Neutron matter based on consistently evolved chiral three-nucleon interactions

Kai Hebeler (OSU)

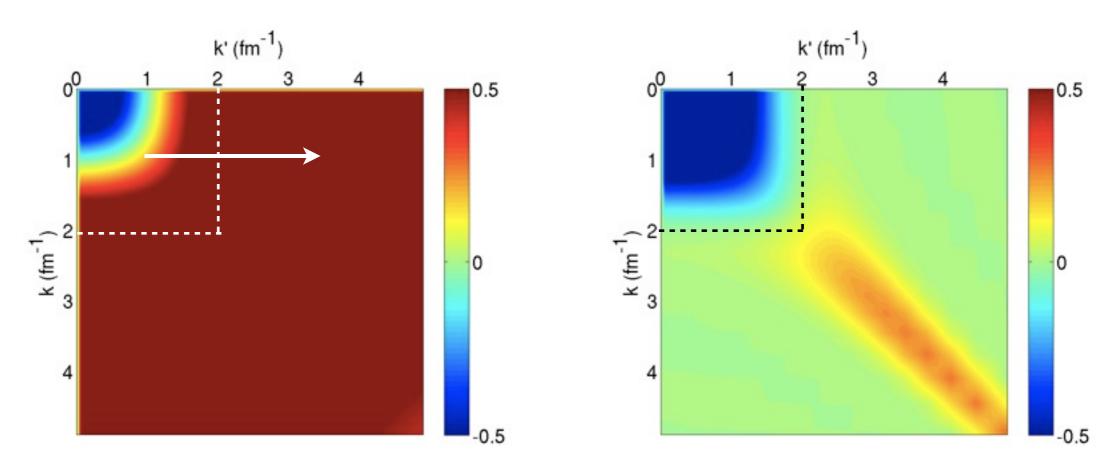
Progress in Ab Initio Techniques in Nuclear Physics

