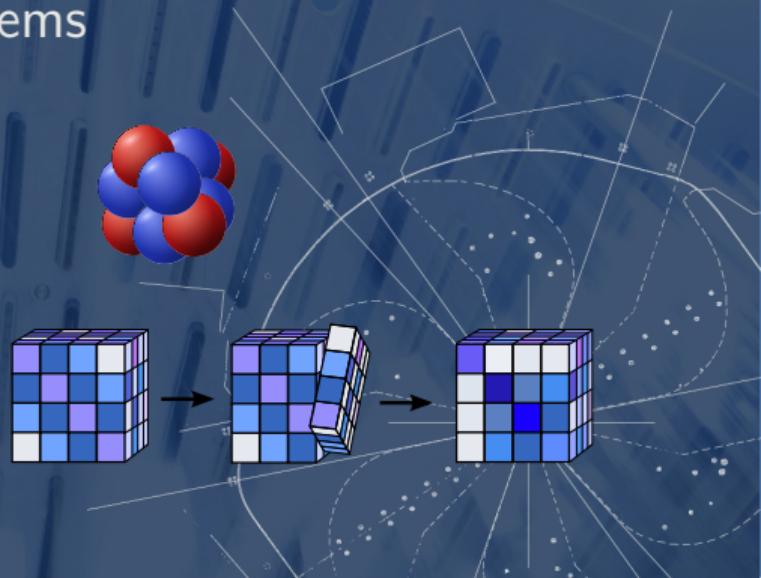


In-medium SRG for fully open-shell systems

Ragnar Stroberg

TRIUMF

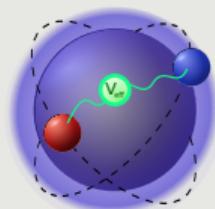
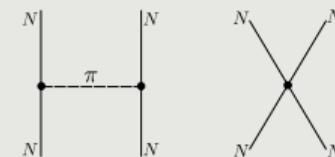
Ab initio workshop
TRIUMF
March 2, 2017





Outline

1. In-medium SRG
2. Ensemble normal ordering
3. Non-standard valence spaces



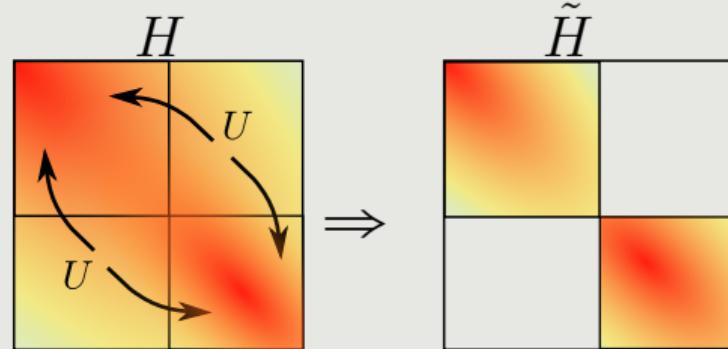
- $H|\Psi\rangle = E|\Psi\rangle$ is too difficult to solve.
- Perform unitary transformation $\tilde{H} = UHU^\dagger$ (implicit change of basis) so SE is easier to solve.

- Iterative/guess-and-check approach.

$$U \equiv e^\Omega = e^{\Omega_n}e^{\Omega_{n-1}}\dots e^{\Omega_2}e^{\Omega_1}$$

- Alternatively, $\Omega_n \rightarrow \eta ds \Rightarrow$ flow equation

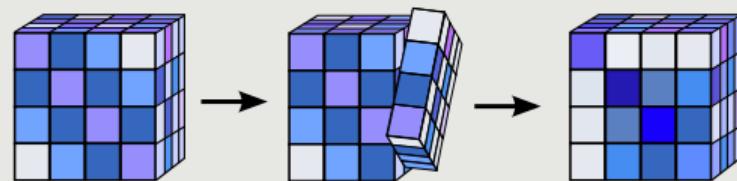
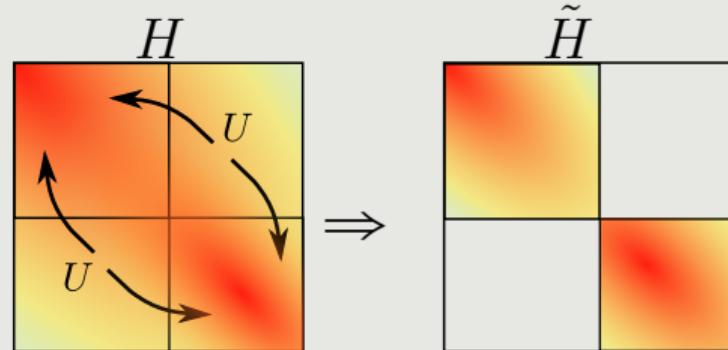
- Computational effort dominated by commutator evaluation.



- $H|\Psi\rangle = E|\Psi\rangle$ is too difficult to solve.
- Perform unitary transformation $\tilde{H} = UHU^\dagger$ (implicit change of basis) so SE is easier to solve.
- Iterative/guess-and-check approach.

$$U \equiv e^\Omega = e^{\Omega_n}e^{\Omega_{n-1}}\dots e^{\Omega_2}e^{\Omega_1}$$

- Alternatively, $\Omega_n \rightarrow \eta ds \Rightarrow$ flow equation
- Computational effort dominated by commutator evaluation.



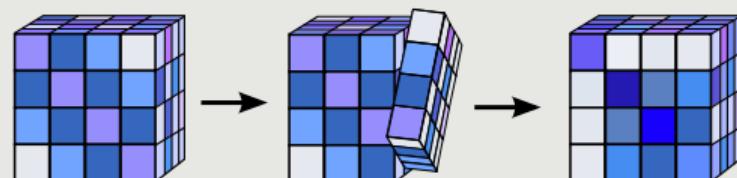
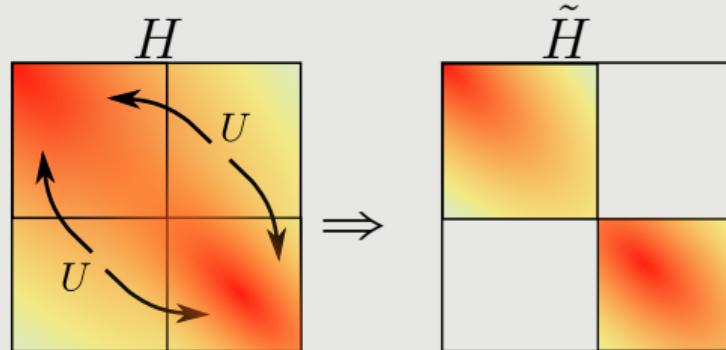
- $H|\Psi\rangle = E|\Psi\rangle$ is too difficult to solve.
- Perform unitary transformation $\tilde{H} = UHU^\dagger$ (implicit change of basis) so SE is easier to solve.
- Iterative/guess-and-check approach.

$$U \equiv e^\Omega = e^{\Omega_n}e^{\Omega_{n-1}}\dots e^{\Omega_2}e^{\Omega_1}$$

- Alternatively, $\Omega_n \rightarrow \eta ds \Rightarrow$ flow equation

$$\frac{dH(s)}{ds} = [\eta(s), H(s)].$$

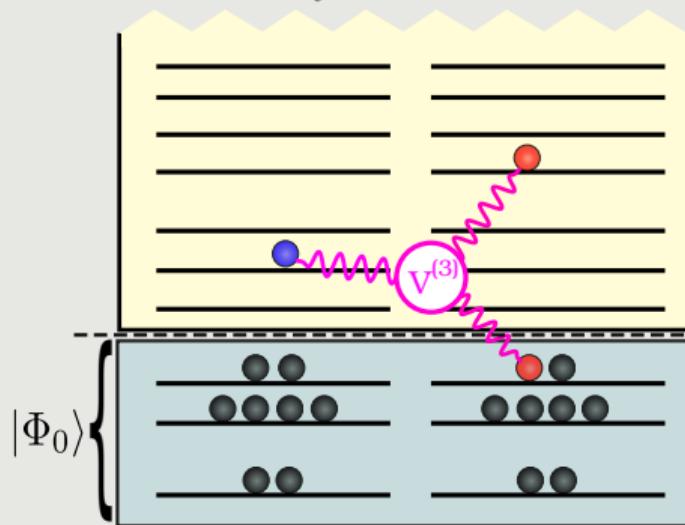
- Computational effort dominated by commutator evaluation.



Why “in-medium”?

$$H = \underbrace{E_0}_{\text{0-body}} + \underbrace{\sum_{ij} H_{ij} \{a_i^\dagger a_j\}}_{\text{1-body}} + \underbrace{\frac{1}{4} \sum_{ijkl} H_{ijkl} \{a_i^\dagger a_j^\dagger a_l a_k\}}_{\text{2-body}} + \underbrace{\frac{1}{36} \sum_{ijklmn} H_{ijklmn} \{a_i^\dagger a_j^\dagger a_k^\dagger a_n a_m a_l\}}_{\text{3-body}} + \dots$$

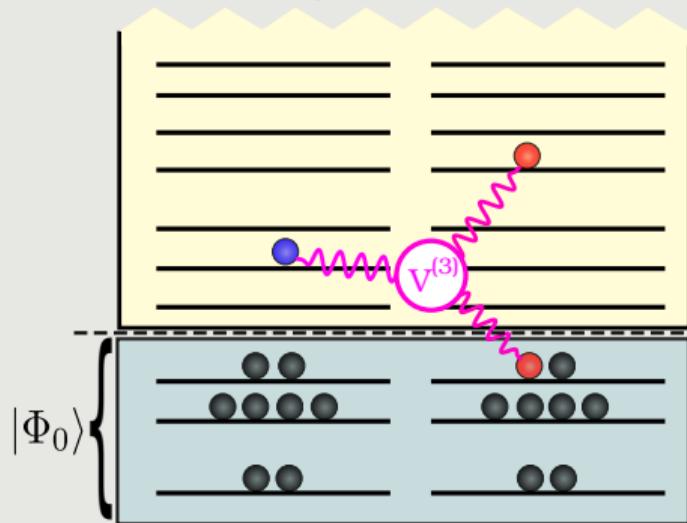
- In general, the transformation U will induce 4-body, 5-body, etc. forces 😱
- Write H in normal-ordered form w.r.t reference $|\Phi_0\rangle$
- $\langle \Phi_0 | \{a_1^\dagger \dots a_N^\dagger a_N \dots a_1\} | \Phi_0 \rangle = 0$
- If $|\Phi_0\rangle \approx |\Psi\rangle$, higher-body terms are negligible
- Truncate all operators at 2 body level



Why “in-medium”?

$$H = \underbrace{E_0}_{\text{0-body}} + \underbrace{\sum_{ij} H_{ij} \{a_i^\dagger a_j\}}_{\text{1-body}} + \underbrace{\frac{1}{4} \sum_{ijkl} H_{ijkl} \{a_i^\dagger a_j^\dagger a_l a_k\}}_{\text{2-body}} + \underbrace{\frac{1}{36} \sum_{ijklmn} H_{ijklmn} \{a_i^\dagger a_j^\dagger a_k^\dagger a_n a_m a_l\}}_{\text{3-body}} + \dots$$

- In general, the transformation U will induce 4-body, 5-body, etc. forces 😱
- Write H in normal-ordered form w.r.t reference $|\Phi_0\rangle$
- $\langle\Phi_0| \{a_1^\dagger \dots a_N^\dagger a_N \dots a_1\} |\Phi_0\rangle = 0$
- If $|\Phi_0\rangle \approx |\Psi\rangle$, higher-body terms are negligible
- Truncate all operators at 2 body level

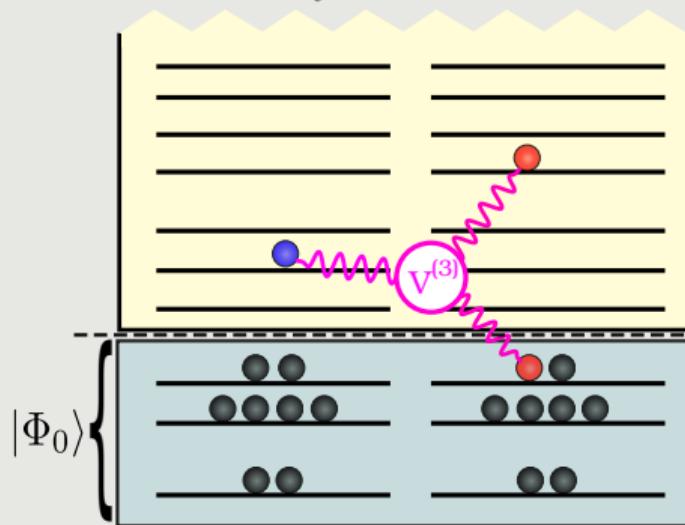


Tskukiyama et al (2011)

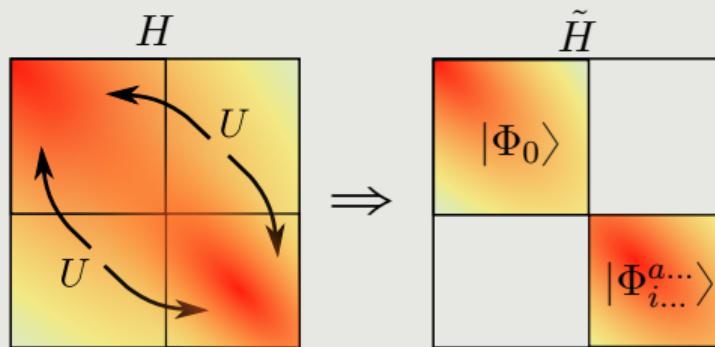
Why “in-medium”?

$$H = \underbrace{E_0}_{\text{0-body}} + \underbrace{\sum_{ij} H_{ij}\{a_i^\dagger a_j\}}_{\text{1-body}} + \underbrace{\frac{1}{4} \sum_{ijkl} H_{ijkl}\{a_i^\dagger a_j^\dagger a_l a_k\}}_{\text{2-body}} + \underbrace{\frac{1}{36} \sum_{ijklmn} H_{ijklmn}\{a_i^\dagger a_j^\dagger a_k^\dagger a_n a_m a_l\}}_{\text{3-body}} + \dots$$

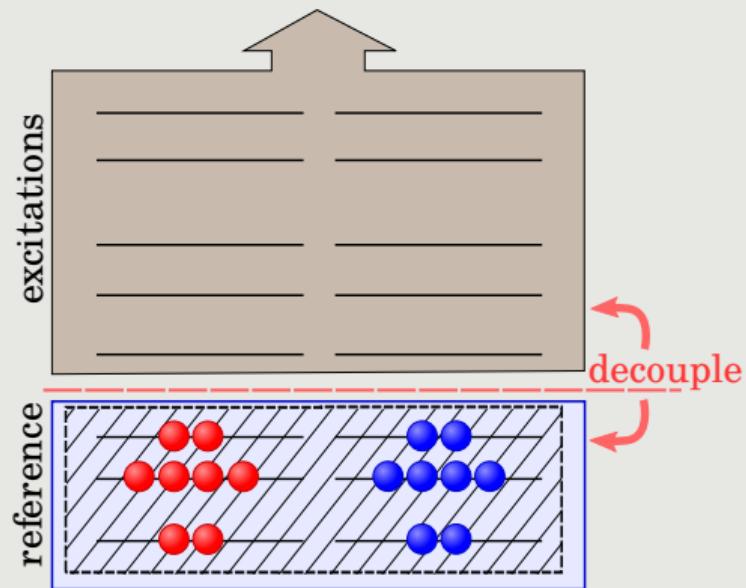
- In general, the transformation U will induce 4-body, 5-body, etc. forces 😱
- Write H in normal-ordered form w.r.t reference $|\Phi_0\rangle$
- $\langle\Phi_0|\{a_1^\dagger \dots a_N^\dagger a_N \dots a_1\}|\Phi_0\rangle = 0$
- If $|\Phi_0\rangle \approx |\Psi\rangle$, higher-body terms are negligible
- **Truncate all operators at 2 body level**



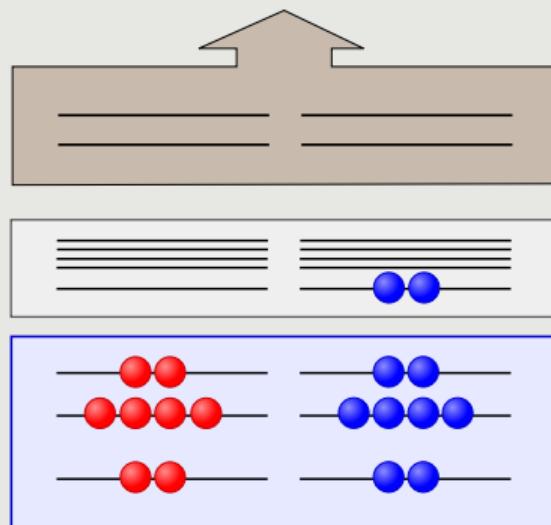
Solving the many-body problem



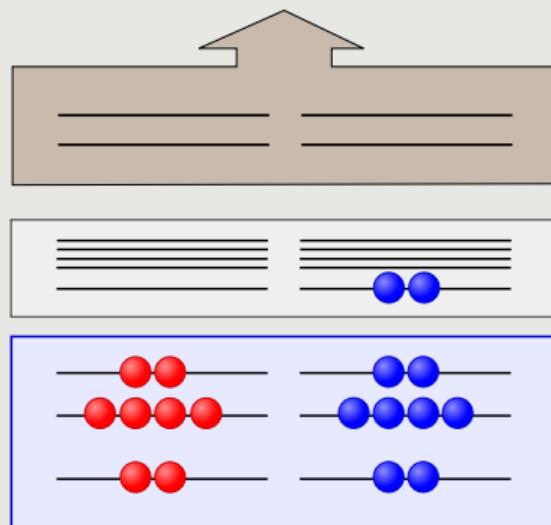
- Decouple a 1×1 sub-block
- Use SRG to suppress excitations out of $|\Phi_0\rangle$
- After decoupling, energy is $E_0 = \langle \Phi_0 | \tilde{H} | \Phi_0 \rangle$



- Open shell systems: multiple (quasi-) degenerate configurations. $|\Phi_0\rangle \approx |\Psi\rangle$
- Single Slater determinant may not have good total angular momentum J
- Large rotation angle \rightarrow induced many-body forces
- Strategies:
 - Break symmetries and restore afterward
 - Construct multi-determinant reference, then decouple (multi-reference IM-SRG)
 - Decouple a subset of determinants, then construct state from them (valence-space IMSRG)

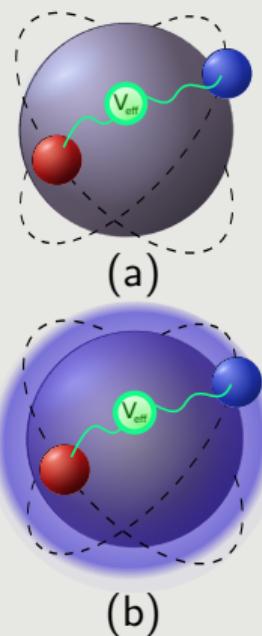
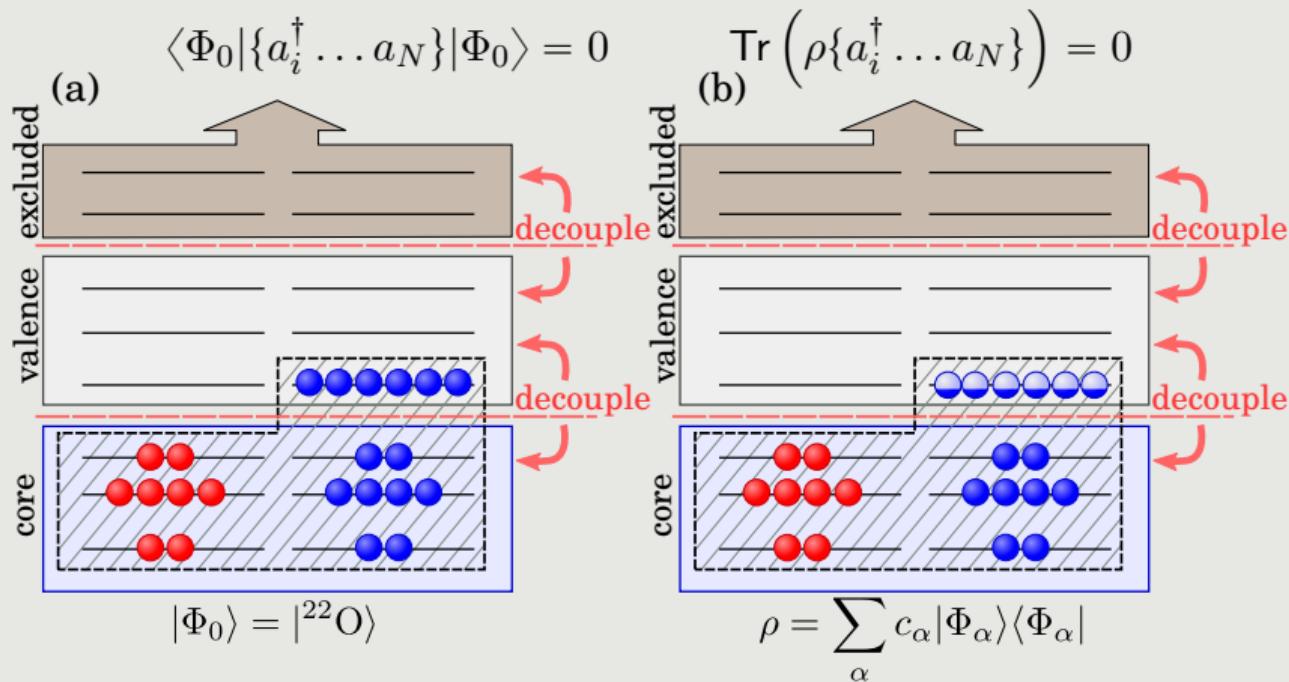


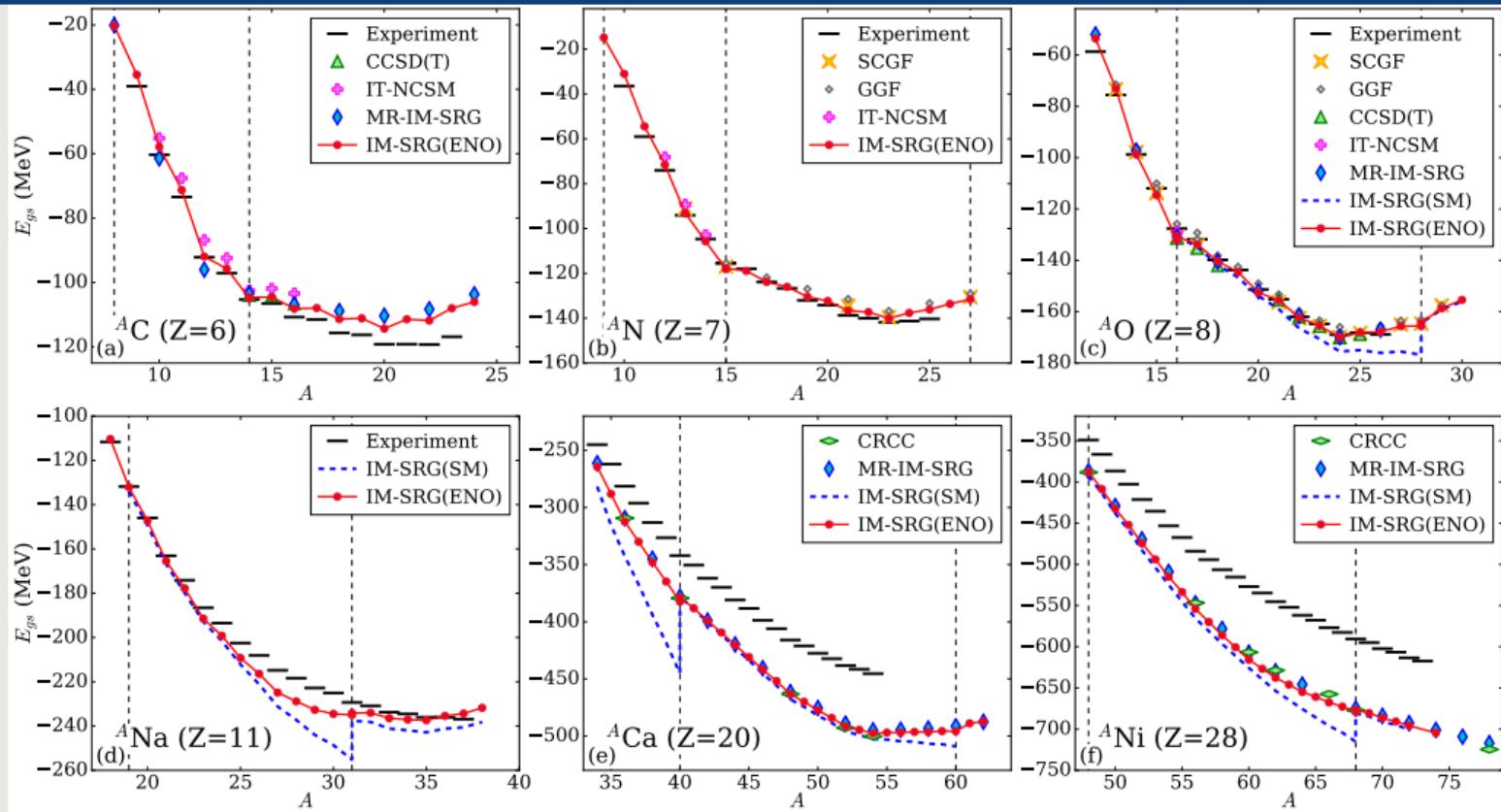
- Open shell systems: multiple (quasi-) degenerate configurations. $|\Phi_0\rangle \approx |\Psi\rangle$
- Single Slater determinant may not have good total angular momentum J
- Large rotation angle \rightarrow induced many-body forces
- Strategies:
 - Break symmetries and restore afterward
 - Construct multi-determinant reference, then decouple (multi-reference IM-SRG)
 - Decouple a subset of determinants, then construct state from them (valence-space IMSRG)

