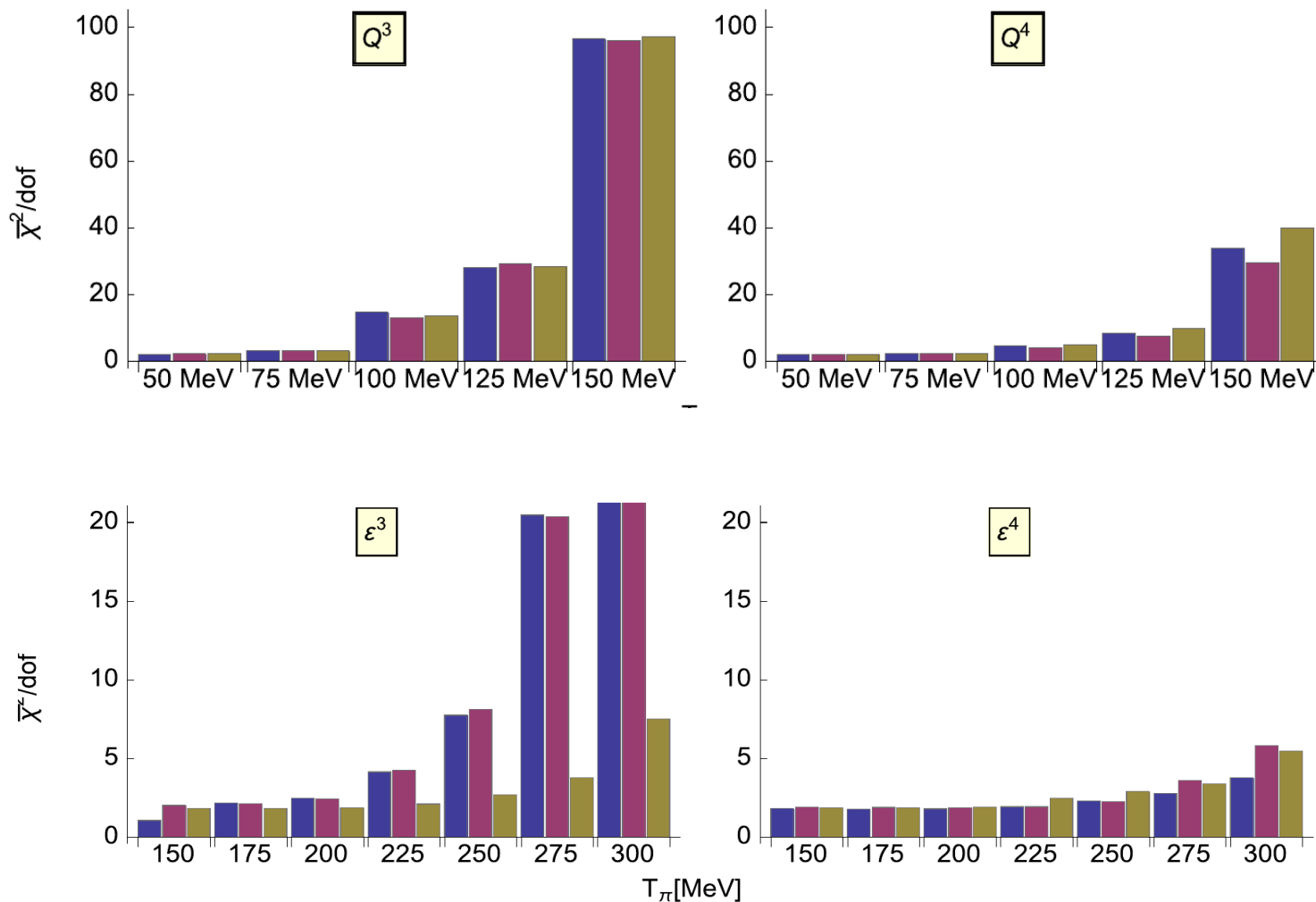
—  $\varepsilon^4$

➔ Theoretical error-bands are narrower

# Quality of the fit to $\pi N$ data

in the  $\Delta$ -less and  $\Delta$ -full  $\chi$ PT (without theoretical errors)