in collaboration with Dick Furnstahl

Vancouver, February 22, 2013



Changing the resolution: The Similarity Renormalization Group

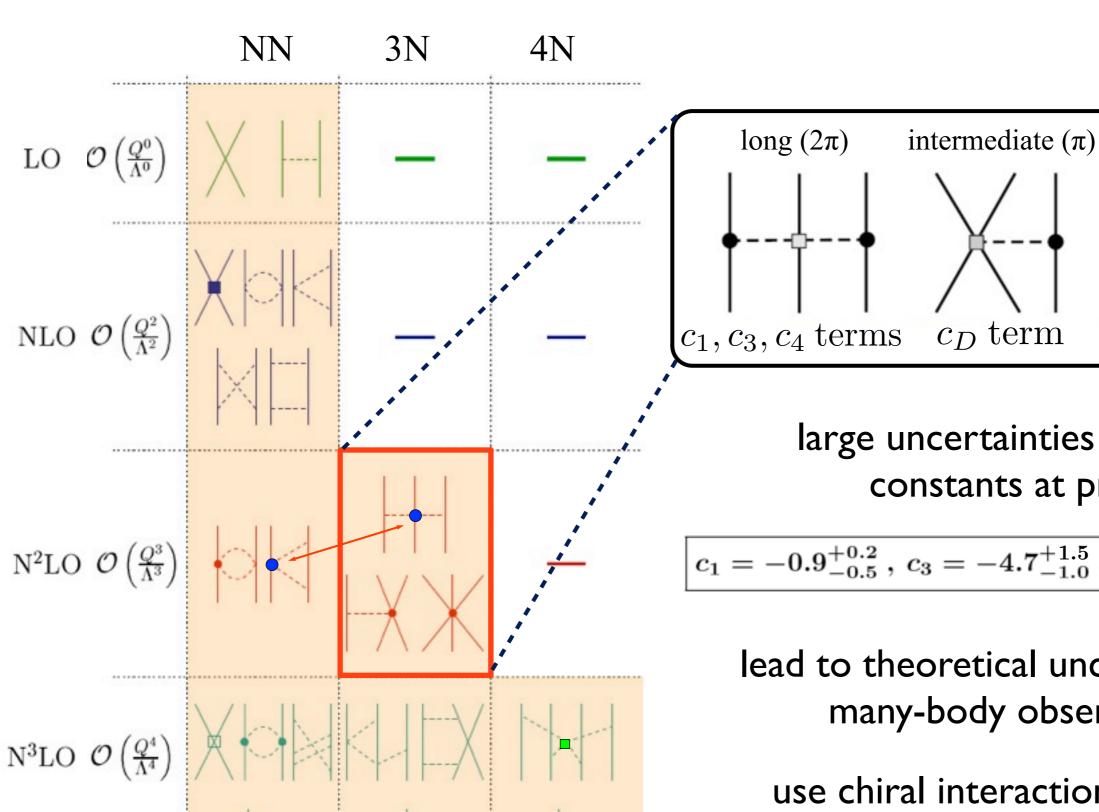


- elimination of coupling between low- and high momentum components, many-body problem more perturbative and tractable
- observables unaffected by resolution change (for exact calculations)
- residual resolution dependences can be used as tool to test calculations

Not the full story:

RG transformation also changes three-body (and higher-body) interactions.

Chiral EFT for nuclear forces, leading order 3N forces



large uncertainties in coupling constants at present:

short-range

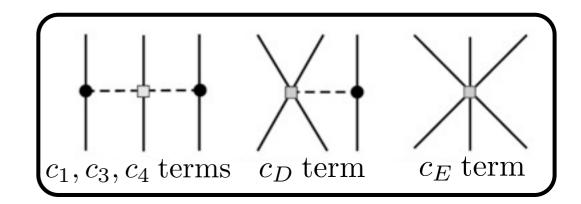
 c_E term

$$c_1 = -0.9^{+0.2}_{-0.5}$$
, $c_3 = -4.7^{+1.5}_{-1.0}$, $c_4 = 3.5^{+0.5}_{-0.2}$

lead to theoretical uncertainties in many-body observables

use chiral interactions as input for RG evolution

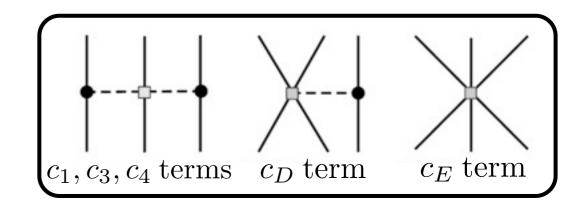
 So far (in momentum basis): intermediate (c_D) and short-range (c_E) 3NF couplings fitted to few-body systems at different resolution scales:



$$E_{^{3}\mathrm{H}} = -8.482\,\mathrm{MeV}$$
 and $r_{^{4}\mathrm{He}} = 1.464\,\mathrm{fm}$

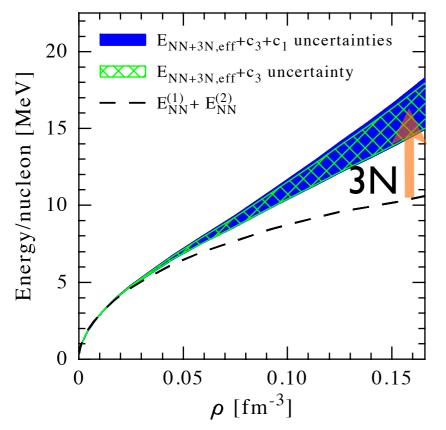
- coupling constants of natural size
- in neutron matter contributions from c_D , c_E and c_4 terms vanish
- ullet long-range 2π contributions assumed to be invariant under RG evolution

 So far (in momentum basis): intermediate (cD) and short-range (c_E) 3NF couplings fitted to few-body systems at different resolution scales:



$$E_{^{3}\mathrm{H}} = -8.482\,\mathrm{MeV}$$
 and $r_{^{4}\mathrm{He}} = 1.464\,\mathrm{fm}$

- coupling constants of natural size
- in neutron matter contributions from c_D , c_E and c_4 terms vanish
- ullet long-range 2π contributions assumed to be invariant under RG evolution