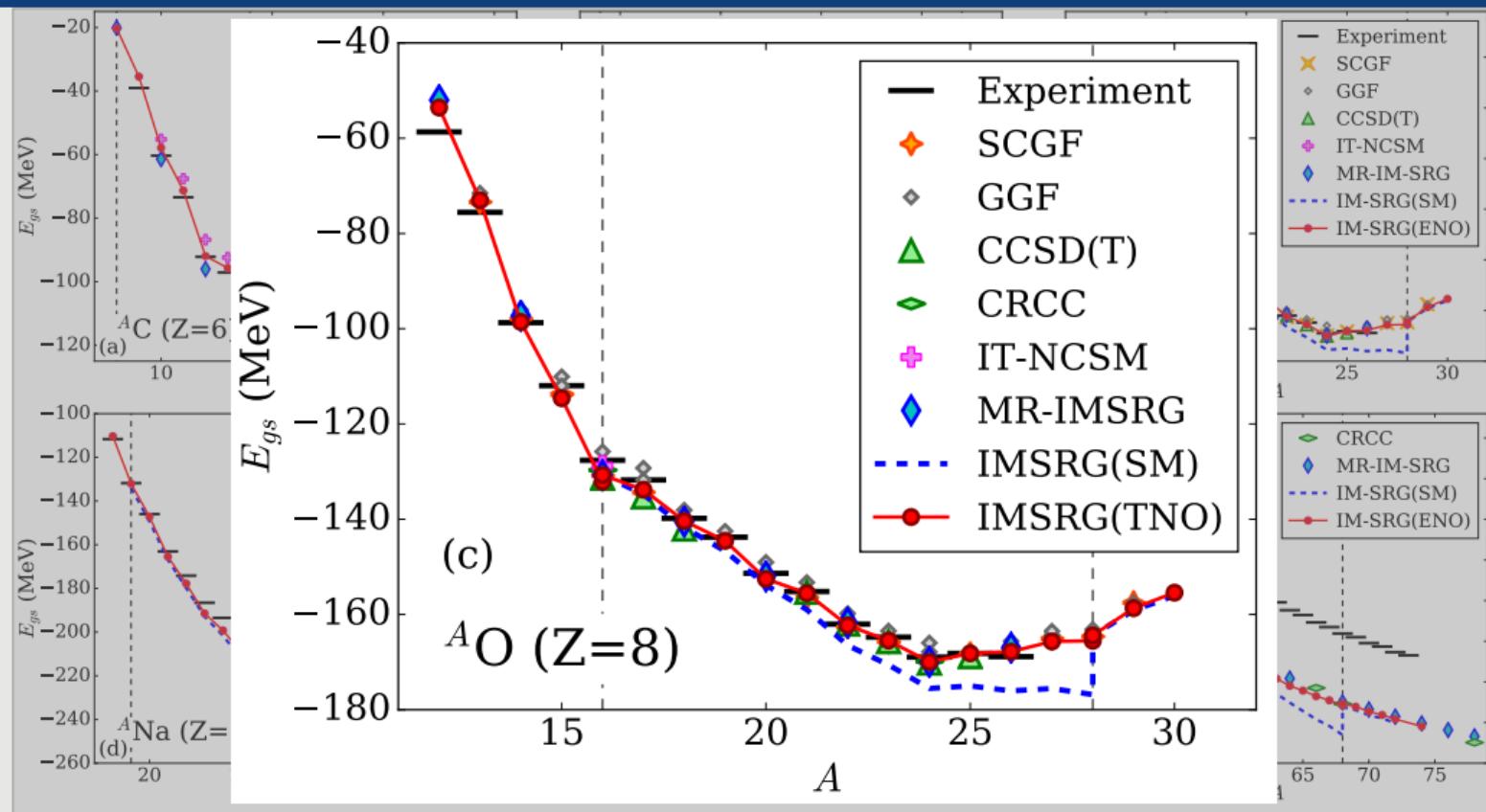


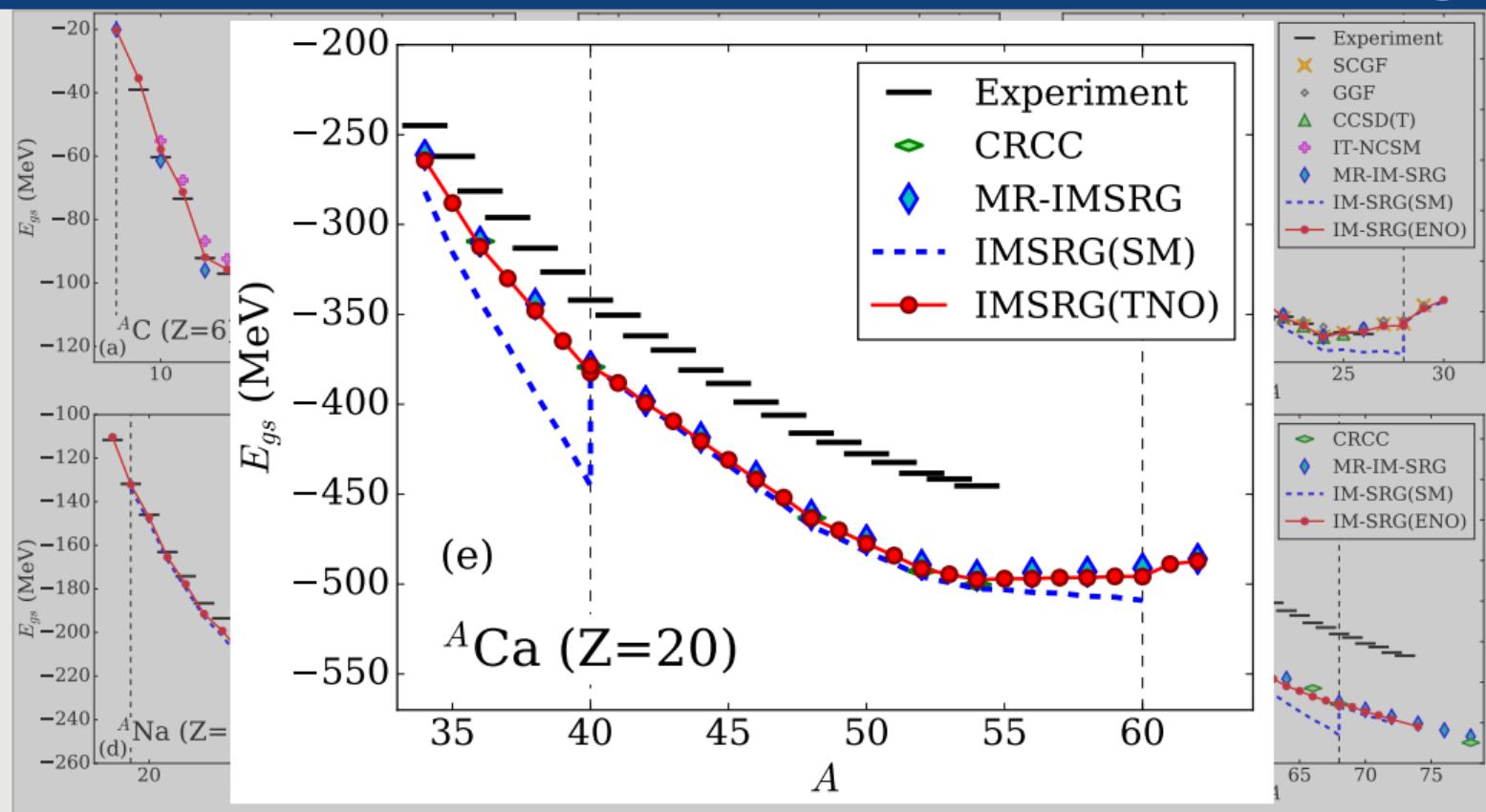
What reference should be used when decoupling a valence space?

$$H = E_0 + \sum_{ij} f_{ij} \{a_i^\dagger a_j\} + \frac{1}{4} \sum_{ijkl} \Gamma_{ijkl} \{a_i^\dagger a_j^\dagger a_l a_k\} + \frac{1}{36} \sum_{ijklmn} W_{ijklmn} \{a_i^\dagger a_j^\dagger a_k^\dagger a_n a_m a_l\}$$

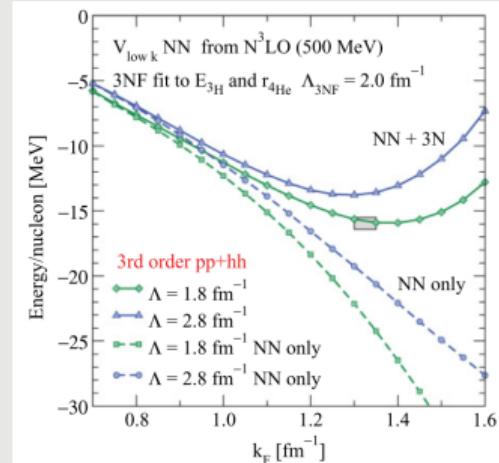




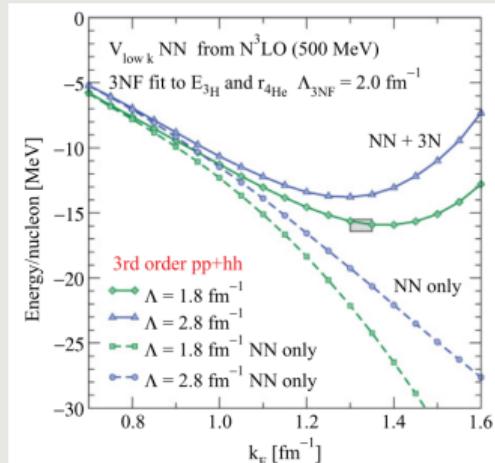
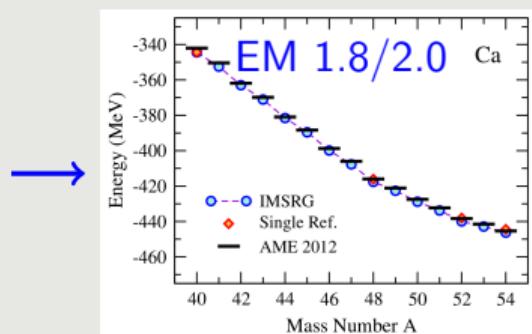
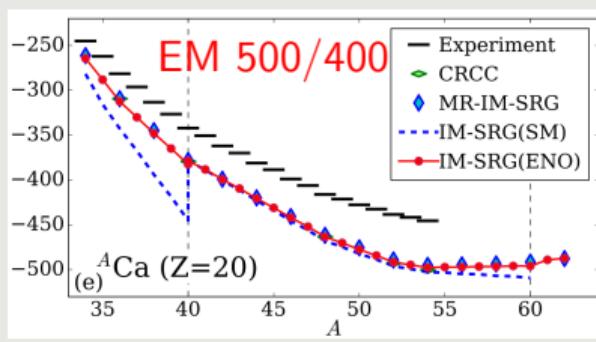




	EM 500/400	EM 1.8/2.0
NN	N ³ LO $\Lambda_{2N} = 500$ MeV non-local regulator fit to NN scattering, ^2H $\lambda_{SRG} = 1.88 \text{ fm}^{-1}$	same same same same \approx same
3N	N ² LO $\Lambda_{3N} = 400$ MeV local regulator fit to ^3H BE, $t_{1/2}$ consistently SRG evolved	same \approx same non-local regulator fit to ^3H BE, ^4He r_{ch} no SRG for 3N

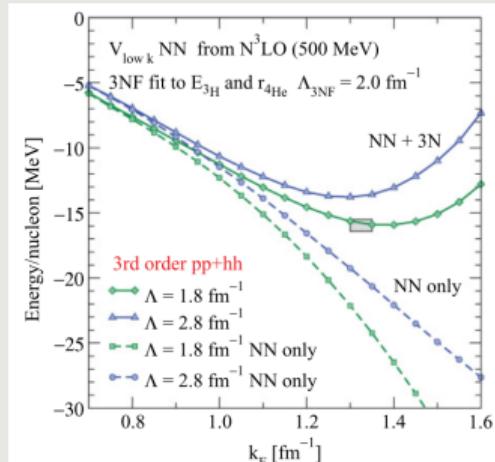
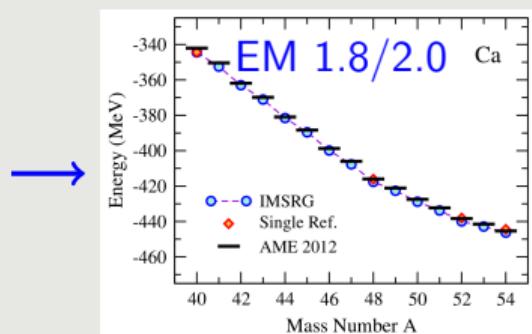
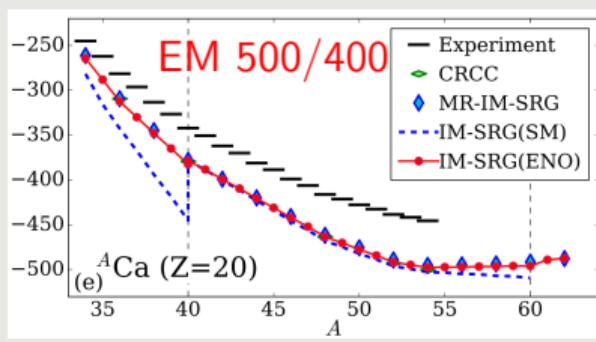


	EM 500/400	EM 1.8/2.0
NN	$N^3\text{LO}$ $\Lambda_{2N} = 500 \text{ MeV}$ non-local regulator fit to NN scattering, ${}^2\text{H}$ $\lambda_{SRG} = 1.88 \text{ fm}^{-1}$	same same same same \approx same
3N	$N^2\text{LO}$ $\Lambda_{3N} = 400 \text{ MeV}$ local regulator fit to ${}^3\text{H}$ BE, $t_{1/2}$ consistently SRG evolved	same \approx same non-local regulator fit to ${}^3\text{H}$ BE, ${}^4\text{He}$ r_{ch} no SRG for 3N



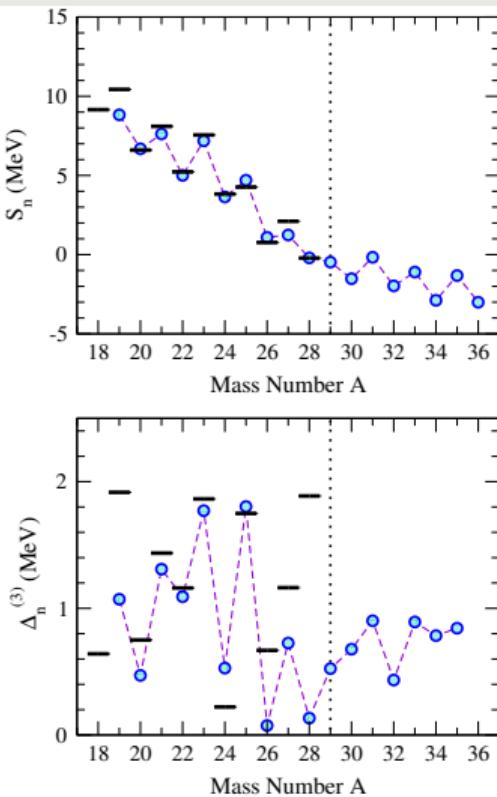
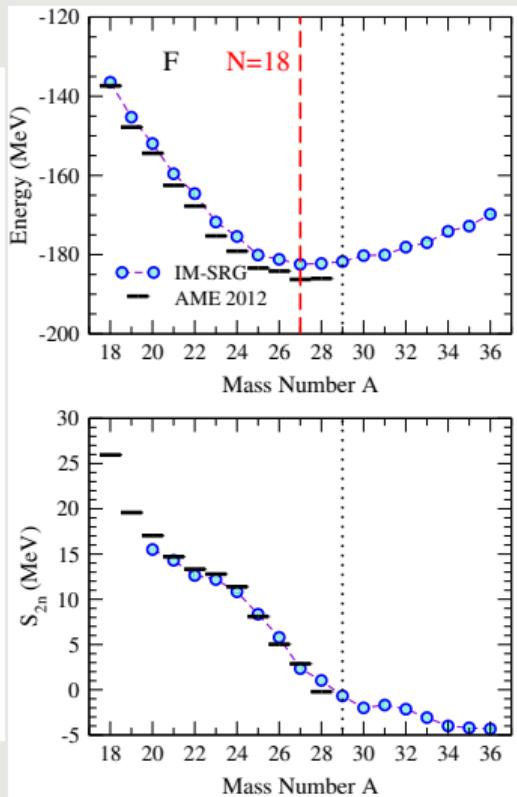
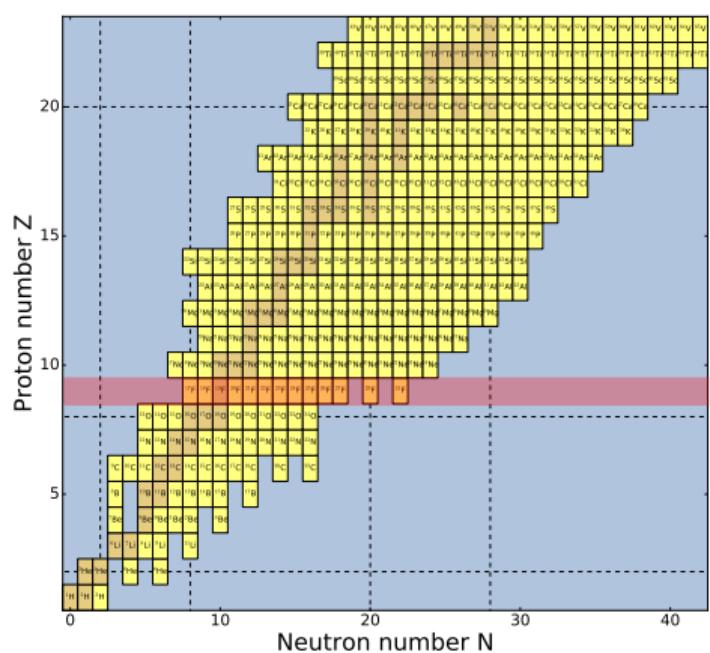
Entem and Machleidt PRC (2003), Gazit et al PRL (2009), Hebeler et al. PRC(R) (2011), Drischler et al. PRC (2016), Simonis et al. (in prep.)

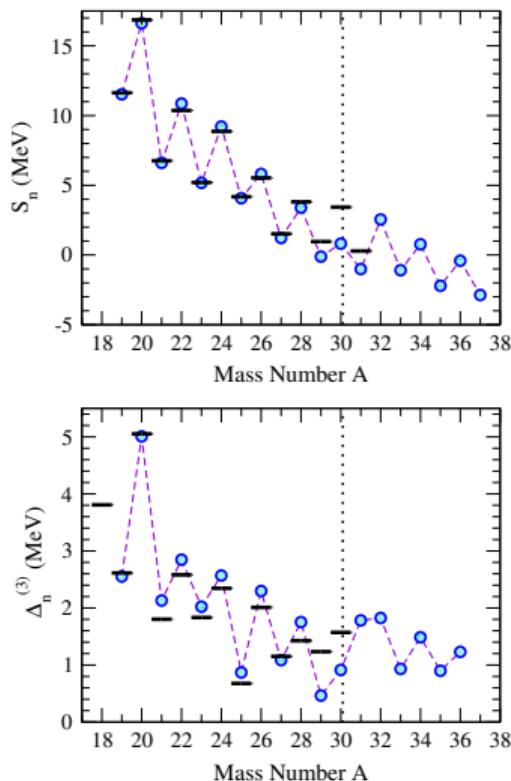
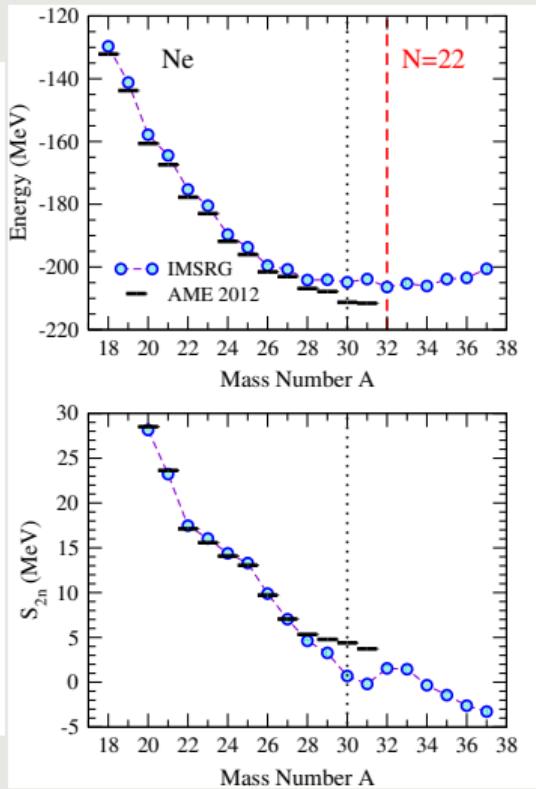
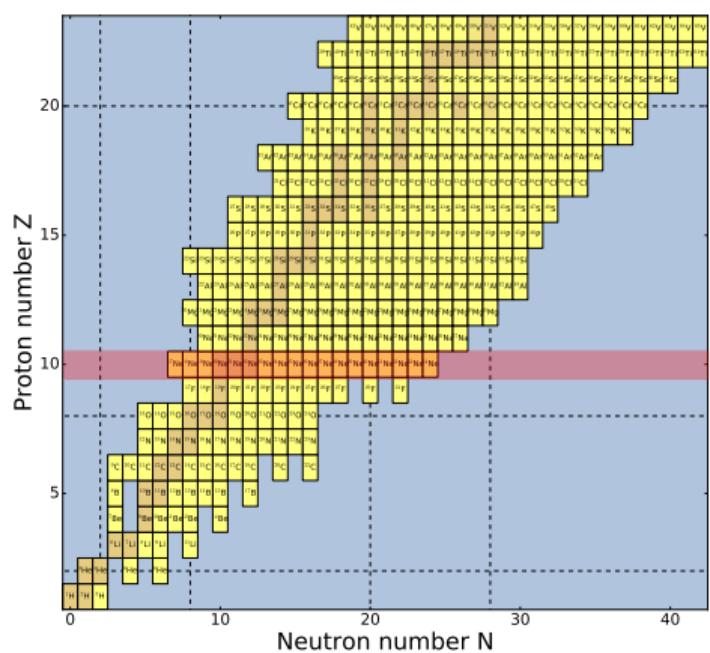
	EM 500/400	EM 1.8/2.0
NN	$N^3\text{LO}$ $\Lambda_{2N} = 500 \text{ MeV}$ non-local regulator fit to NN scattering, ${}^2\text{H}$ $\lambda_{SRG} = 1.88 \text{ fm}^{-1}$	same same same same \approx same
3N	$N^2\text{LO}$ $\Lambda_{3N} = 400 \text{ MeV}$ local regulator fit to ${}^3\text{H}$ BE, $t_{1/2}$ consistently SRG evolved	same \approx same non-local regulator fit to ${}^3\text{H}$ BE, ${}^4\text{He}$ r_{ch} no SRG for 3N