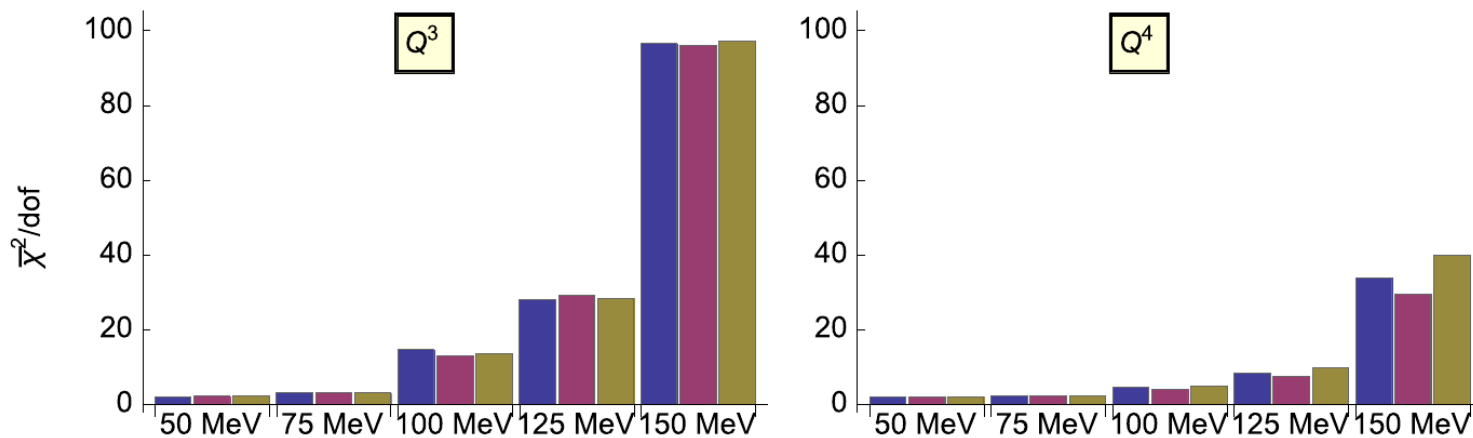
HB-NN HB- $\pi$ N covariant



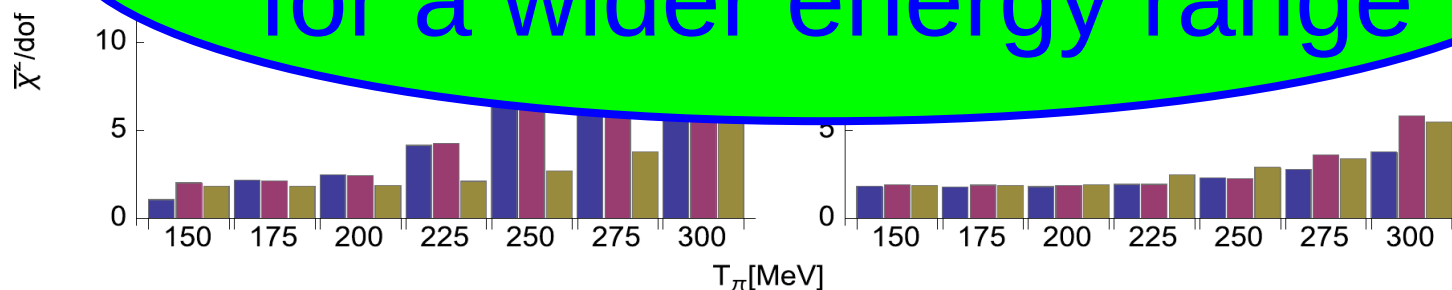
# Quality of the fit to $\pi N$ data

in the  $\Delta$ -less and  $\Delta$ -full  $\chi$ PT (without theoretical errors)

HB-NN HB- $\pi N$  covariant



sizable reduction of the  $\chi^2$   
for a wider energy range



# Summary

- Preliminary results for  $\Delta$ -full chiral 2-nucleon and 3-nucleon forces at  $N^3\text{LO}$  are presented
- 2-nucleon forces (peripheral phases): significant improvement compared to the  $\Delta$ -less case
- 3-nucleon forces: indication of a better convergence; sizable  $\Delta$ -contributions missing in  $\Delta$ -less  $N^4\text{LO}$  3NF  $\sim O(1/\Delta^2)$
- New results for  $\pi\text{N}$  scattering at order  $\varepsilon^4$ : much better fit to data

## Outlook

- Completing construction of  $\Delta$ -full chiral 2N and 3N forces at  $N^3\text{LO}$  and moving forward to even more precise nuclear forces.