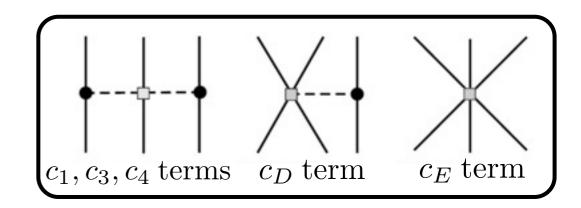


pure neutron matter

KH and Schwenk PRC 82, 014314 (2010)

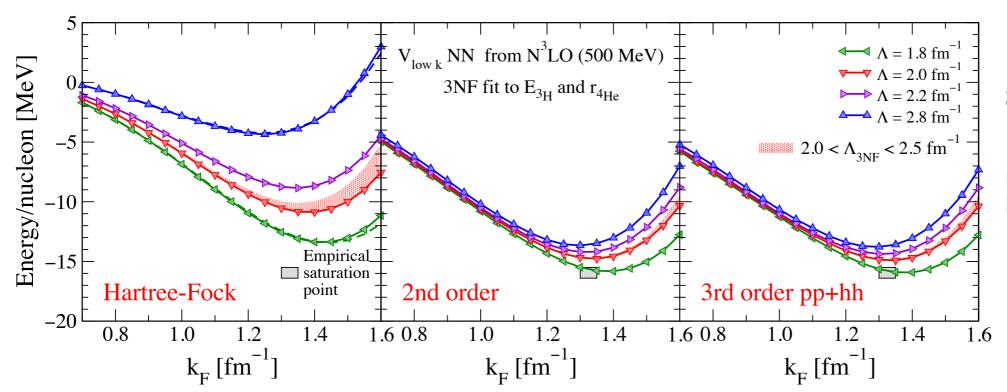
NN only, EGM $[10^{33} \mathrm{dyne/cm}^2]$ 2.0 1.5 1.0 0.5 neutron star matter 0.2 0.4 1.0 KH, Lattimer, Pethick, Schwenk, $\rho \left[\rho_0 \right]$ PRL 105, 161102 (2010)

 So far (in momentum basis): intermediate (cD) and short-range (cE) 3NF couplings fitted to few-body systems at different resolution scales:



$$E_{^{3}\mathrm{H}} = -8.482\,\mathrm{MeV}$$
 and $r_{^{4}\mathrm{He}} = 1.464\,\mathrm{fm}$

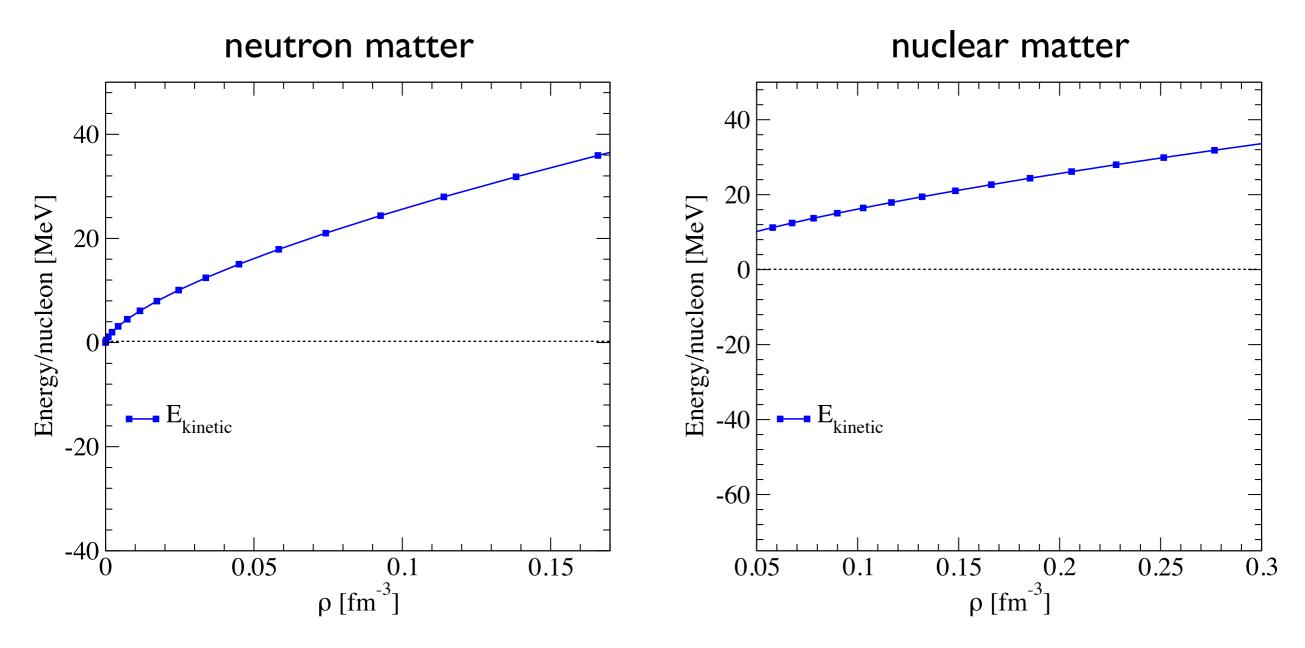
- coupling constants of natural size
- in neutron matter contributions from c_D , c_E and c_4 terms vanish
- ullet long-range 2π contributions assumed to be invariant under RG evolution