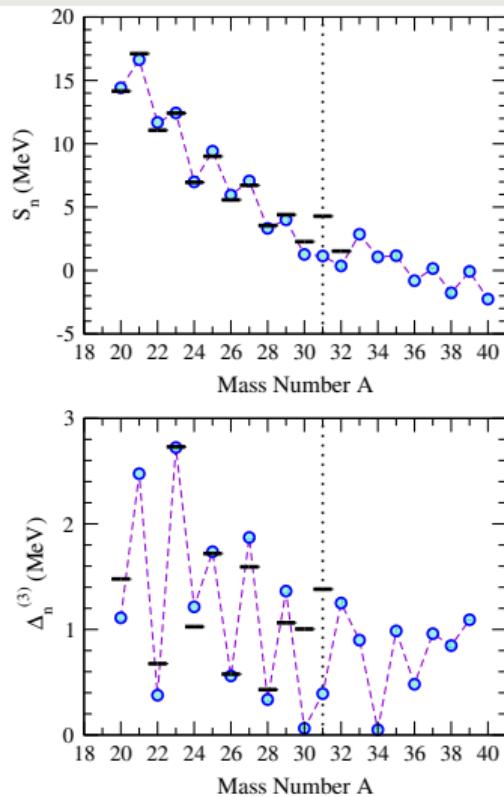
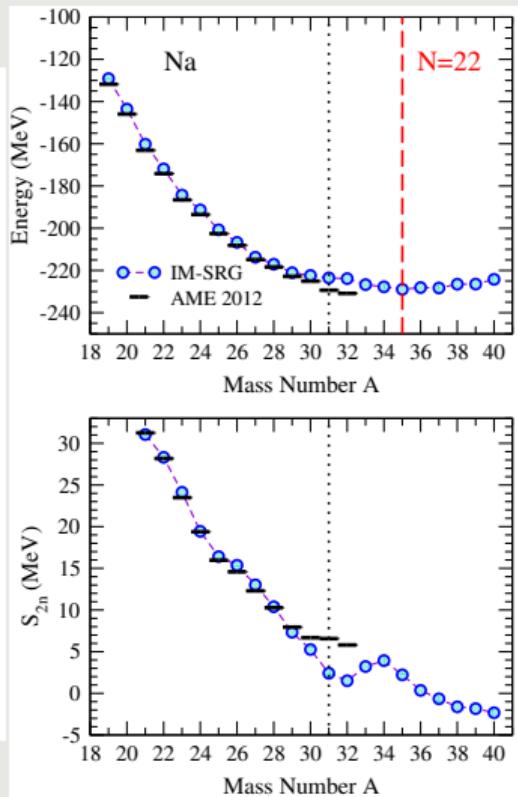
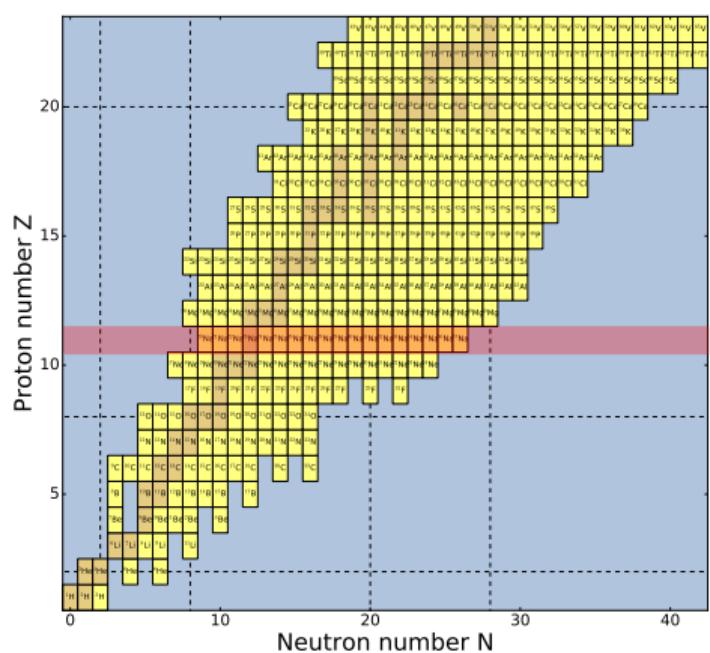


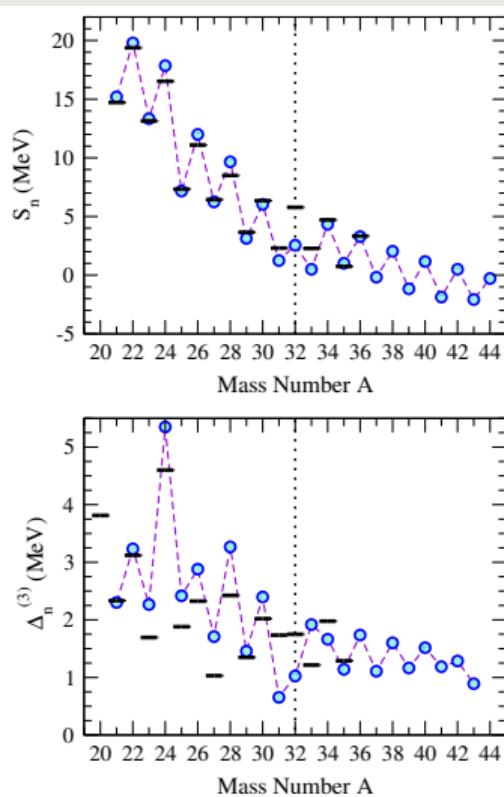
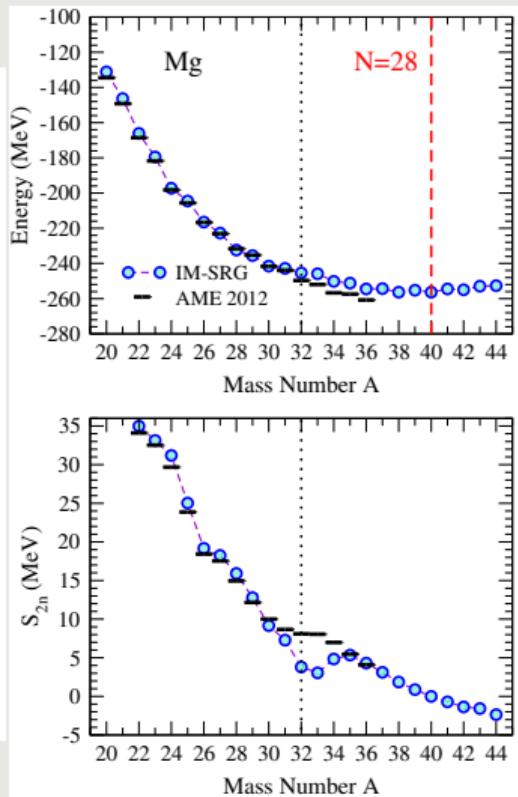
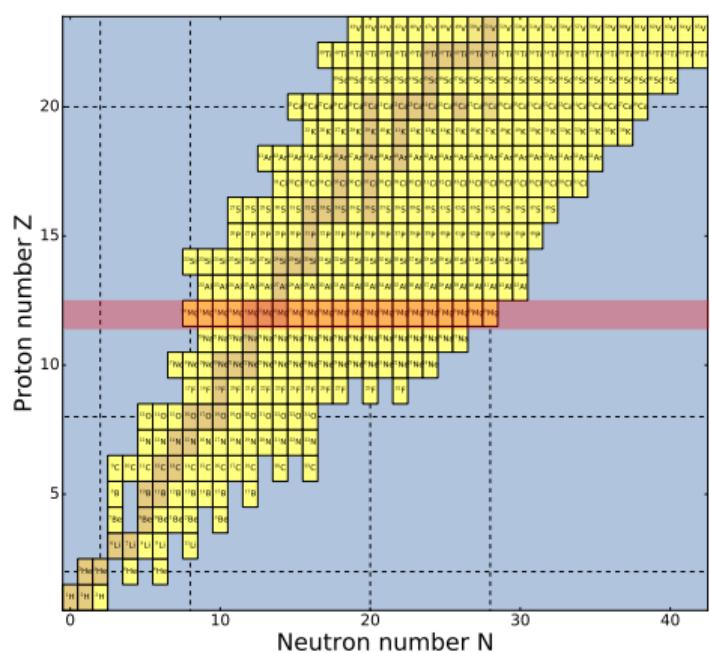
- Neither interaction is fully consistent
however...
- Saturation properties are important for finite nuclei

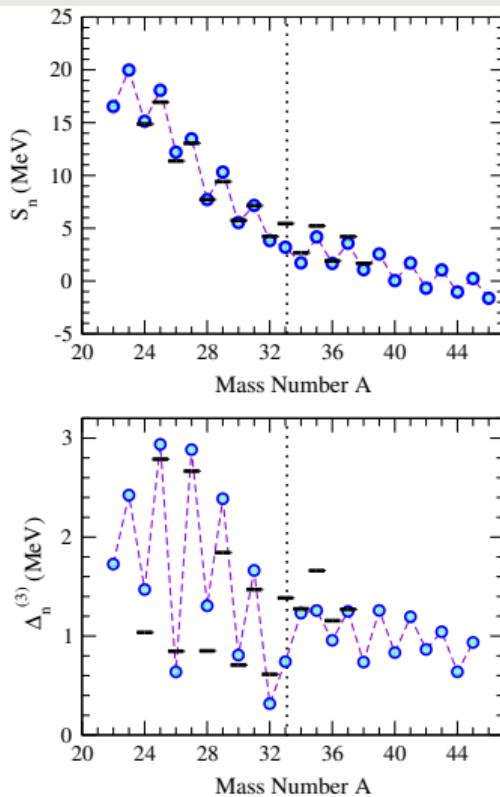
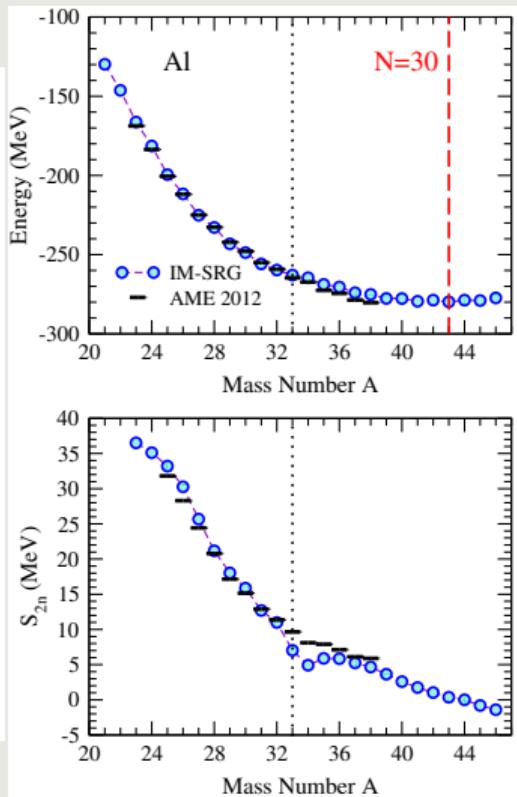
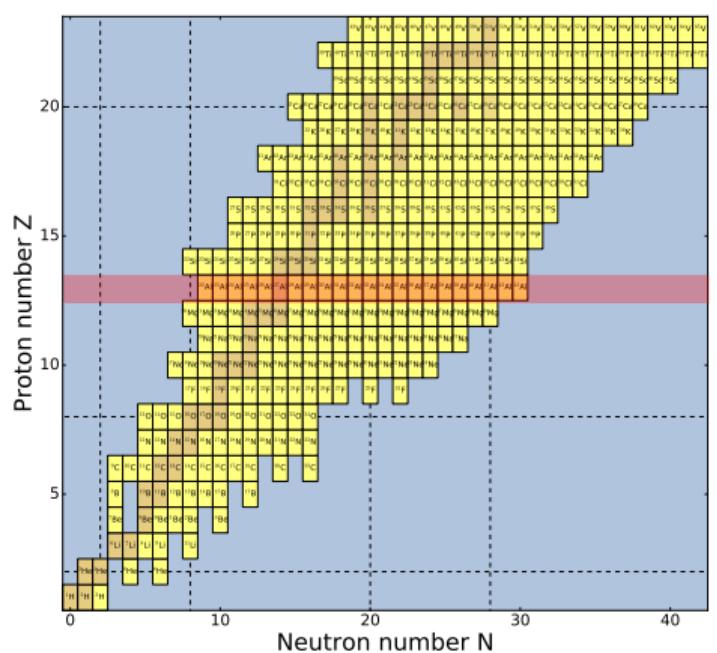
Entem and Machleidt PRC (2003), Gazit et al PRL (2009), Hebeler et al. PRC(R) (2011), Drischler et al. PRC (2016), Simonis et al. (in prep.)

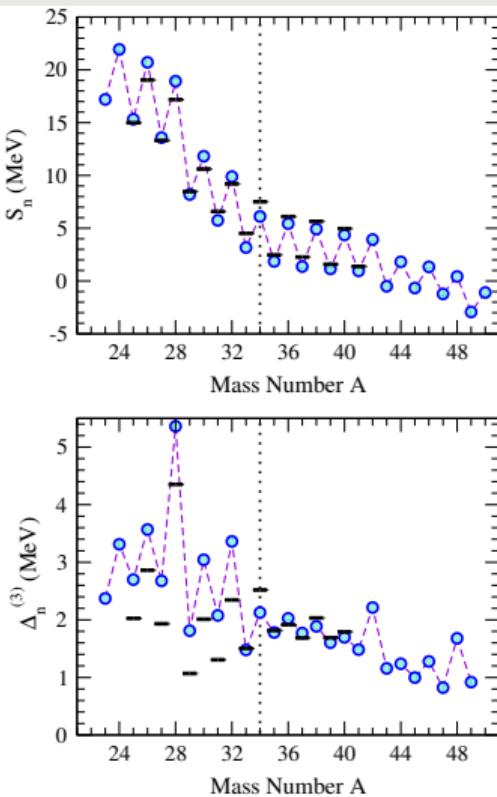
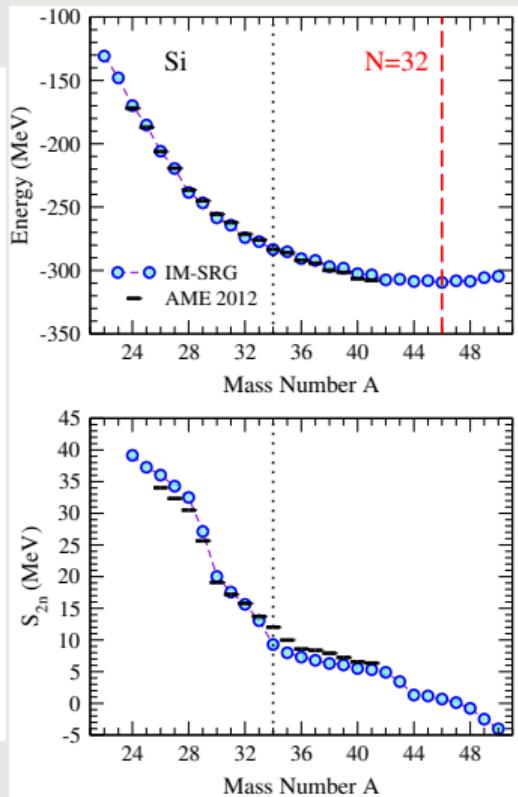
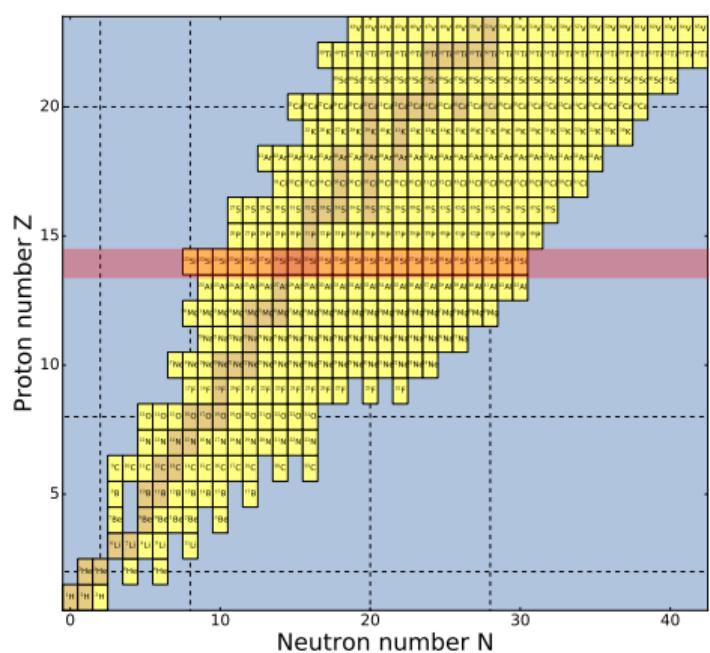


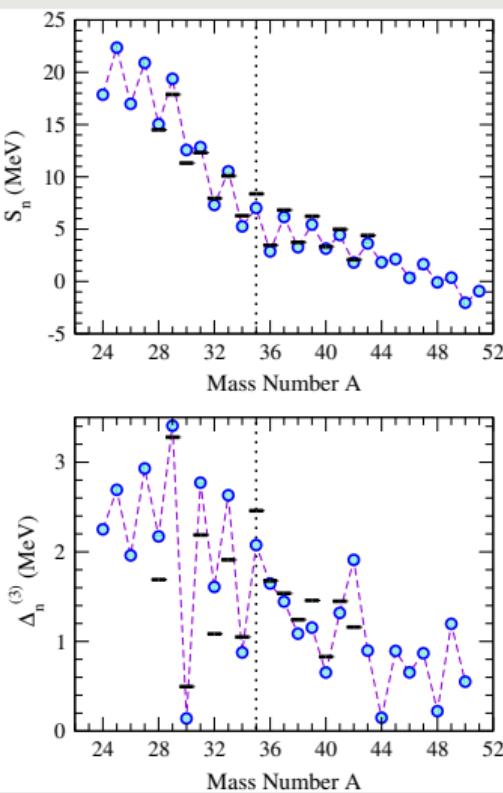
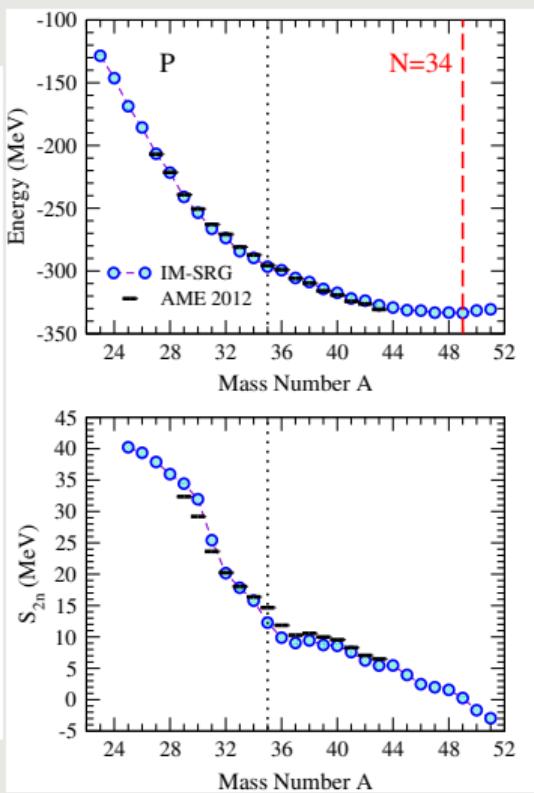
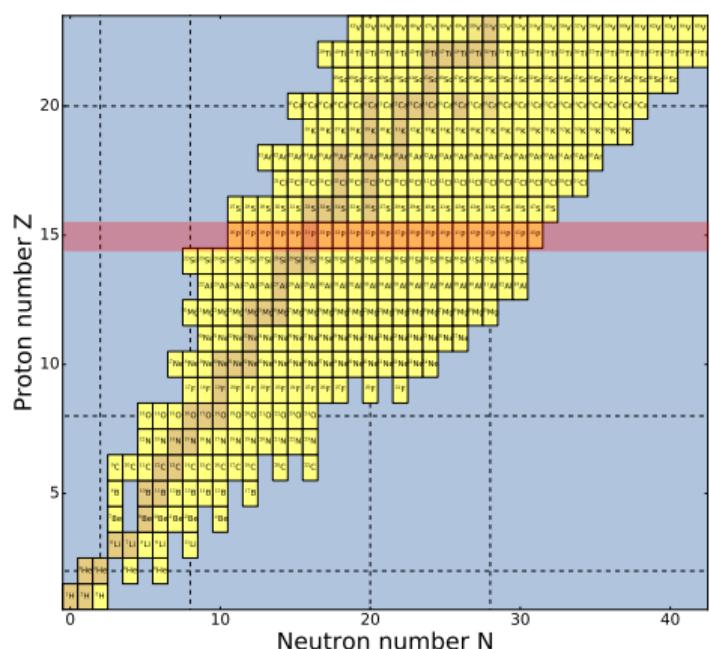


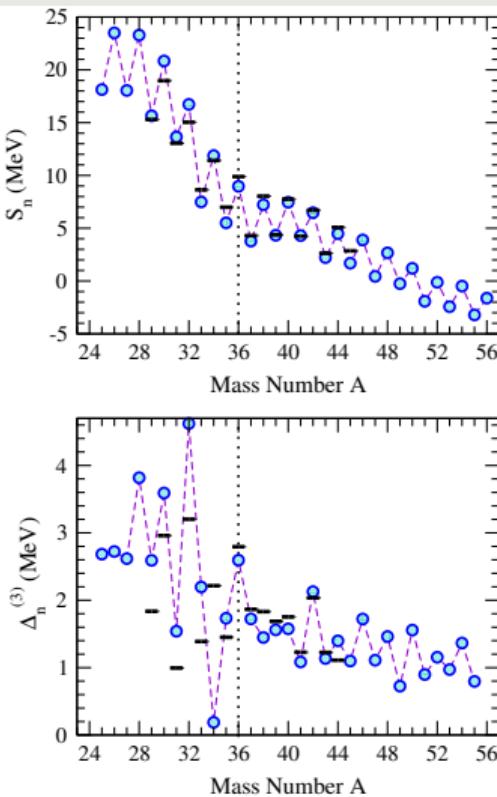
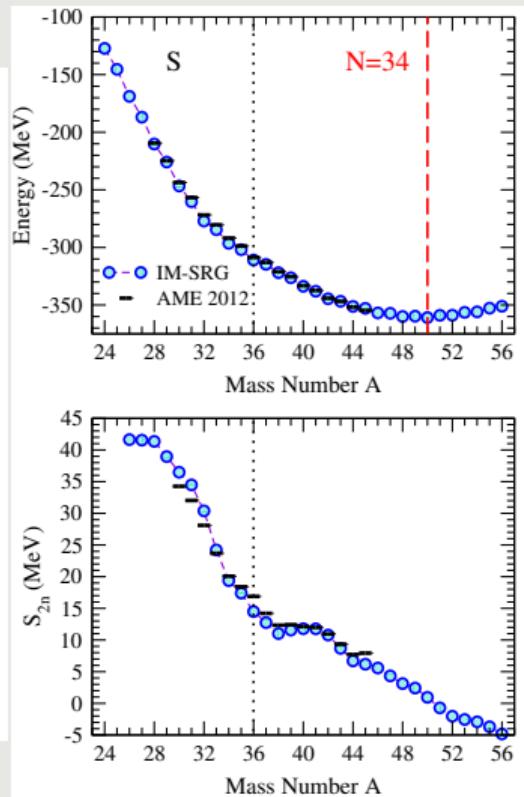
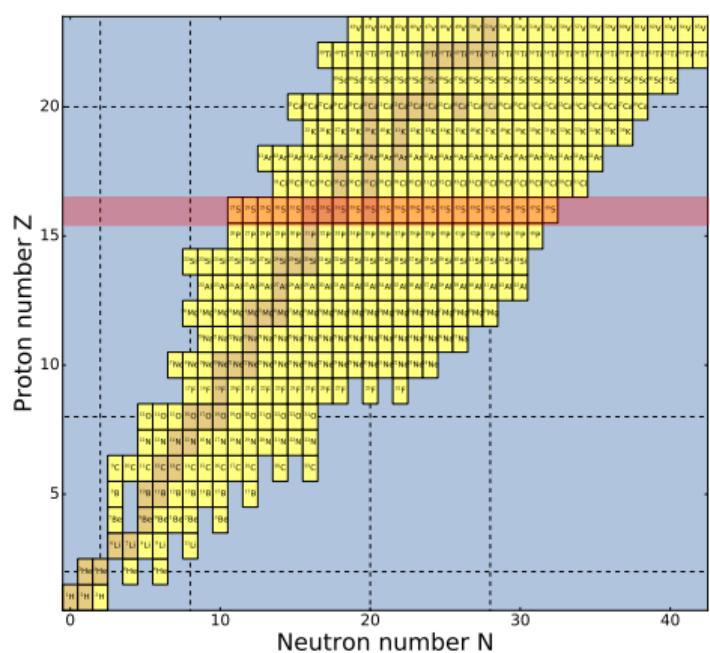


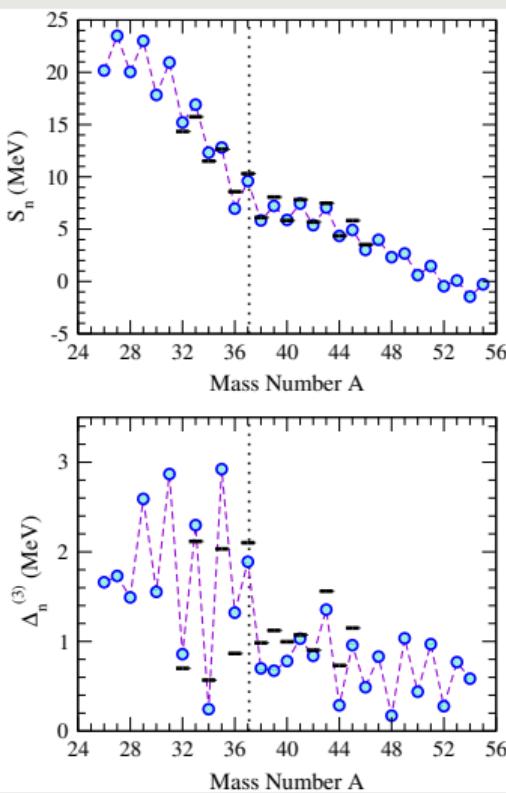
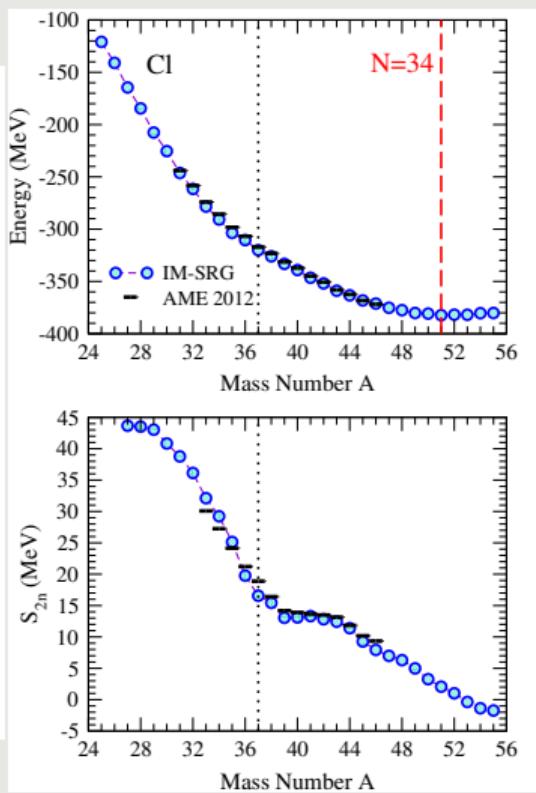
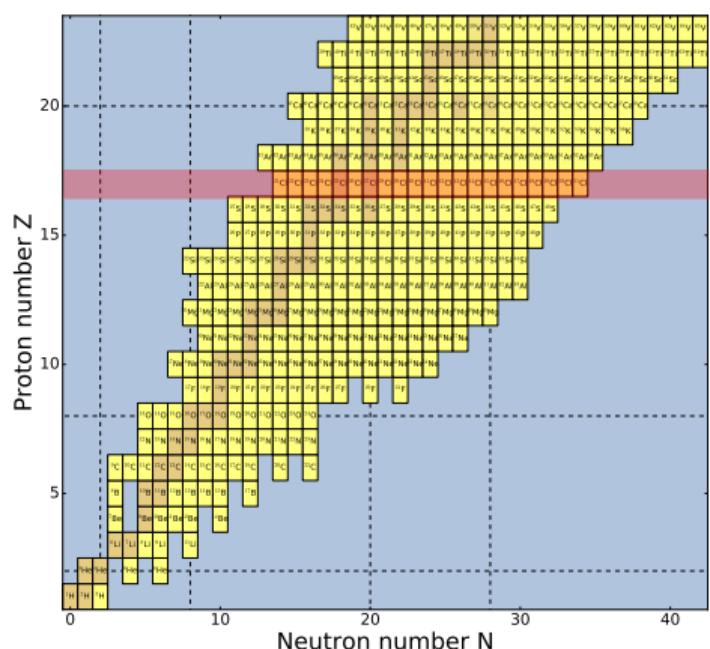


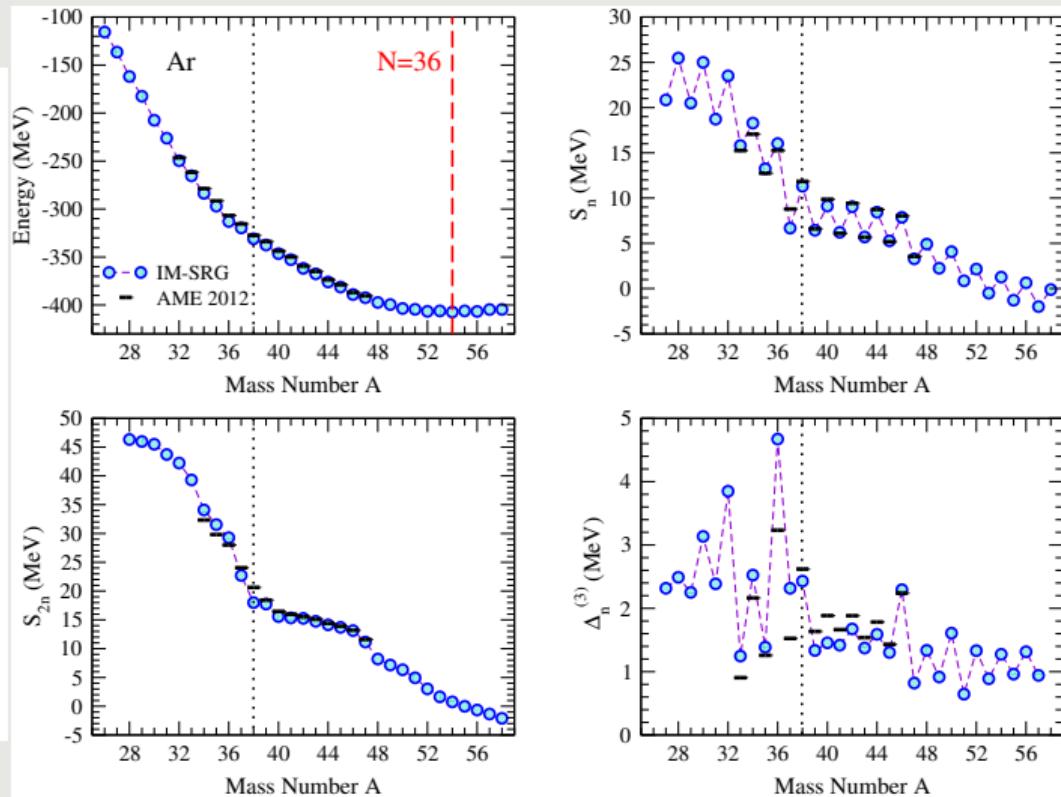
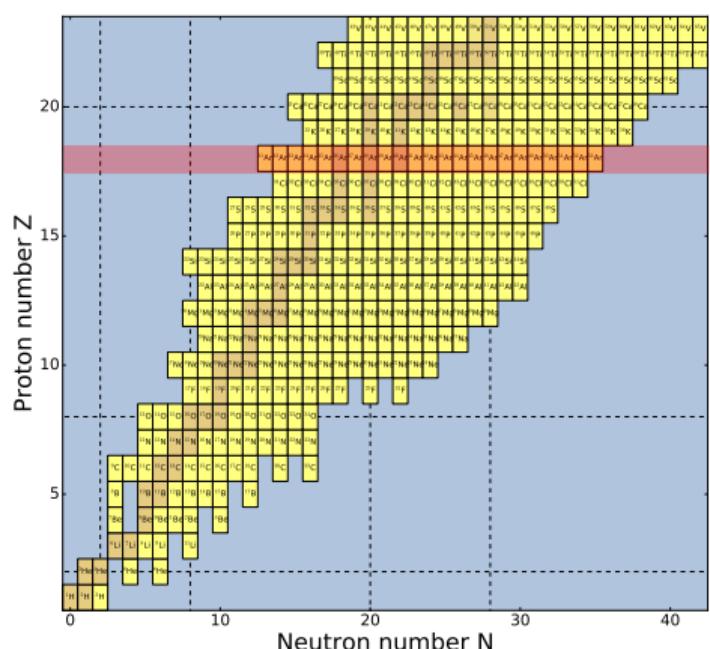


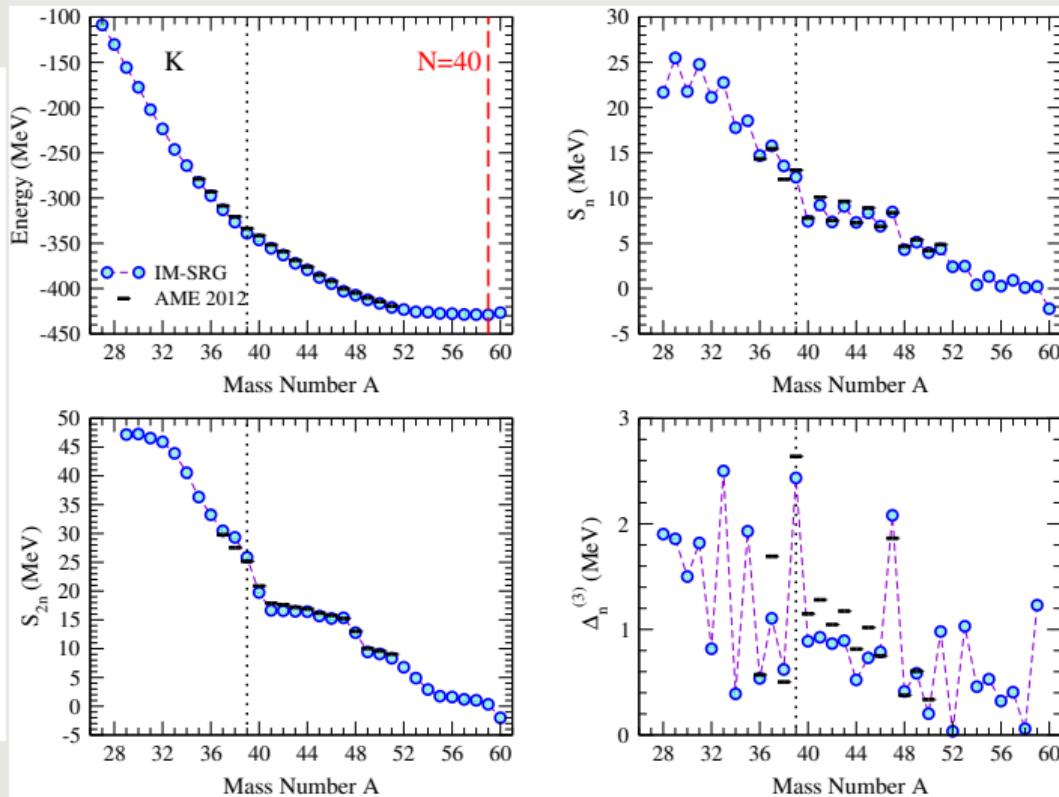
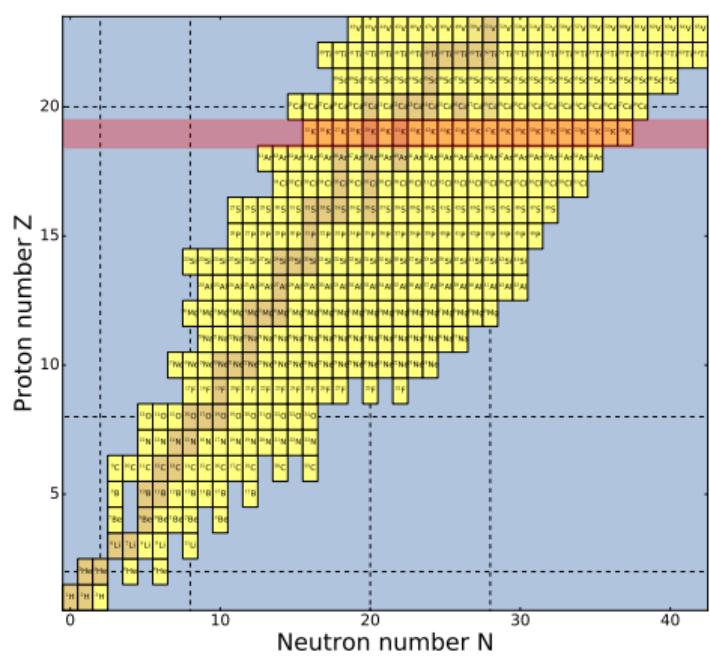


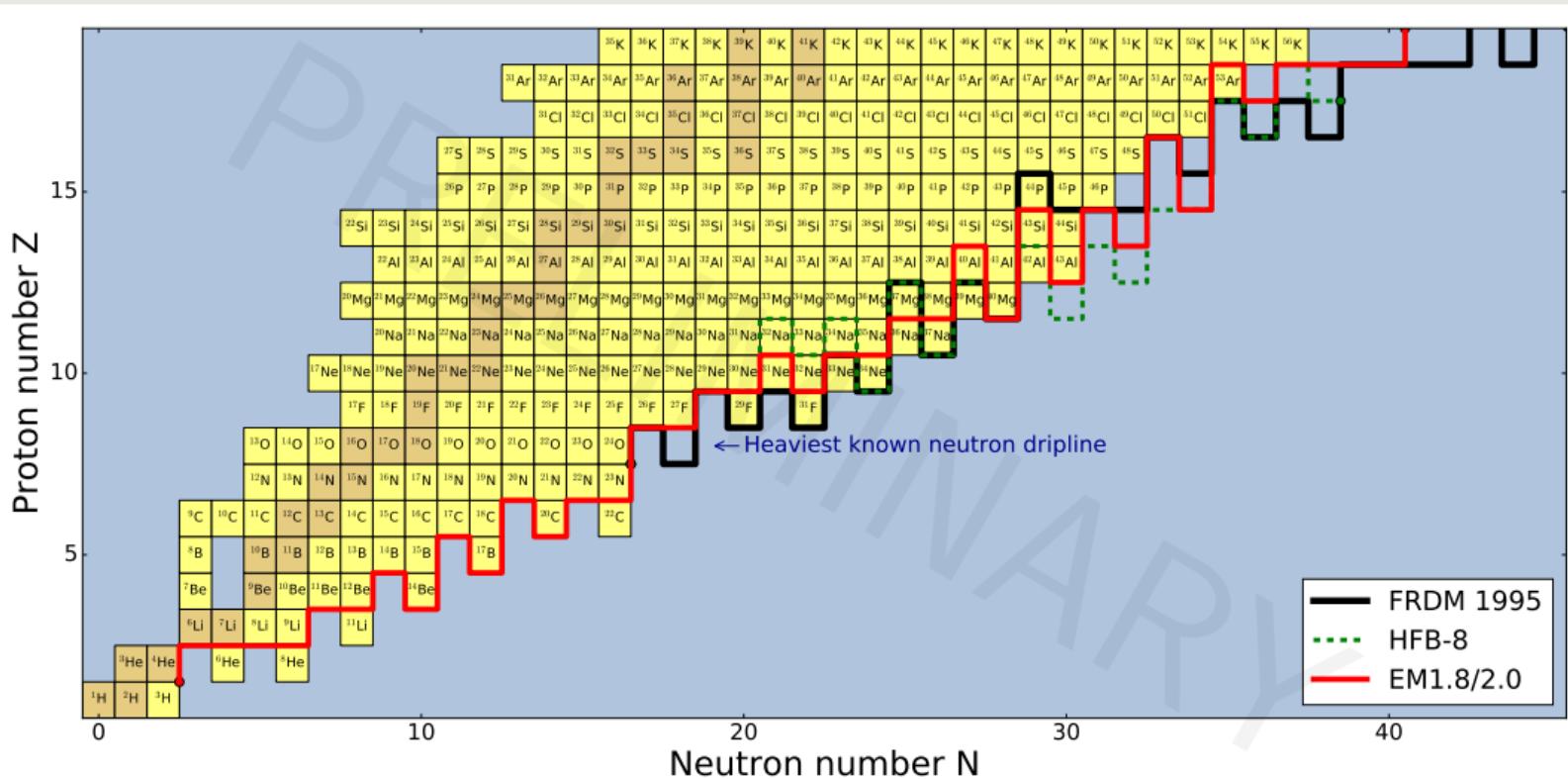




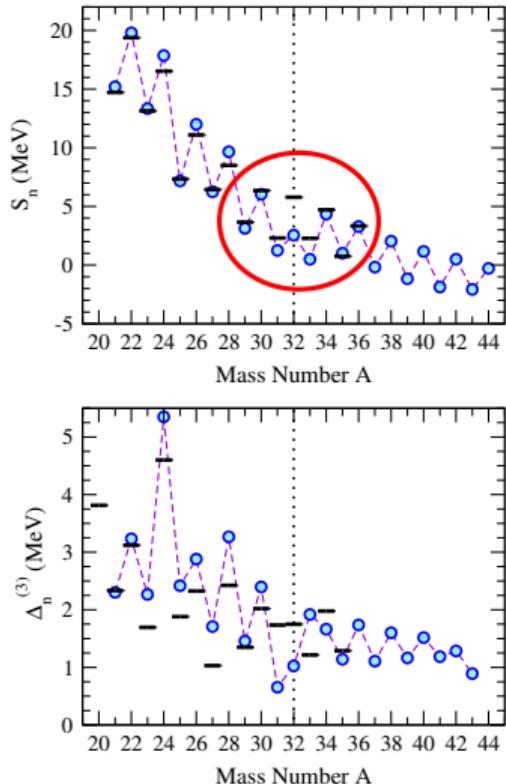
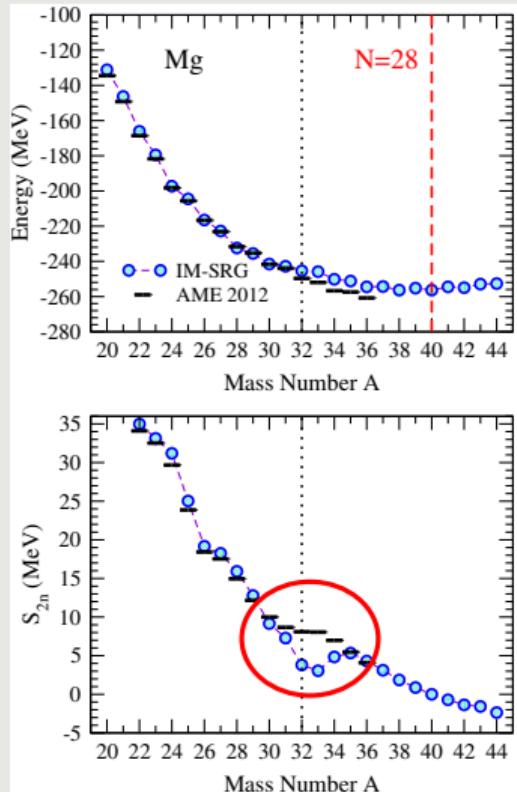
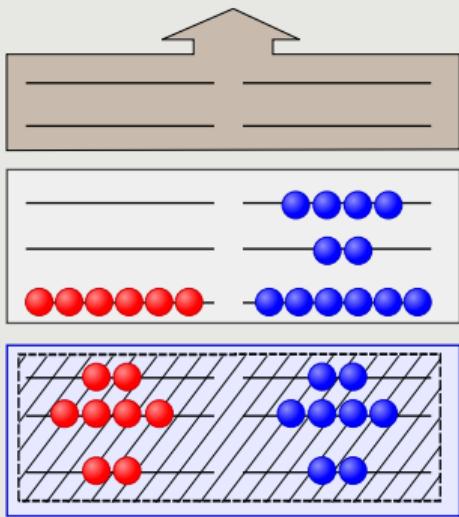






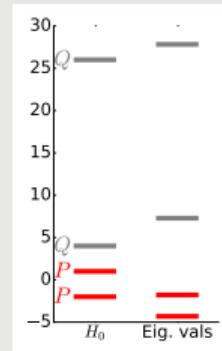


Baumann et al. Nature (2007), Möller et al. (1995), Samyn et al. (2004), Holt et al. (in prep.)



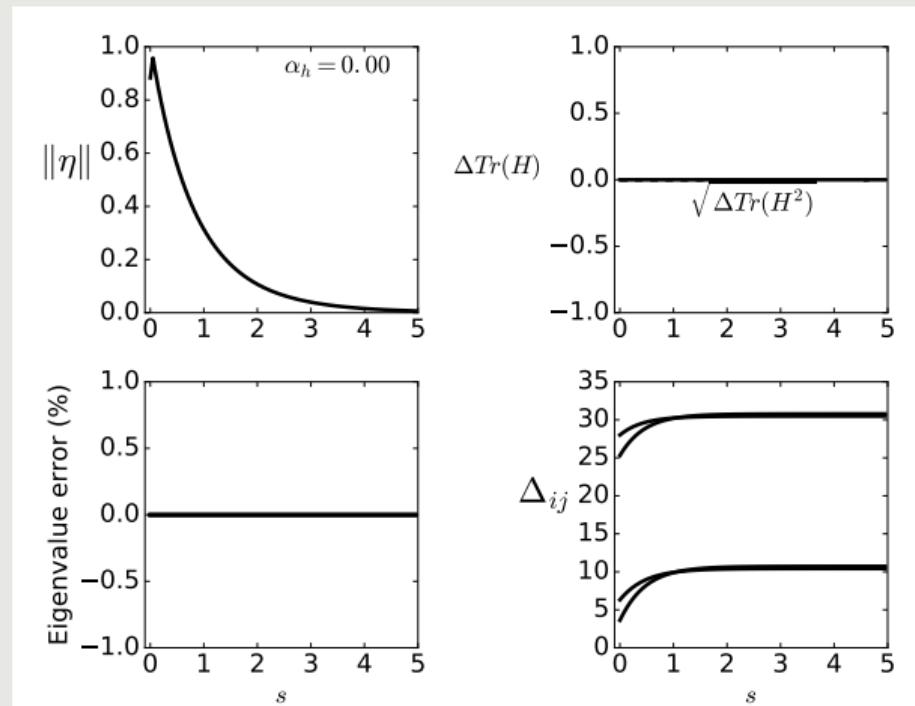
First, a toy problem:

$$H = \begin{pmatrix} 1 & 5 & 0 & 5 \\ 5 & 26 & 5 & 0 \\ 0 & 5 & -2 & 1 \\ 5 & 0 & 1 & 4 \end{pmatrix}$$



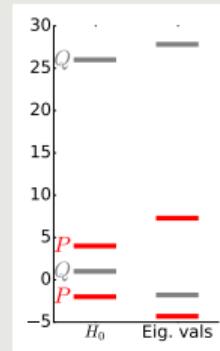
$$\eta = \frac{1}{2} \text{atan} \left(\frac{2H_{qp}}{H_{qq} - H_{pp}} \right) - h.c.$$

$$H(s + \delta s) = e^{\eta_s \delta s} H(s) e^{\eta_s^\dagger \delta s}$$



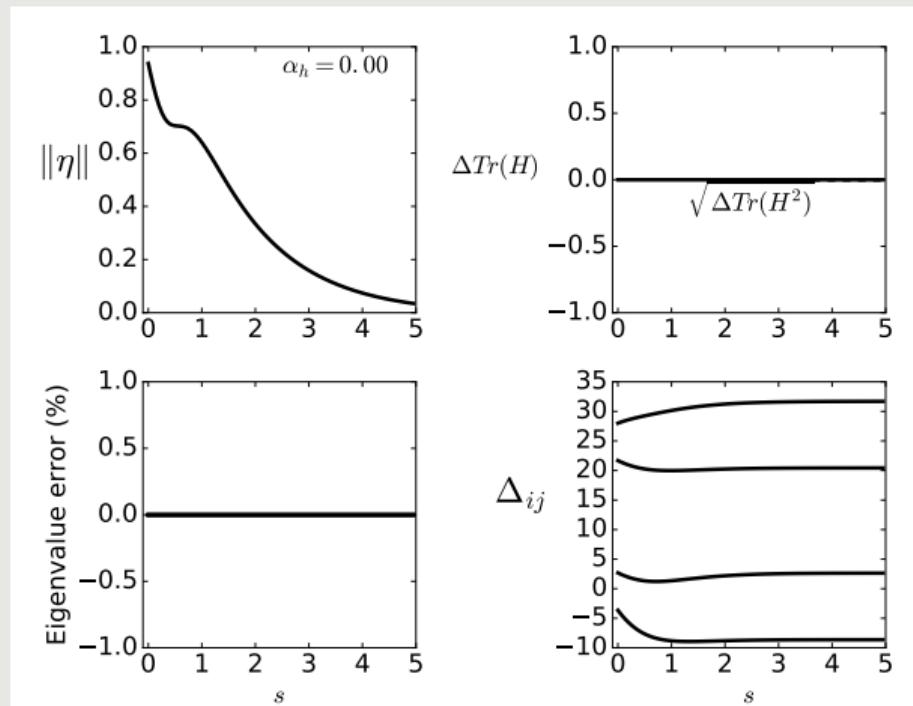
First, a toy problem:

$$H = \begin{pmatrix} 1 & 5 & 0 & 5 \\ 5 & 26 & 5 & 0 \\ 0 & 5 & -2 & 1 \\ 5 & 0 & 1 & 4 \end{pmatrix}$$



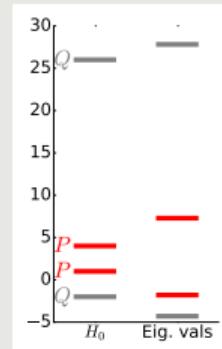
$$\eta = \frac{1}{2} \text{atan} \left(\frac{2H_{qp}}{H_{qq} - H_{pp}} \right) - h.c.$$

$$H(s + \delta s) = e^{\eta_s \delta s} H(s) e^{\eta_s^\dagger \delta s}$$



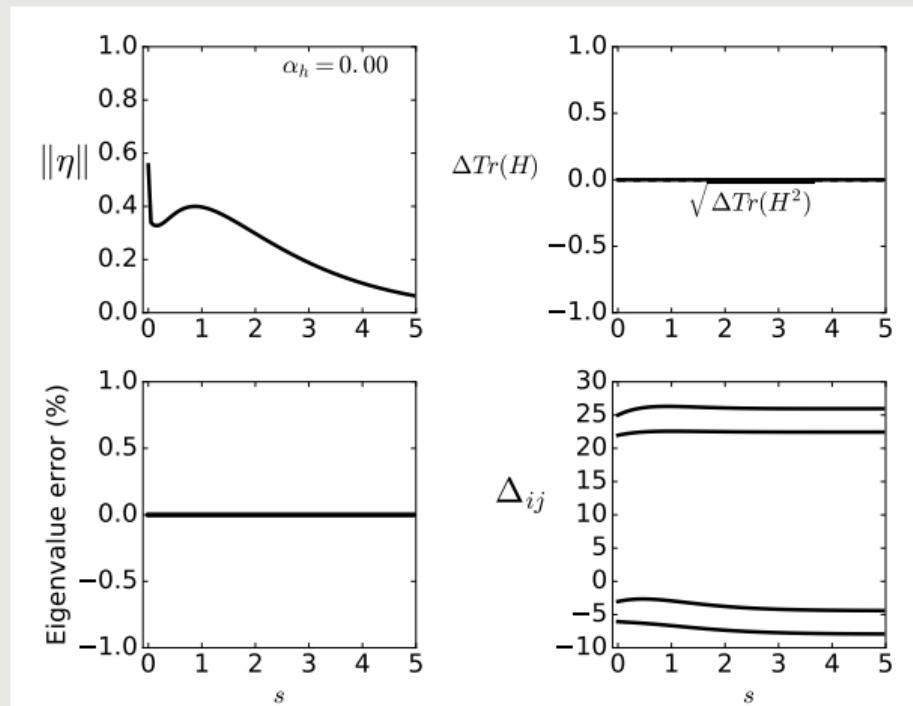
First, a toy problem:

$$H = \begin{pmatrix} 1 & 5 & 0 & 5 \\ 5 & 26 & 5 & 0 \\ 0 & 5 & -2 & 1 \\ 5 & 0 & 1 & 4 \end{pmatrix}$$



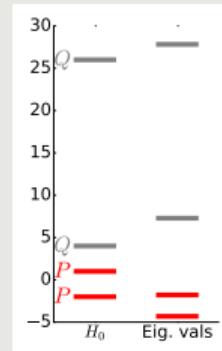
$$\eta = \frac{1}{2} \text{atan} \left(\frac{2H_{qp}}{H_{qq} - H_{pp}} \right) - h.c.$$

$$H(s + \delta s) = e^{\eta_s \delta s} H(s) e^{\eta_s^\dagger \delta s}$$



First, a toy problem:

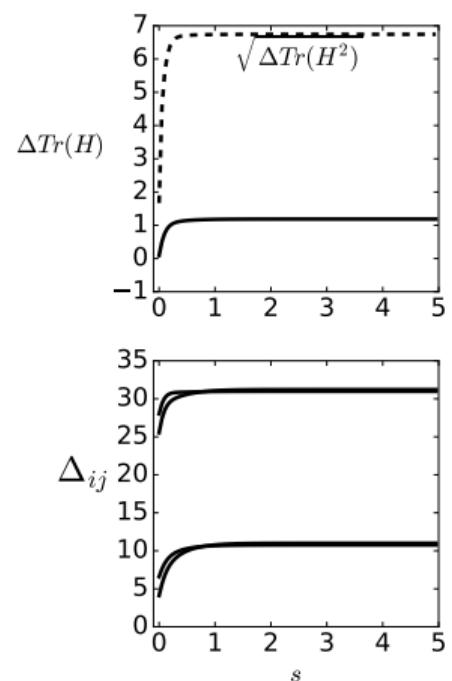
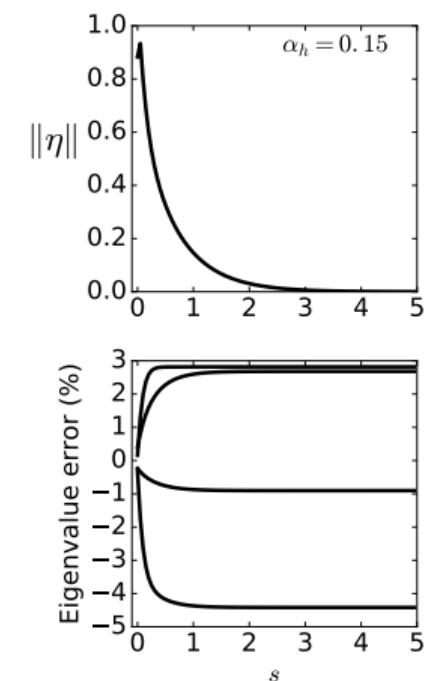
$$H = \begin{pmatrix} 1 & 5 & 0 & 5 \\ 5 & 26 & 5 & 0 \\ 0 & 5 & -2 & 1 \\ 5 & 0 & 1 & 4 \end{pmatrix}$$



$$\eta = \frac{1}{2} \text{atan} \left(\frac{2H_{qp}}{H_{qq} - H_{pp}} \right) - h.c.$$

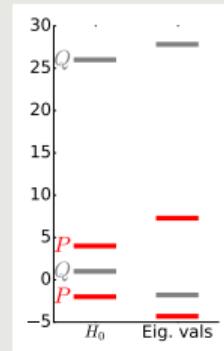
$$H(s + \delta s) = e^{\eta_s \delta s} H(s) e^{\eta_s^\dagger \delta s}$$

break unitarity: $\eta \rightarrow \eta + \alpha_h [\eta, H]$



First, a toy problem:

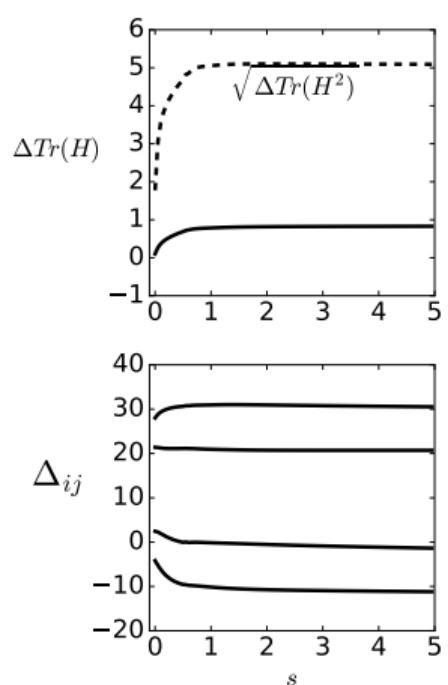
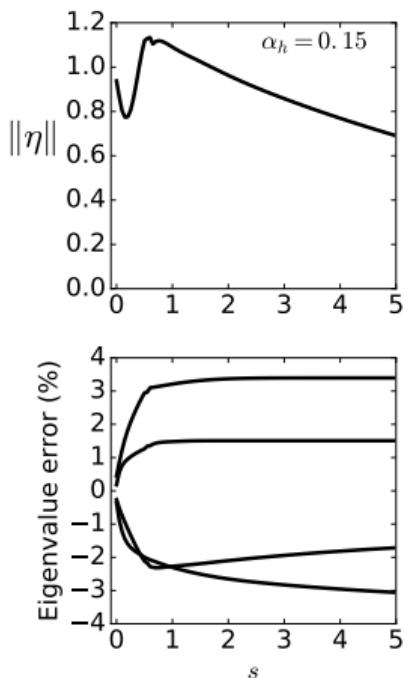
$$H = \begin{pmatrix} 1 & 5 & 0 & 5 \\ 5 & 26 & 5 & 0 \\ 0 & 5 & -2 & 1 \\ 5 & 0 & 1 & 4 \end{pmatrix}$$



$$\eta = \frac{1}{2} \text{atan} \left(\frac{2H_{qp}}{H_{qq} - H_{pp}} \right) - h.c.$$

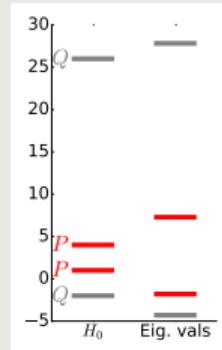
$$H(s + \delta s) = e^{\eta_s \delta s} H(s) e^{\eta_s^\dagger \delta s}$$

break unitarity: $\eta \rightarrow \eta + \alpha_h [\eta, H]$



First, a toy problem:

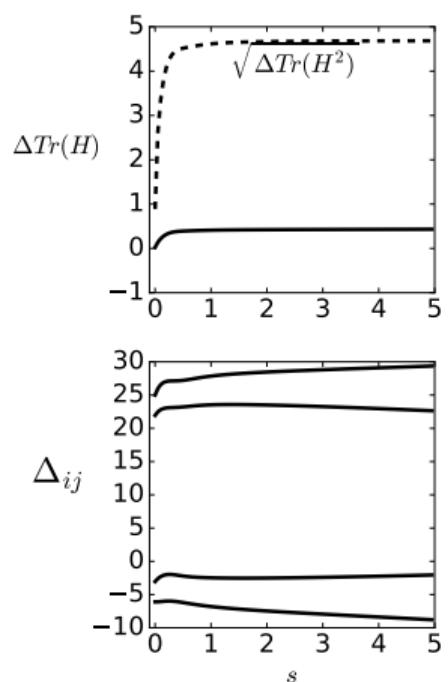
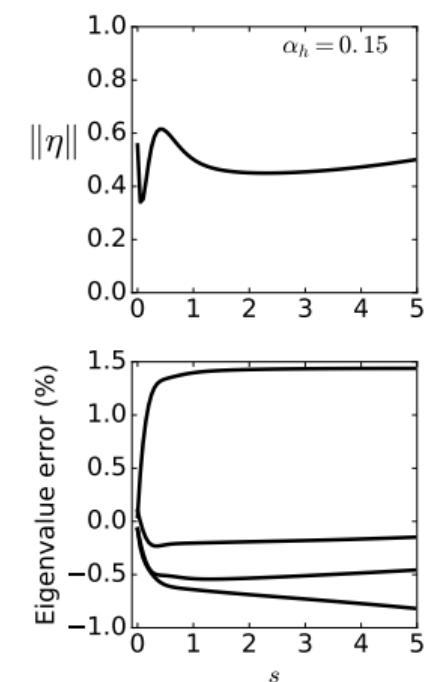
$$H = \begin{pmatrix} 1 & 5 & 0 & 5 \\ 5 & 26 & 5 & 0 \\ 0 & 5 & -2 & 1 \\ 5 & 0 & 1 & 4 \end{pmatrix}$$

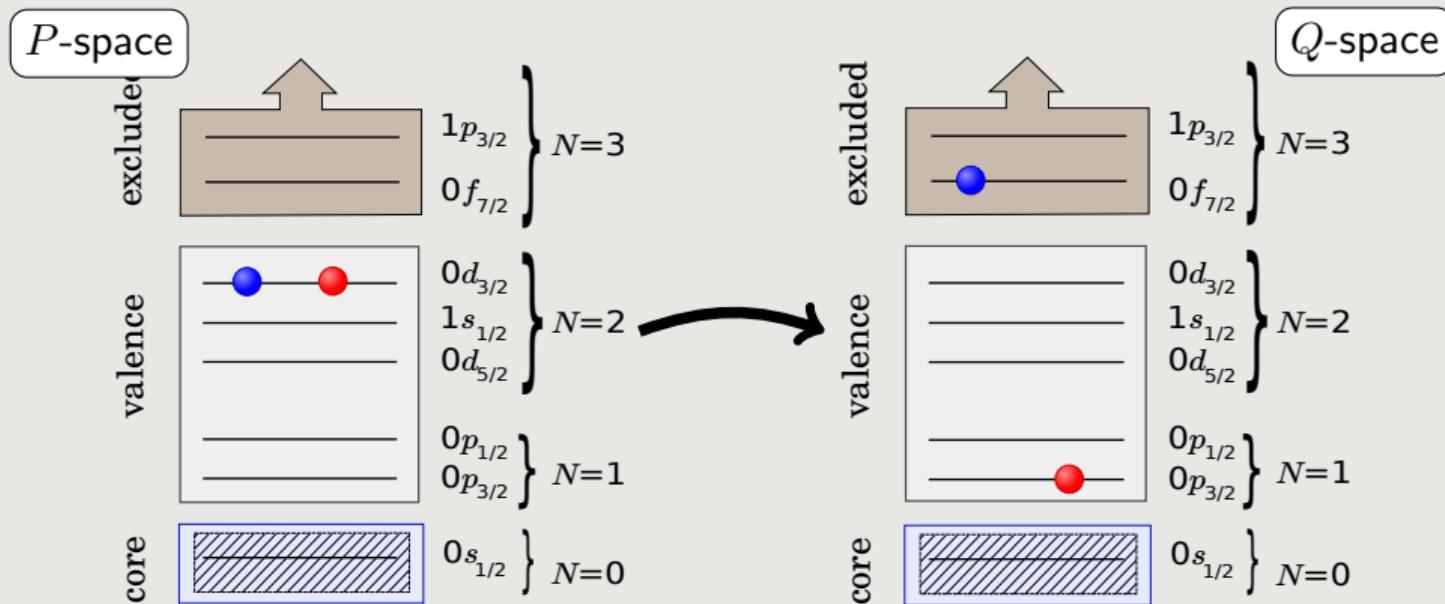


$$\eta = \frac{1}{2} \text{atan} \left(\frac{2H_{qp}}{H_{qq} - H_{pp}} \right) - h.c.$$

$$H(s + \delta s) = e^{\eta_s \delta s} H(s) e^{\eta_s^\dagger \delta s}$$

break unitarity: $\eta \rightarrow \eta + \alpha_h [\eta, H]$

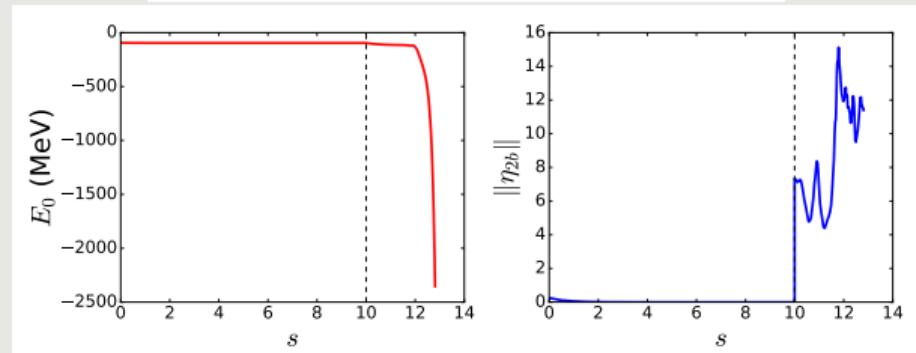
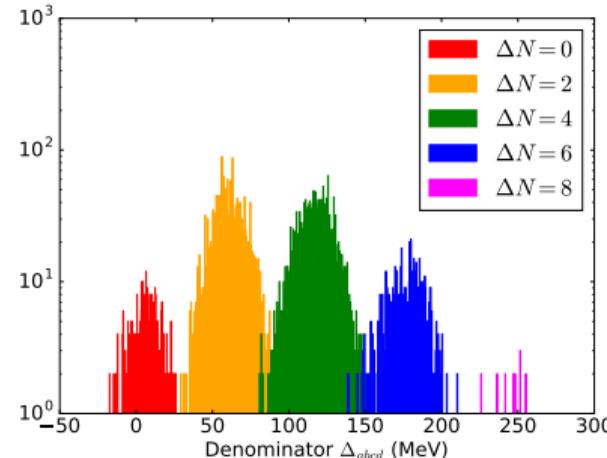
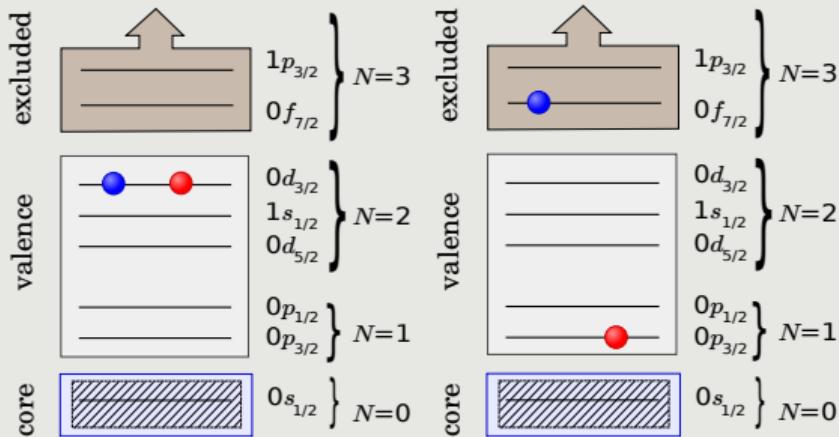




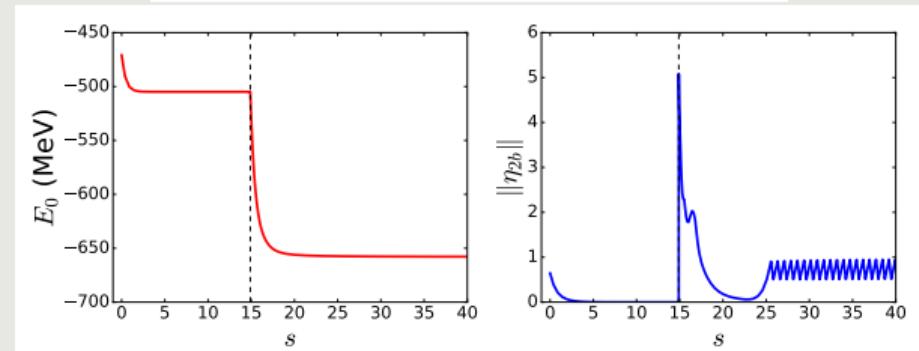
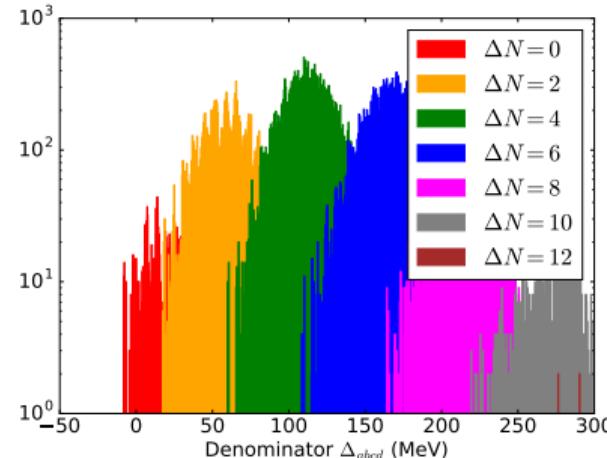
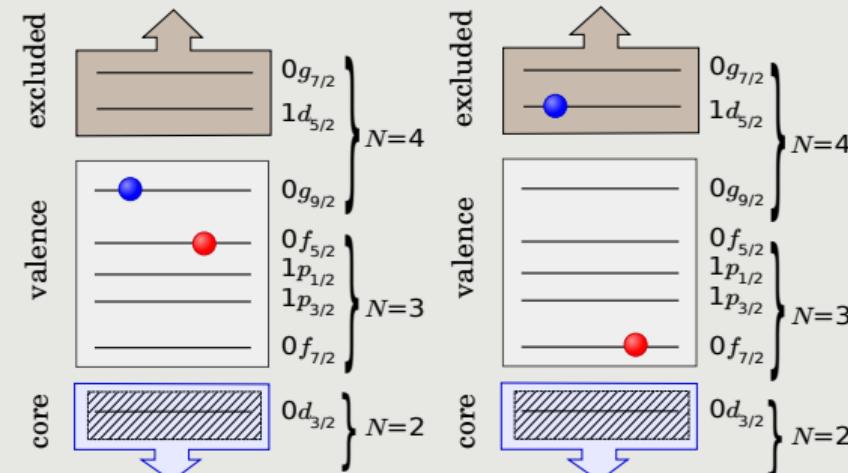
$$\eta_s = \frac{1}{2} \text{atan} \left(\frac{2H_{qp}(s)}{H_{qq}(s) - H_{pp}(s)} \right) - h.c.$$

$\Delta N = 0 \rightarrow \text{negative denominators}$

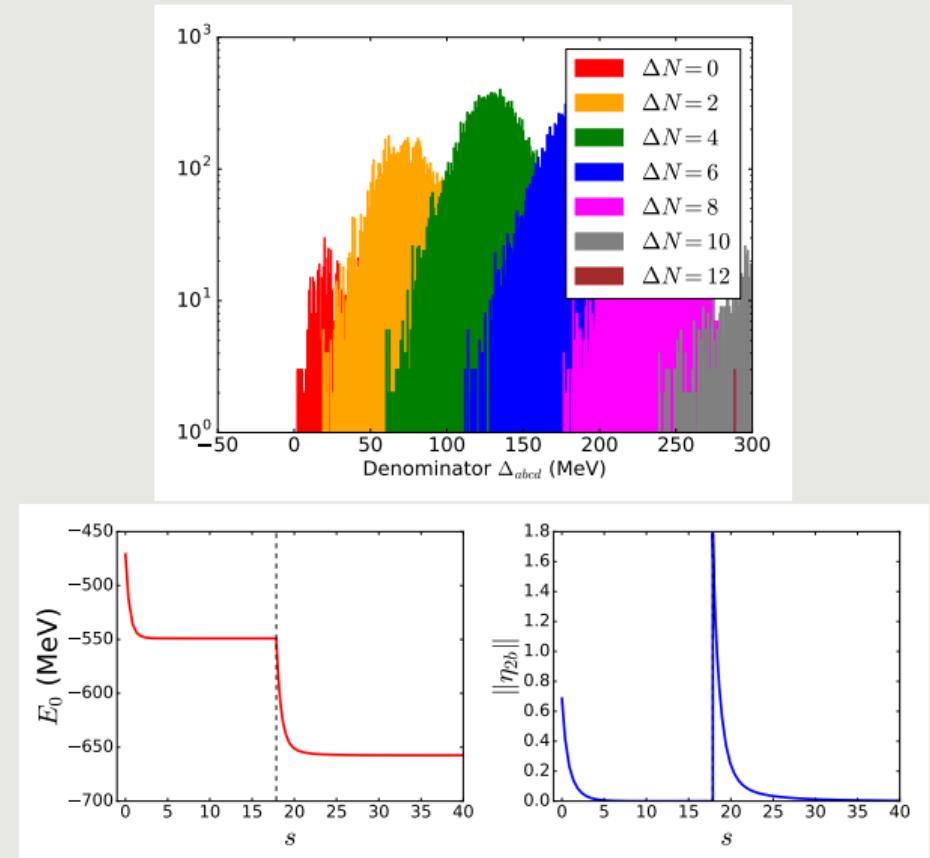
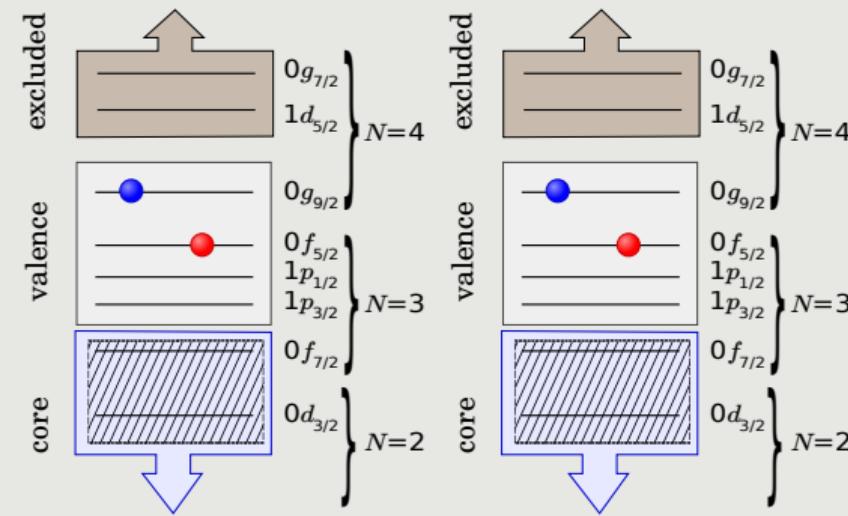
psd space, ^{16}O reference

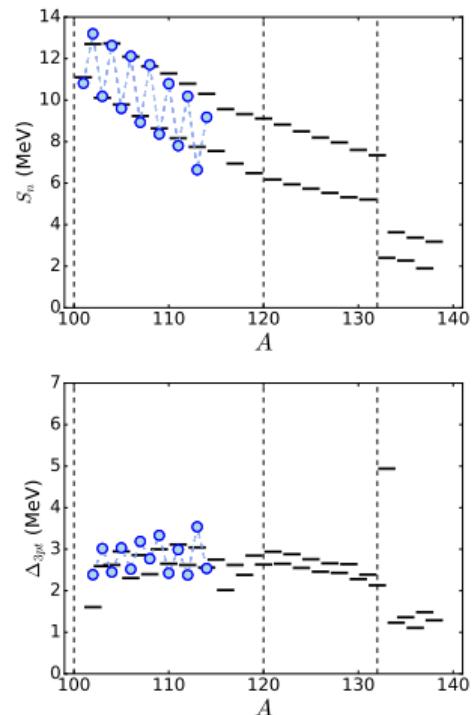
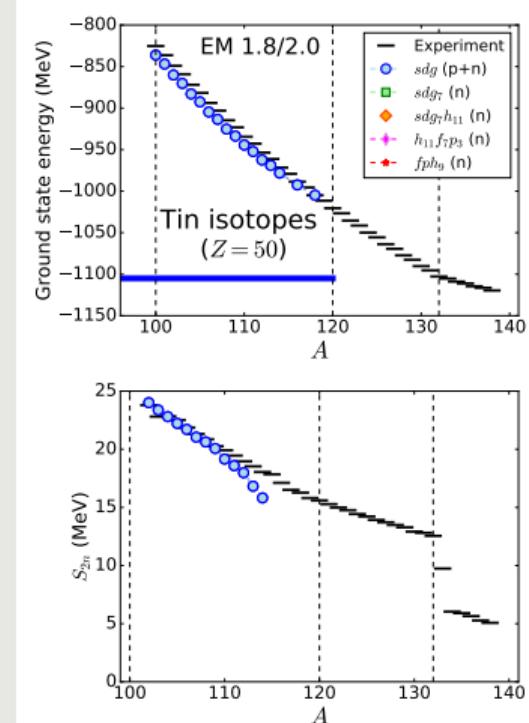
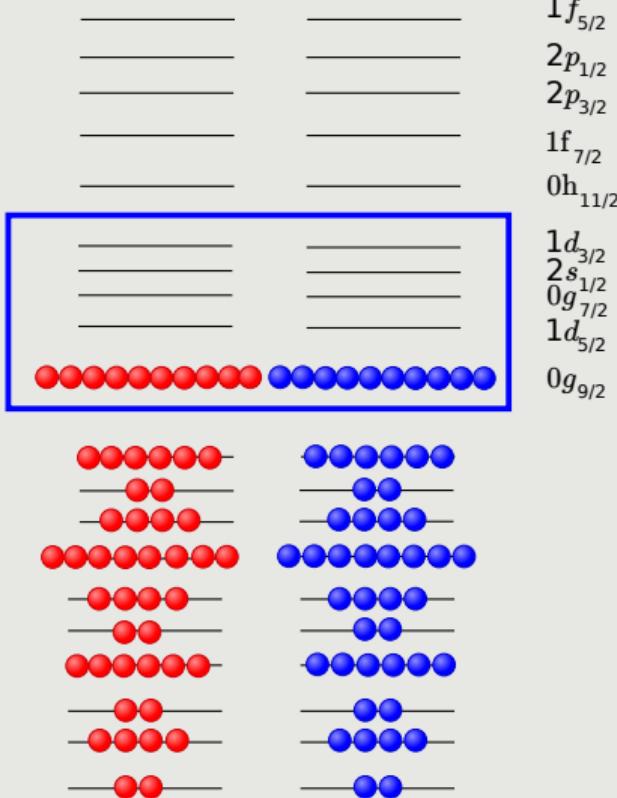


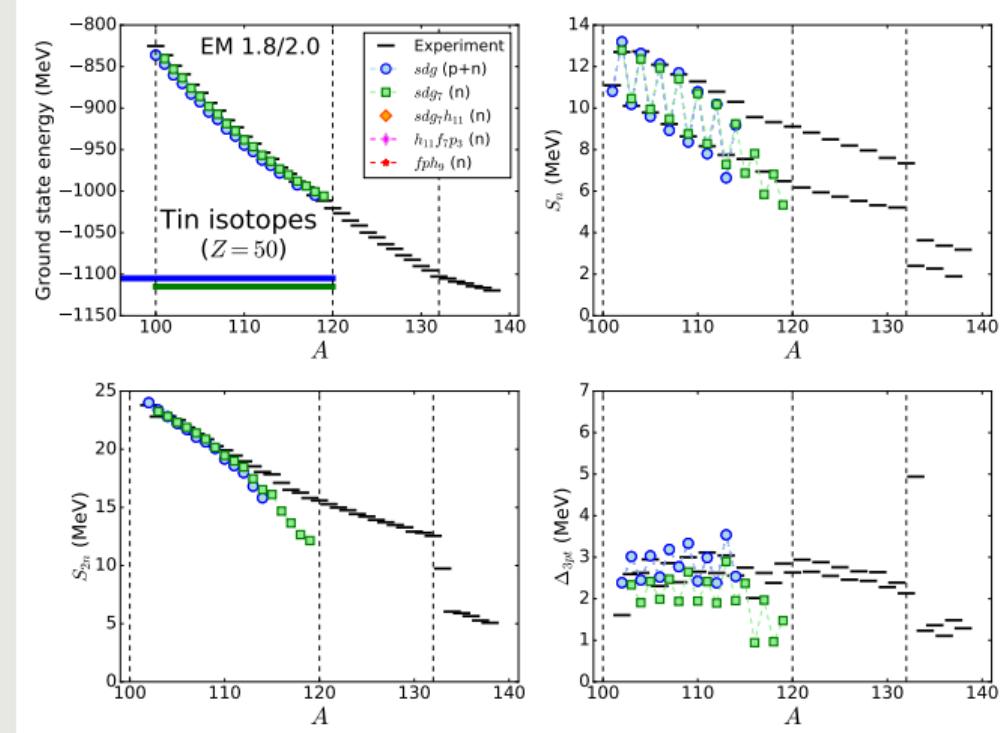
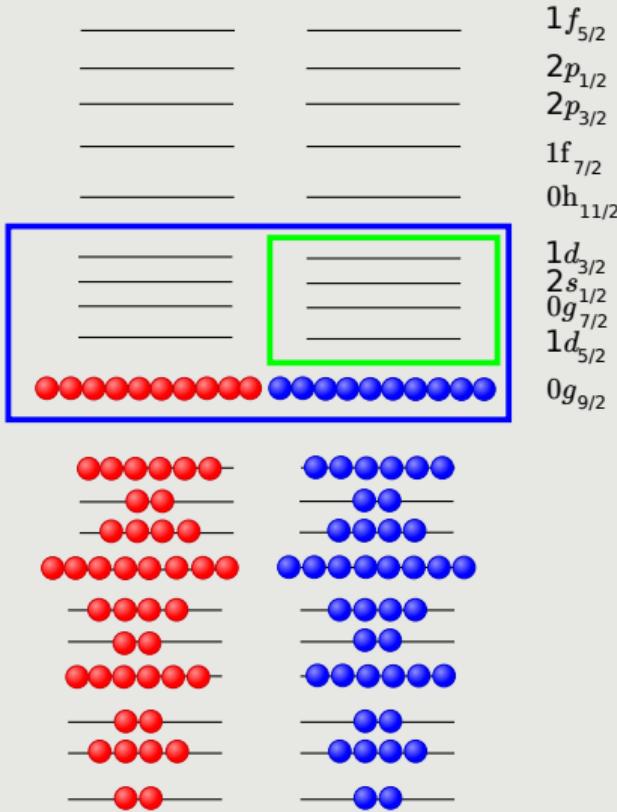
pfg_9 space, ^{76}Ge reference

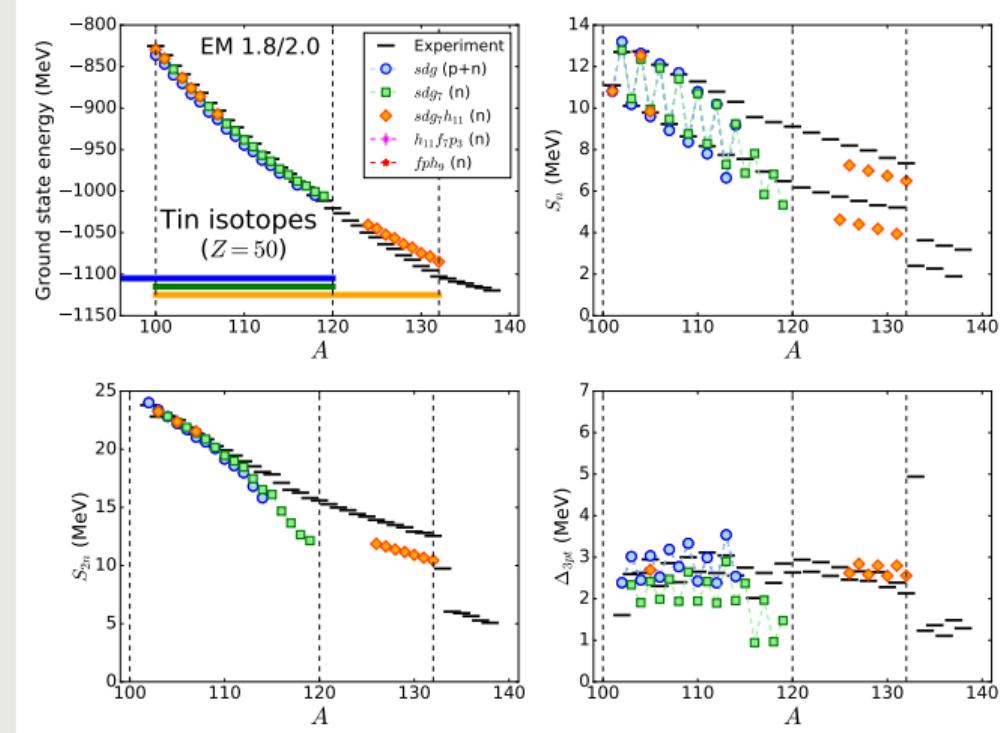
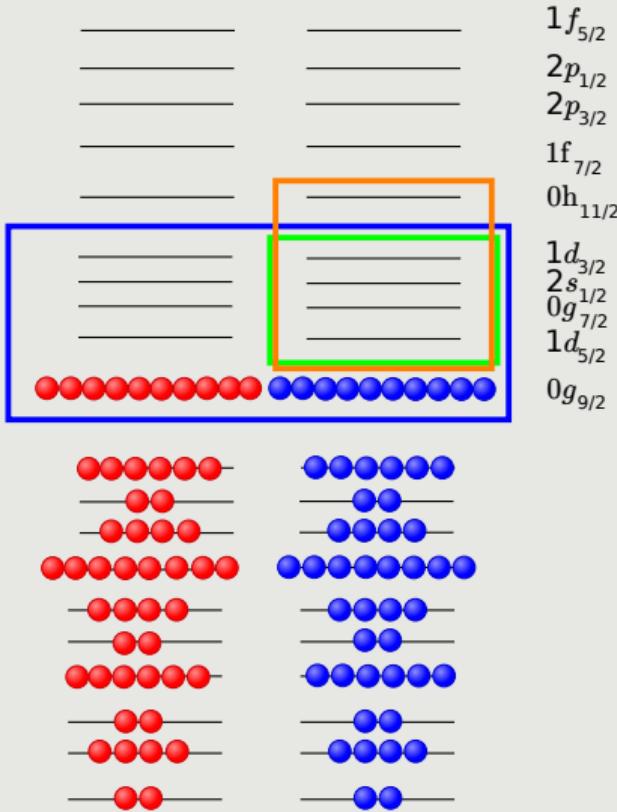


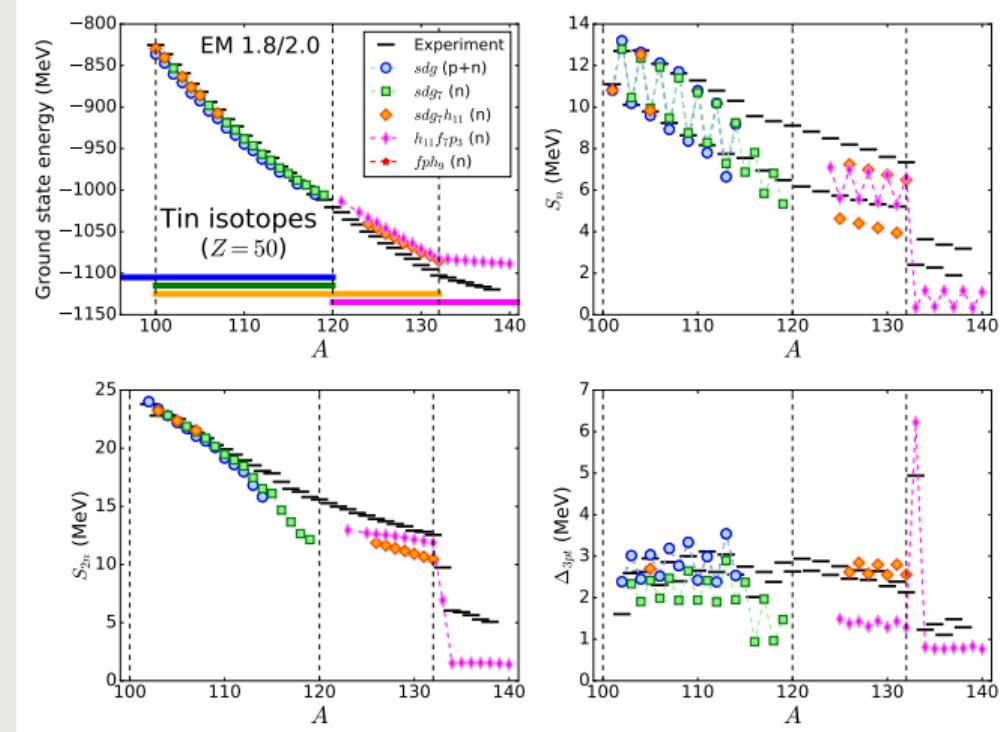
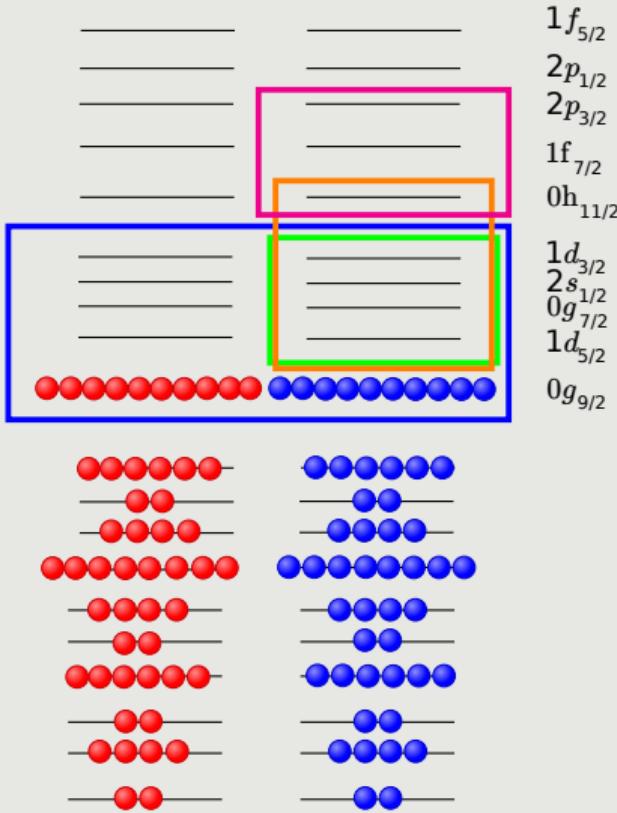
pf_5g_9 space, ^{76}Ge reference

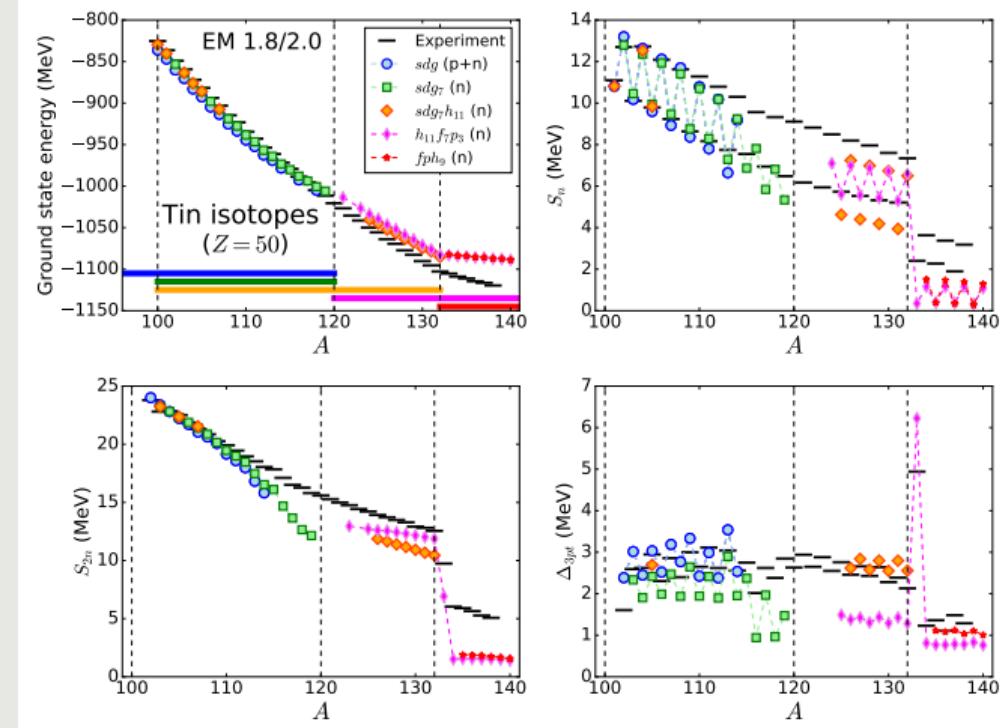
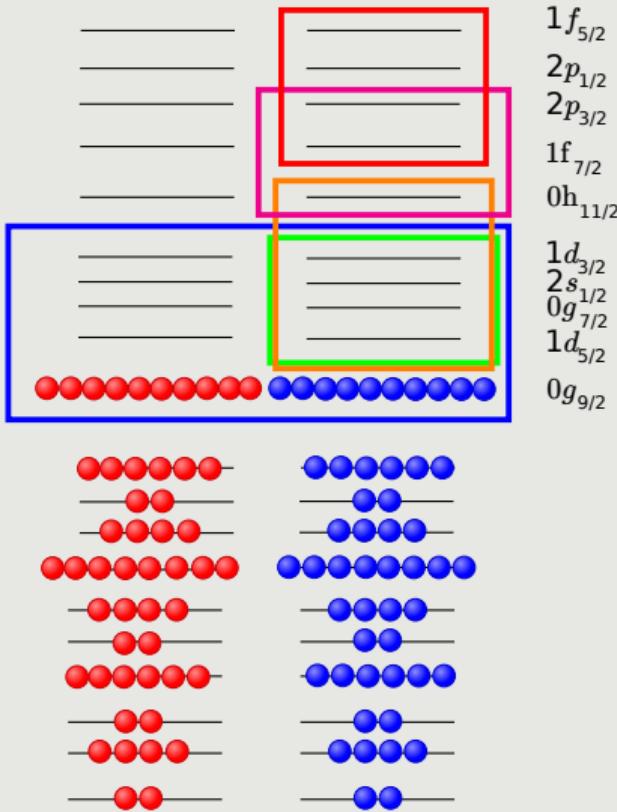


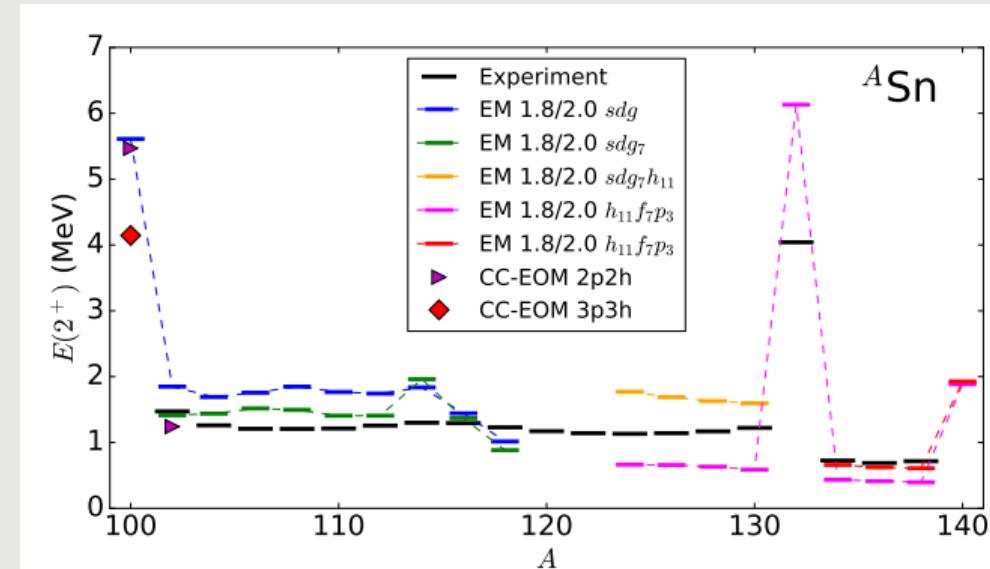
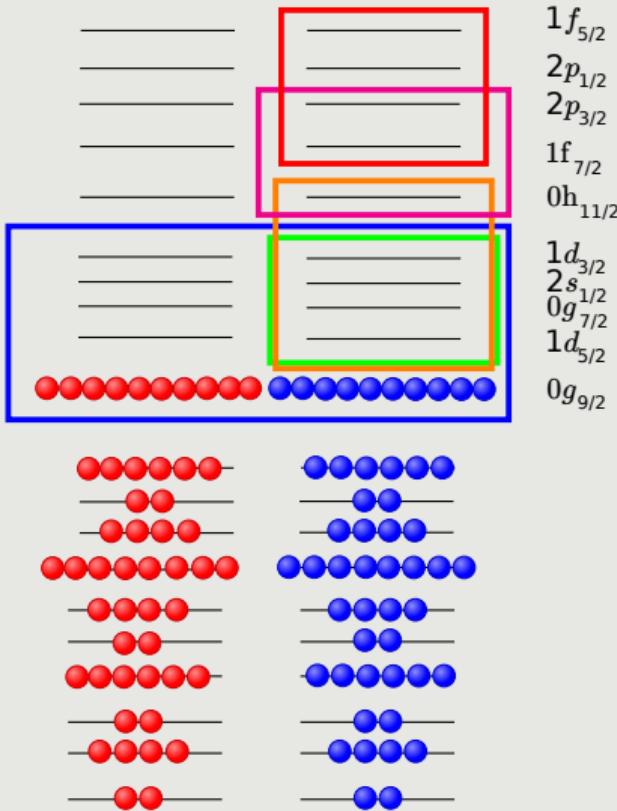


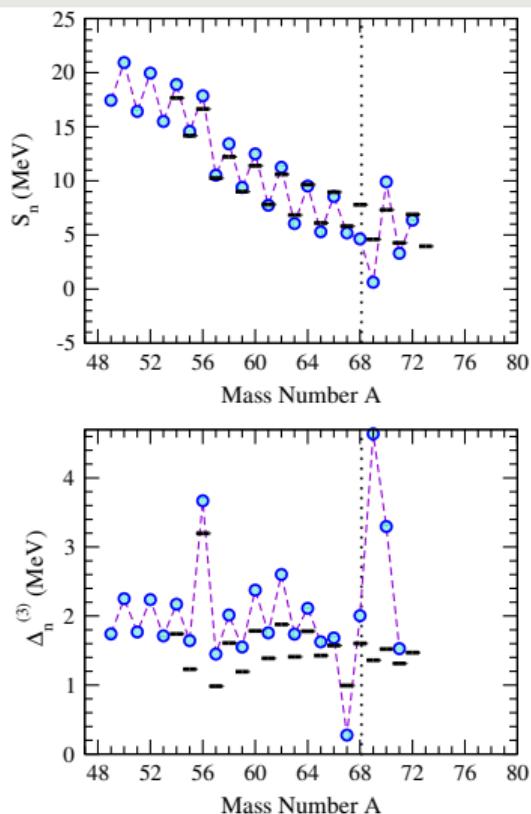
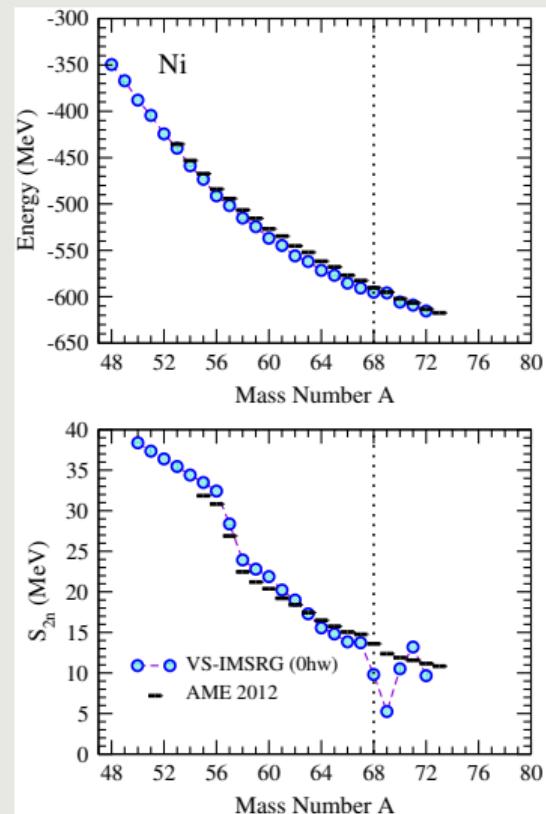
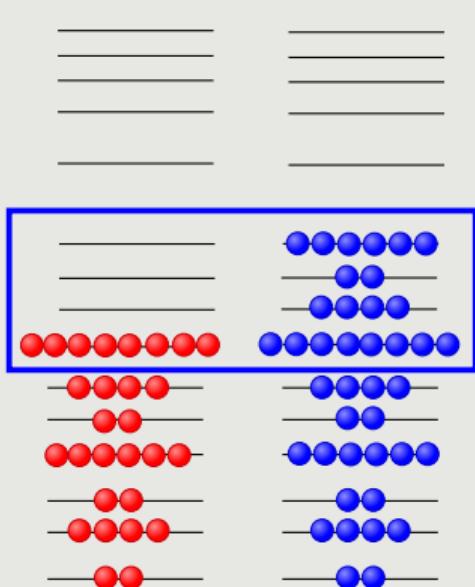


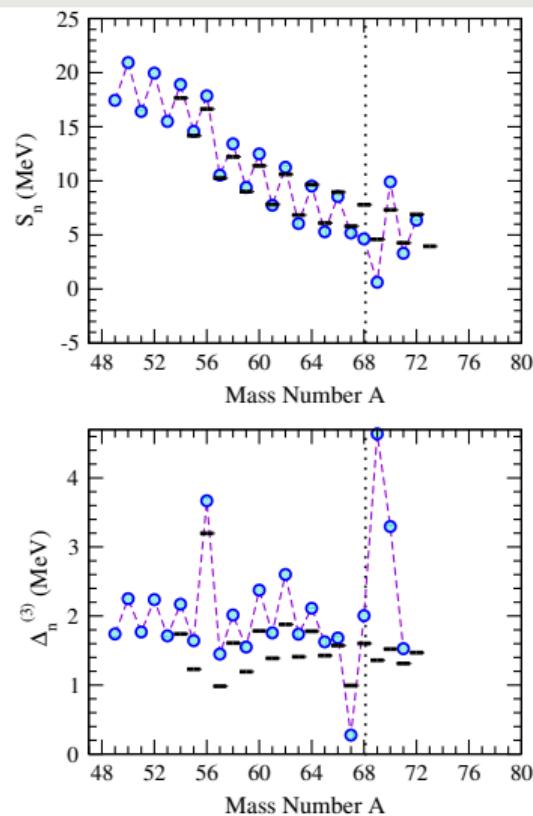
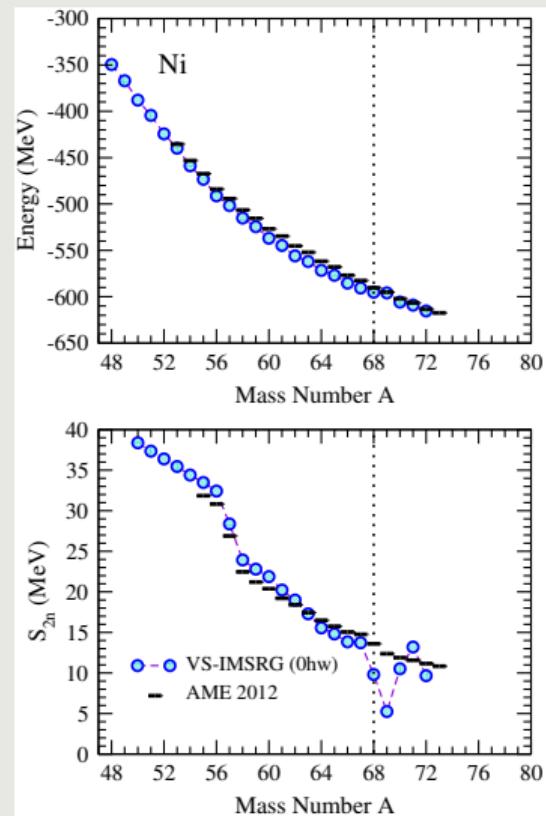
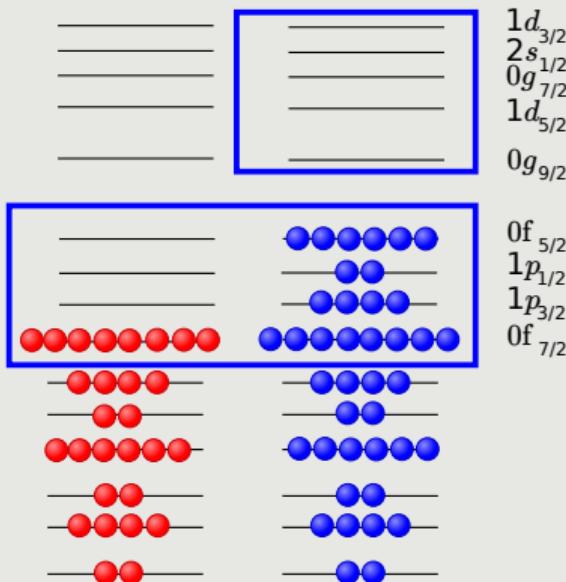


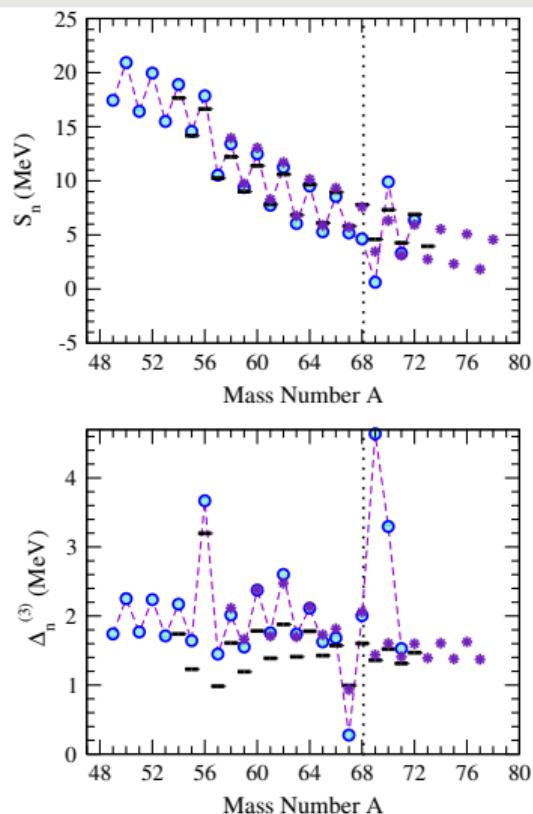
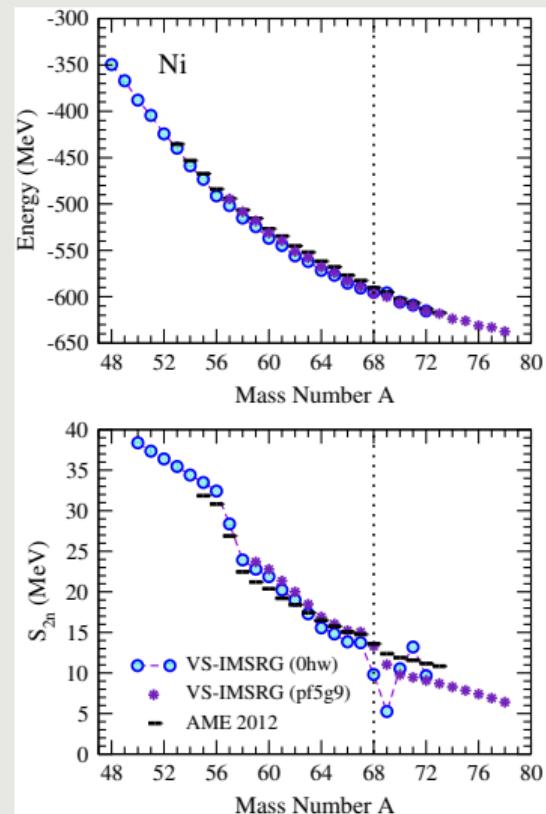
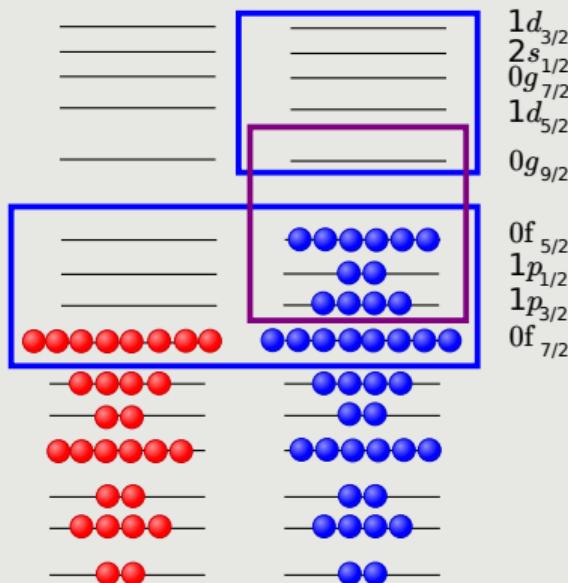




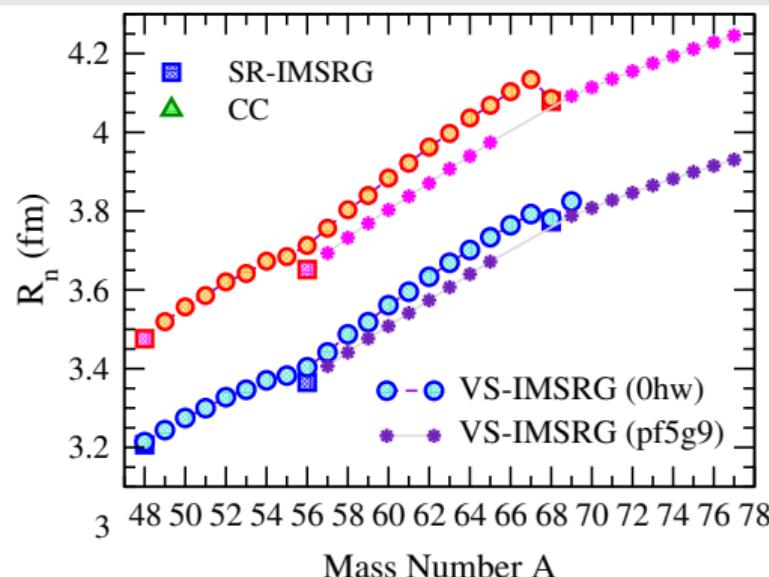
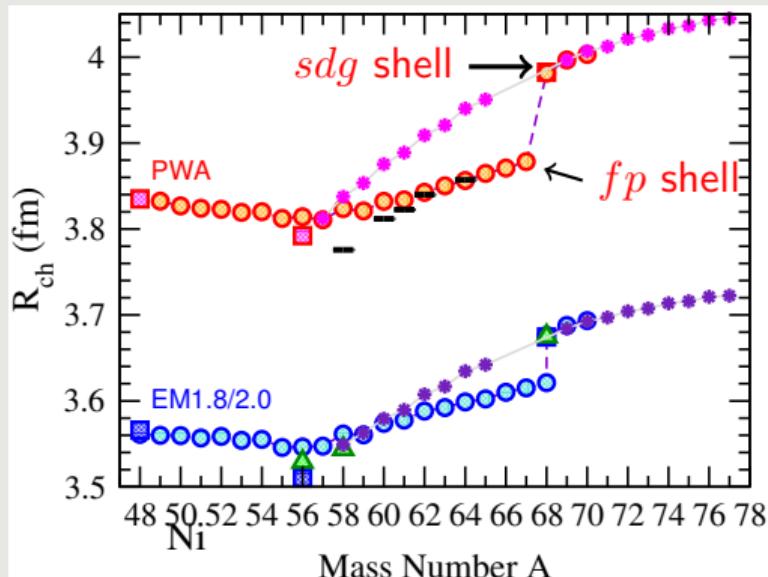




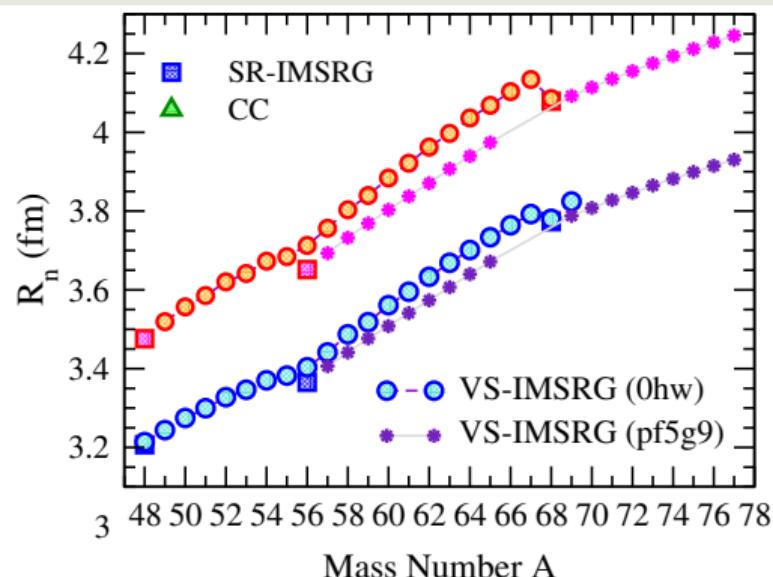
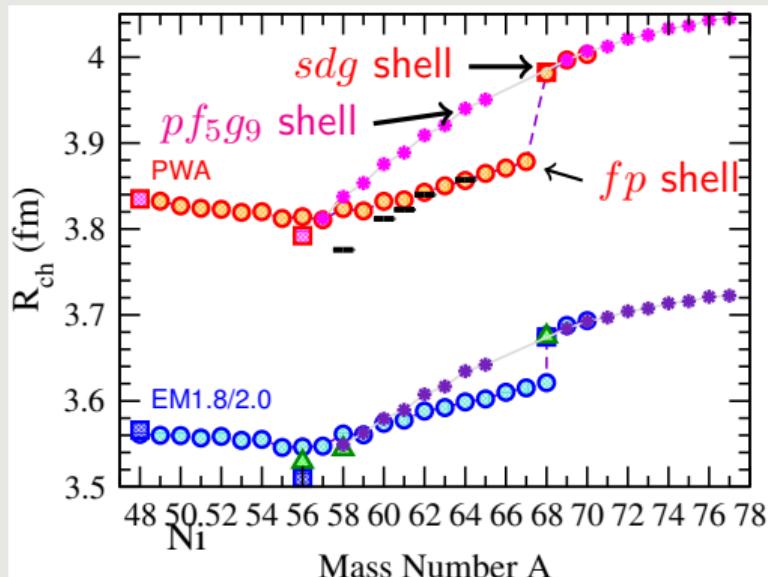


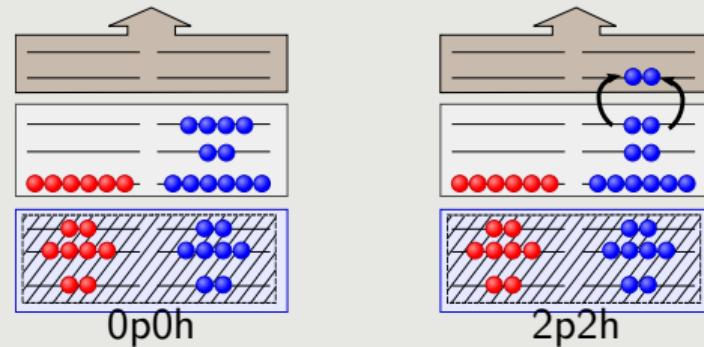


Radii of nickel isotopes

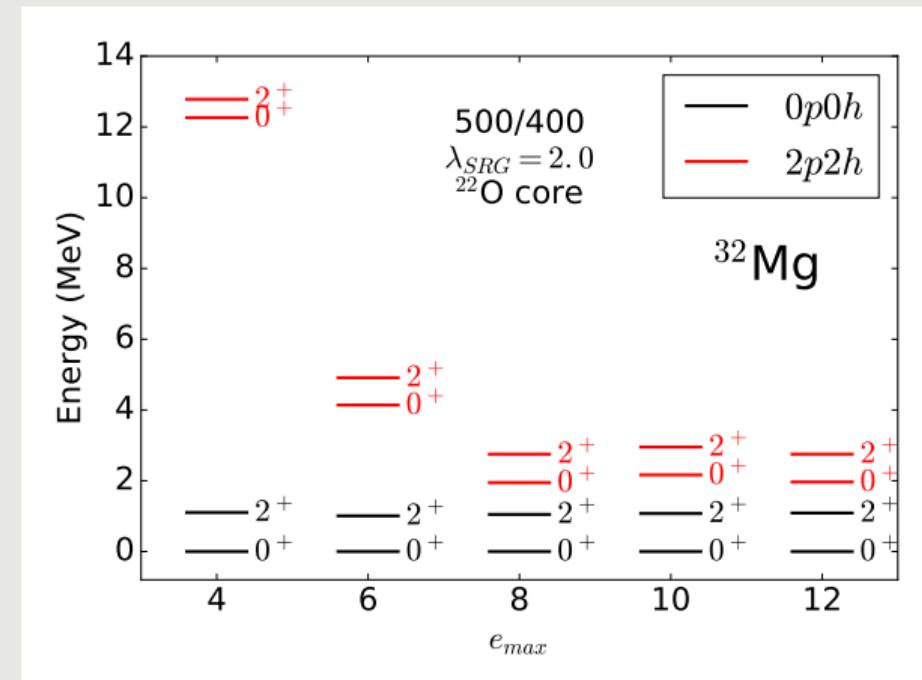


Radii of nickel isotopes



 ^{22}O core

$sd_3f_7p_3$ valence space for neutrons
 sd valence space for protons



Approximations:

- e_{max}, E_{3max}
- NO2B, IM-SRG(2)

“Arbitrary choices”:

- $\hbar\omega$ for oscillator basis
- Generator η
- Reference $|\Phi_0\rangle$ (or ρ)
- Valence space

- The “choices” should be made to minimize the amount left out by the approximations.
- Eliminate approximations → result independent of choices.
- Dependence on choices → estimate of approximation error
- Other ways to estimate error:
 - Extrapolations
 - Perturbative estimate of e.g. 3-body terms
 - Invariant trace?

Approximations:

- e_{max}, E_{3max}
- NO2B, IM-SRG(2)

“Arbitrary choices”:

- $\hbar\omega$ for oscillator basis
- Generator η
- Reference $|\Phi_0\rangle$ (or ρ)
- Valence space

- The “choices” should be made to minimize the amount left out by the approximations.
- Eliminate approximations → result independent of choices.
- Dependence on choices → estimate of approximation error
- Other ways to estimate error:
 - Extrapolations
 - Perturbative estimate of e.g. 3-body terms
 - Invariant trace?

Approximations:

- e_{max}, E_{3max}
- NO2B, IM-SRG(2)

“Arbitrary choices”:

- $\hbar\omega$ for oscillator basis
- Generator η
- Reference $|\Phi_0\rangle$ (or ρ)
- Valence space

- The “choices” should be made to minimize the amount left out by the approximations.
- Eliminate approximations → result independent of choices.
- Dependence on choices → estimate of approximation error
- Other ways to estimate error:
 - Extrapolations
 - Perturbative estimate of e.g. 3-body terms
 - Invariant trace?

Approximations:

- e_{max}, E_{3max}
- NO2B, IM-SRG(2)

“Arbitrary choices”:

- $\hbar\omega$ for oscillator basis
- Generator η
- Reference $|\Phi_0\rangle$ (or ρ)
- Valence space

- The “choices” should be made to minimize the amount left out by the approximations.
- Eliminate approximations → result independent of choices.
- Dependence on choices → estimate of approximation error
- Other ways to estimate error:
 - Extrapolations
 - Perturbative estimate of e.g. 3-body terms
 - Invariant trace?

- Valence space IM-SRG with ensemble normal ordering allows access to all nuclei up to $A \sim 100$
- Cost and accuracy comparable to closed-shell IM-SRG
- Consistent operators for transitions (see poster by N. Parzuchowski)
- Development of non-standard valence spaces extends this reach and improves results at the edge of the valence space
- Limit in A is due to E_{3max} truncation
- Calculations with multiple valence spaces probes truncation error

Collaborators:



A. Calci, J. Holt, P. Navrátil, C. Payne, O. Drozdowski,
D. Fullerton, C. Gwak, L. Kemmler, S. Leutheusser, D. Livermore



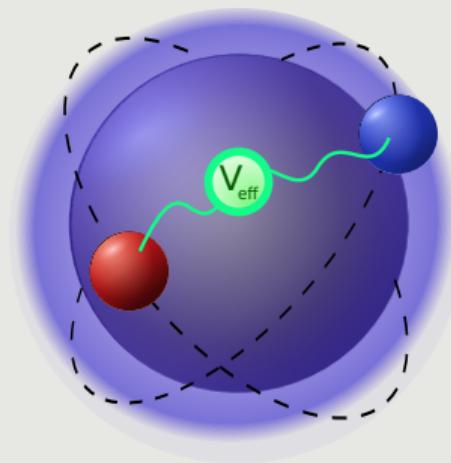
S. Bogner, H. Hergert, N. Parzuchowski



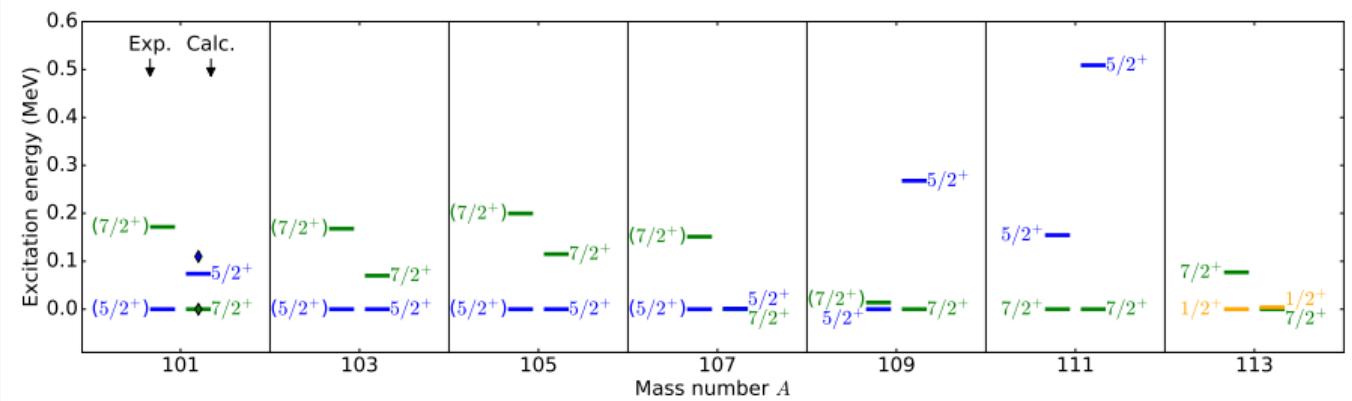
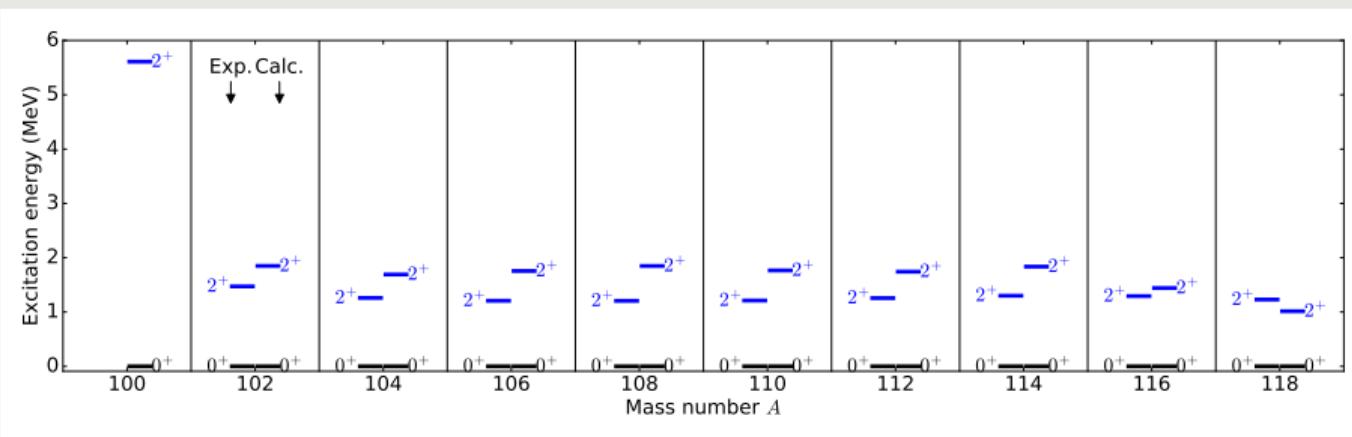
TU Darmstadt R. Roth, A. Schwenk, J. Simonis, C. Stumpf



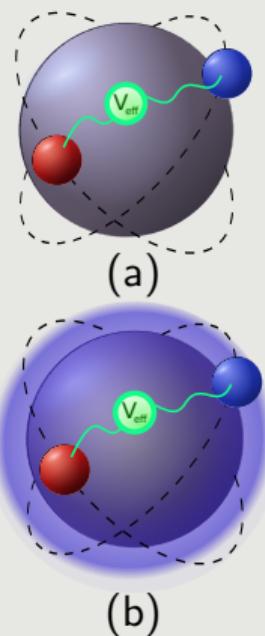
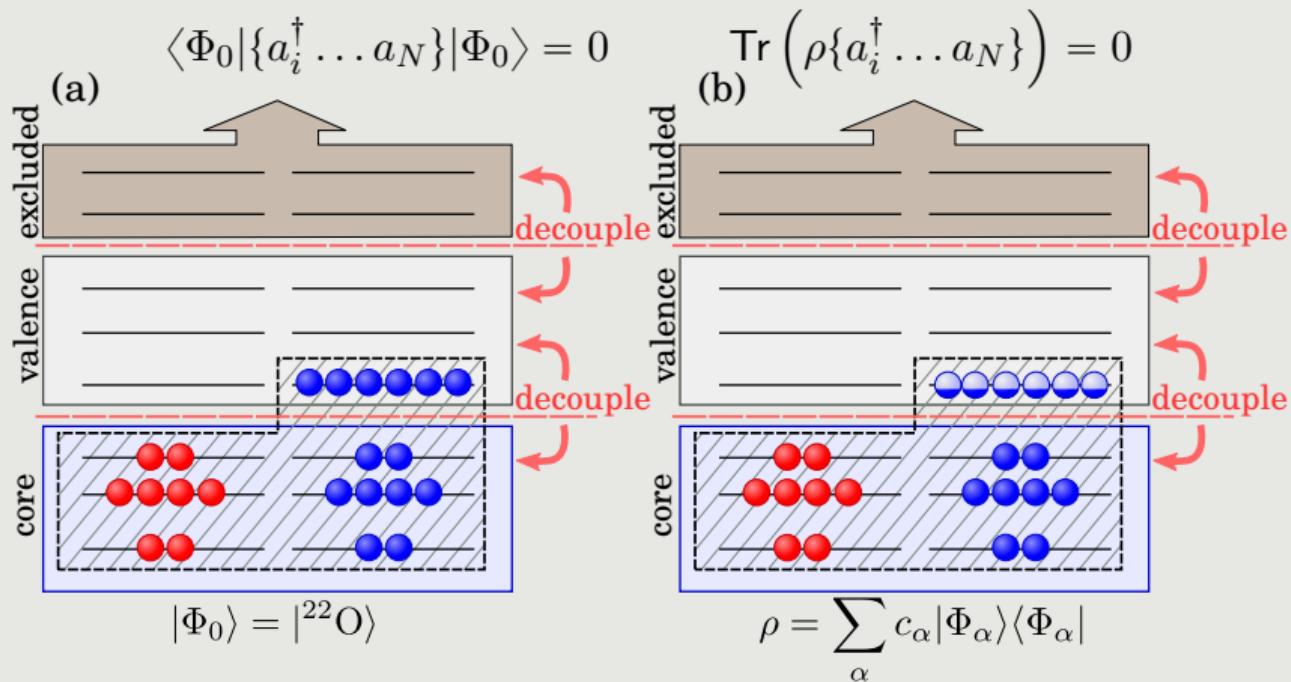
ORNL/UT G. Hagen, T. Morris



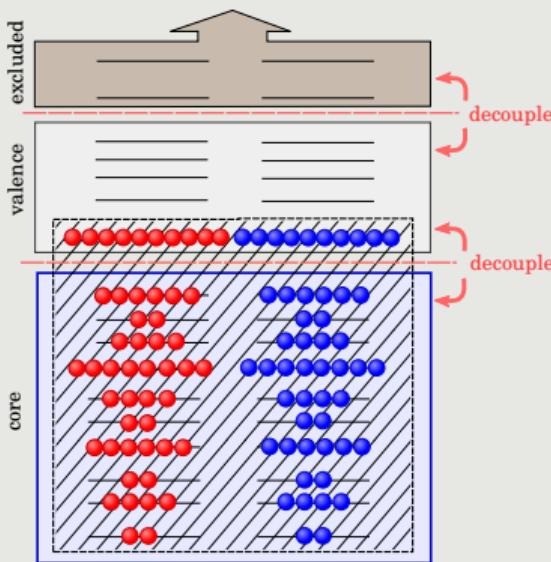
Backup slides



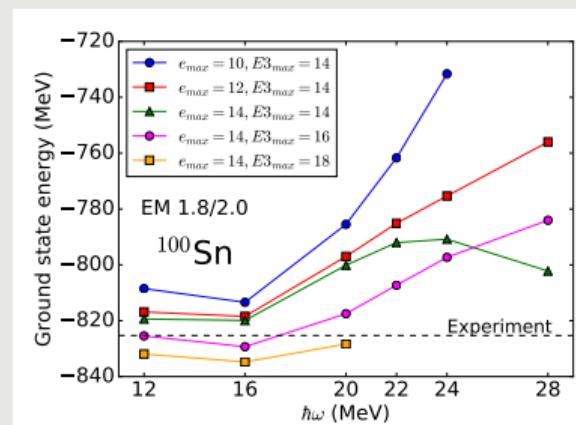
$$H = E_0 + \sum_{ij} f_{ij} \{a_i^\dagger a_j\} + \frac{1}{4} \sum_{ijkl} \Gamma_{ijkl} \{a_i^\dagger a_j^\dagger a_l a_k\} + \frac{1}{36} \sum_{ijklmn} W_{ijklmn} \{a_i^\dagger a_j^\dagger a_k^\dagger a_n a_m a_l\}$$



The tin isotopes ($Z = 50$)

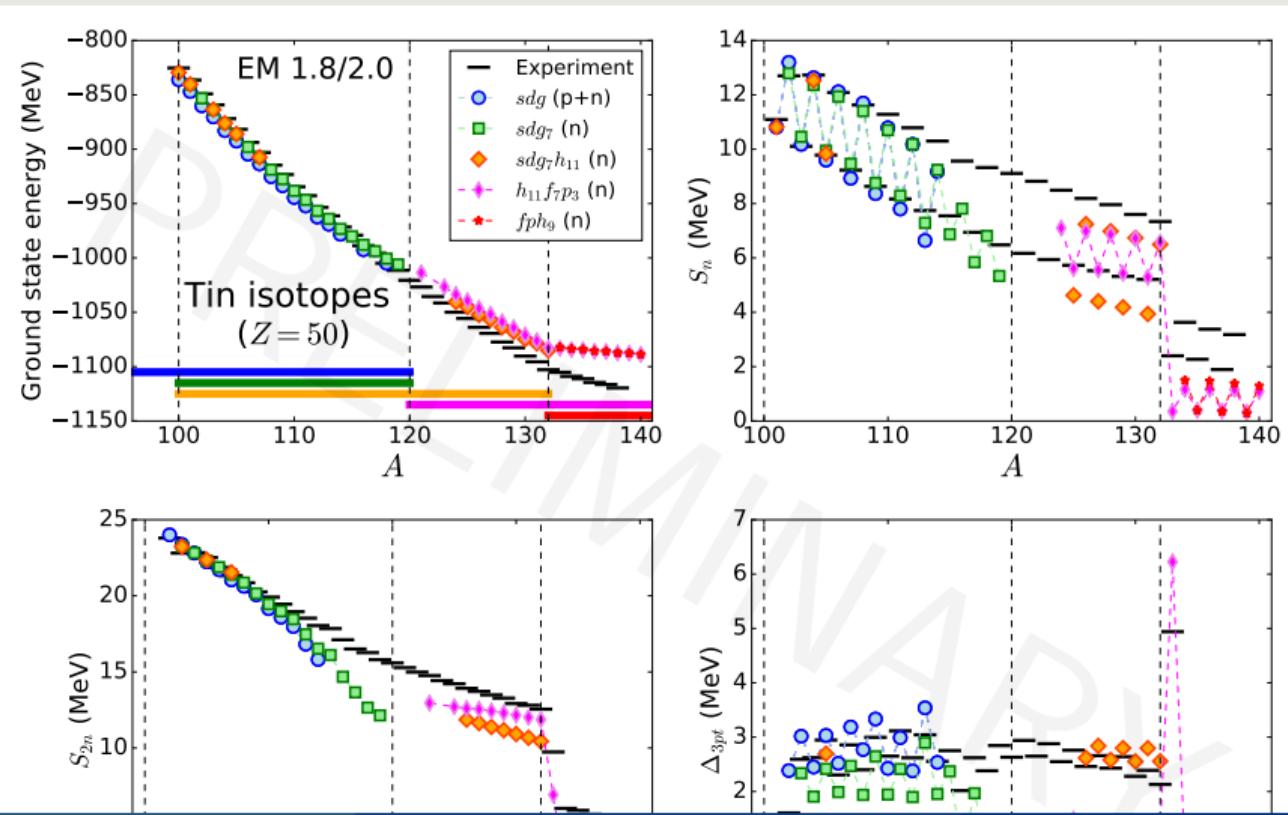


- ^{100}Sn : 50 protons, 50 neutrons
- Open shell valence space: full gds shell
- m -scheme dimension $\sim 10^{12}$
- Need importance truncation to diagonalize!

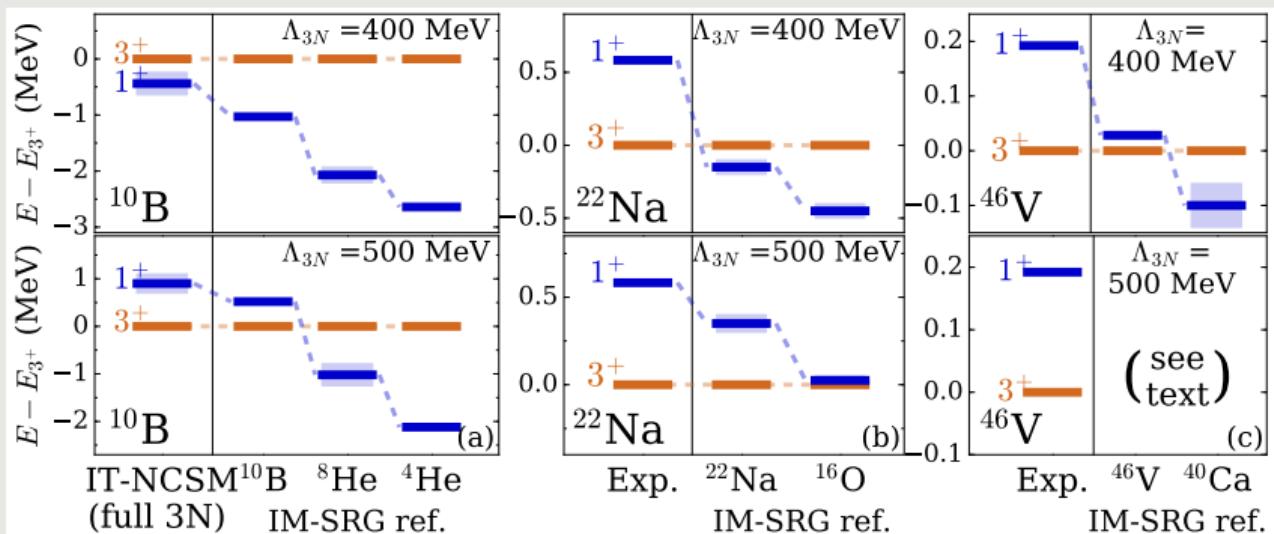
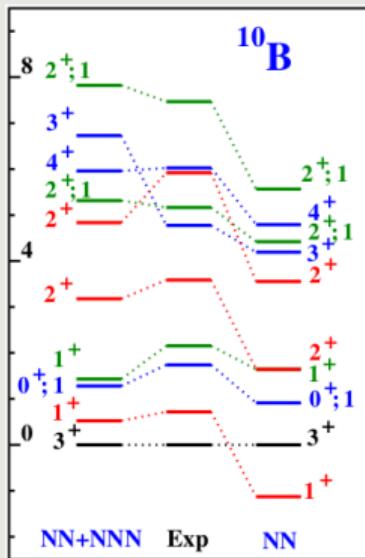


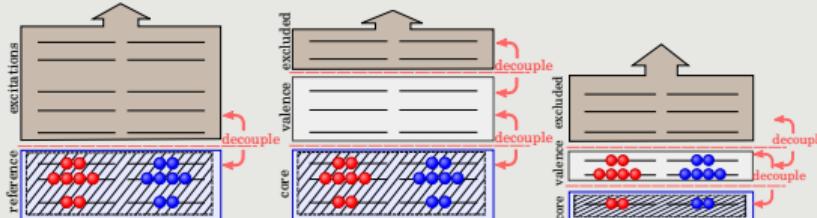
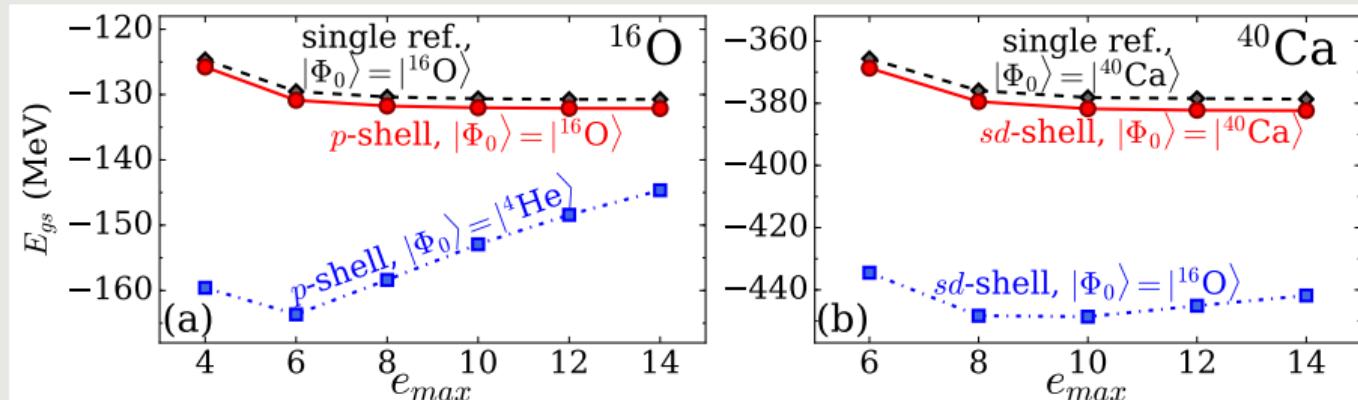
E_{3max}	Storage (GB)
14	5
16	20
18	100

Isotopic chain with $\hbar\omega = 16$, $e_{max} = 14$, $E3_{max} = 16$



Capturing valence 3N effects w/ NN machinery:

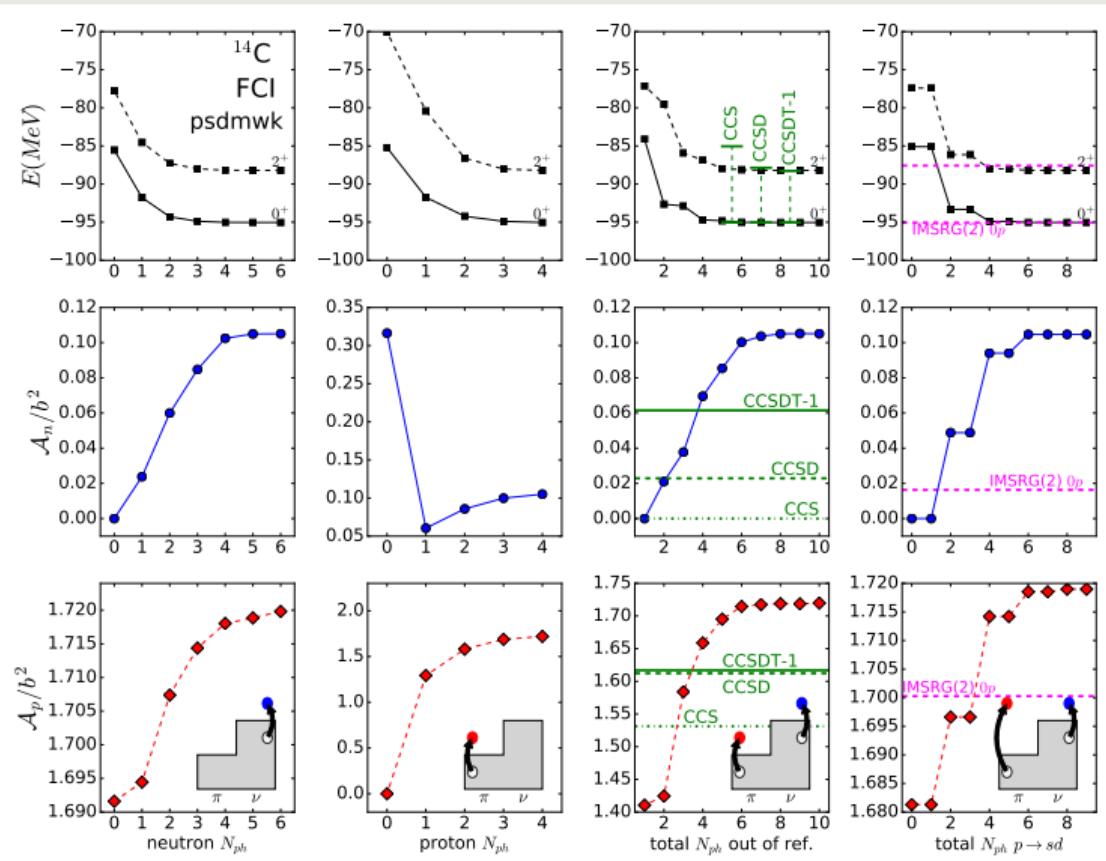




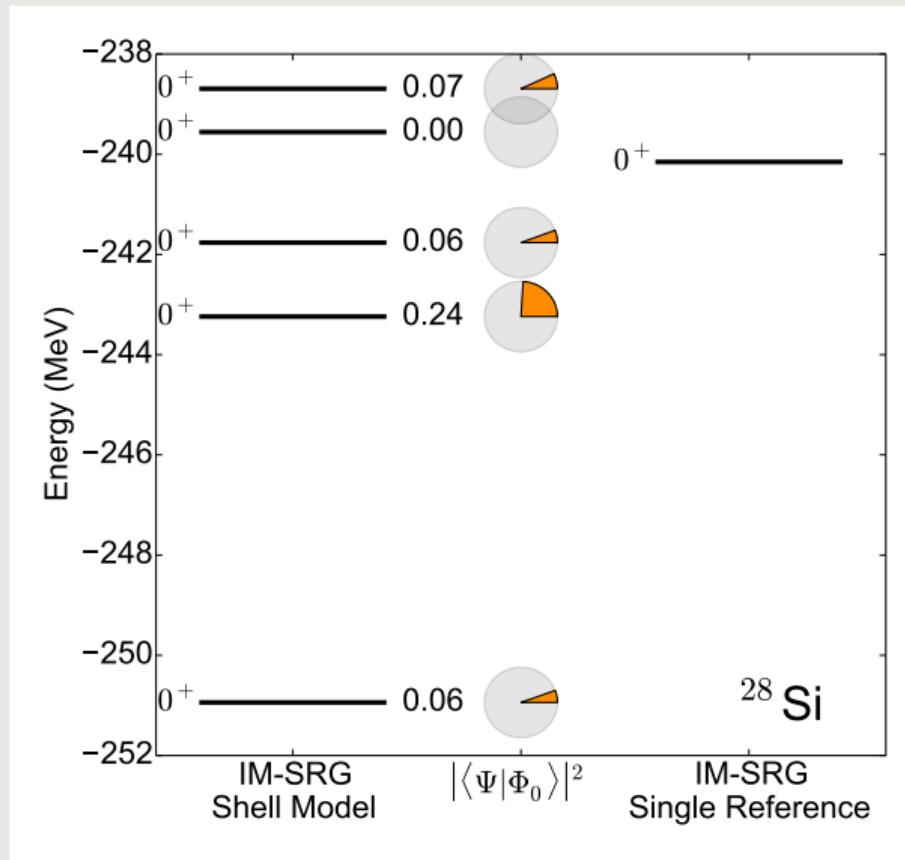
- Convergence not possible without proper normal ordering reference
- Two competing effects
 - Missing 3N forces
 - Bad single particle basis
- $\sim 1\%$ error due to additional decoupling

What's going on?

- Toy problem: ^{14}C , $p\text{-}sd$ space
- $A_p \equiv |\langle 0^+ \| r_p^2 Y_p^{(2)} \| 2^+ \rangle|$
- Truncate FCI in $N\text{p}-N\text{h}$ excitations
- Compare FCI with coupled cluster (from Gaute Hagen) and IM-SRG
 - CCSDT-1: singles+doubles+ \approx triples
- Missing p-h excitations: unimportant for energy, important for $E2$



Picking out spherical excited states



What *can't* we do, and why isn't everything perfect?

- **So far, limited to valence spaces defined by a single major oscillator shell**
 - No intruder states
 - No “island of inversion” states
 - No excited states of ^4He , ^{16}O , ^{28}O , ^{40}Ca , ^{60}Ca , ^{80}Zr (\leftarrow EOM can do these)
 - Max. 70 protons, 70 neutrons (oscillator magic numbers: 2, 8, 20, 40, 70, 112...)
- Large space limited to \sim 15 major oscillator shells (usually sufficient)
- **Limited to IM-SRG(2) approximation**
- Continuum states not included
- **Current input chiral interactions are not perfect**

Technical aside:

Recall the transformed Hamiltonian:

$$\tilde{H} = UHU^\dagger$$

Other operators may be transformed consistently.

If the operator \mathcal{O}^λ carries angular momentum λ , then

$$\tilde{\mathcal{O}}^\lambda = U\mathcal{O}^\lambda U^\dagger$$

$$e^{\Omega}\mathcal{O}^\lambda e^{-\Omega} = \mathcal{O}^\lambda + [\Omega, \mathcal{O}^\lambda] + \frac{1}{2}[\Omega, [\Omega, \mathcal{O}^\lambda]] + \dots$$

Only additional work is to derive angular momentum coupled commutator expressions (done).

Work in progress:

- Understand (and remedy) lack of $E2$ strength
- Understand quenching of Gamow-Teller strength
- Neutrinoless double beta decay
(C. Payne[†])
- Dark matter scattering
(S. Leutheusser*)
- Improve IM-SRG(2) approximation
- Applications to atomic systems
(D. Livermore*)

Potential projects:

- Can medium-mass nuclei provide a filter for chiral interactions?
- Unify reaction and structure – ab initio optical potentials, electron scattering
- Connections to DFT
 - Can ab initio calculations provide additional “data” for fitting?
- Can we explicitly calculate collective model parameters?

[†]M.Sc. student, *Undergraduate