symmetric nuclear matter

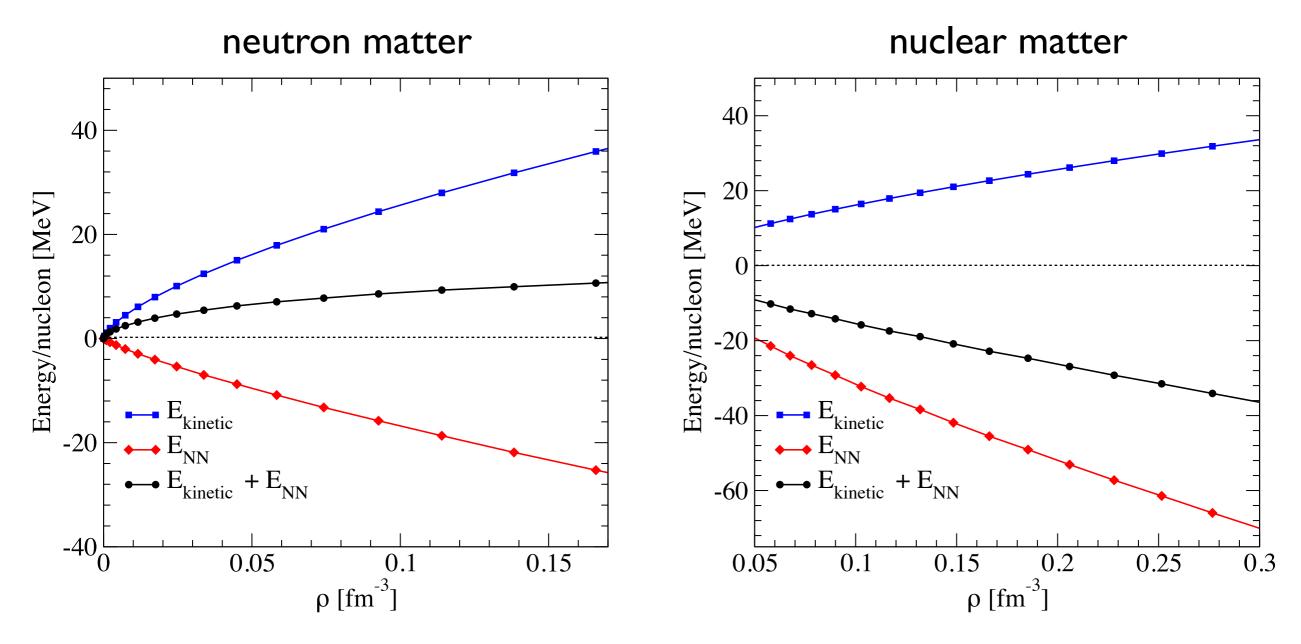
KH, Bogner, Furnstahl, Nogga, PRC(R) 83, 031301 (2011)

Hierarchy of many-body contributions



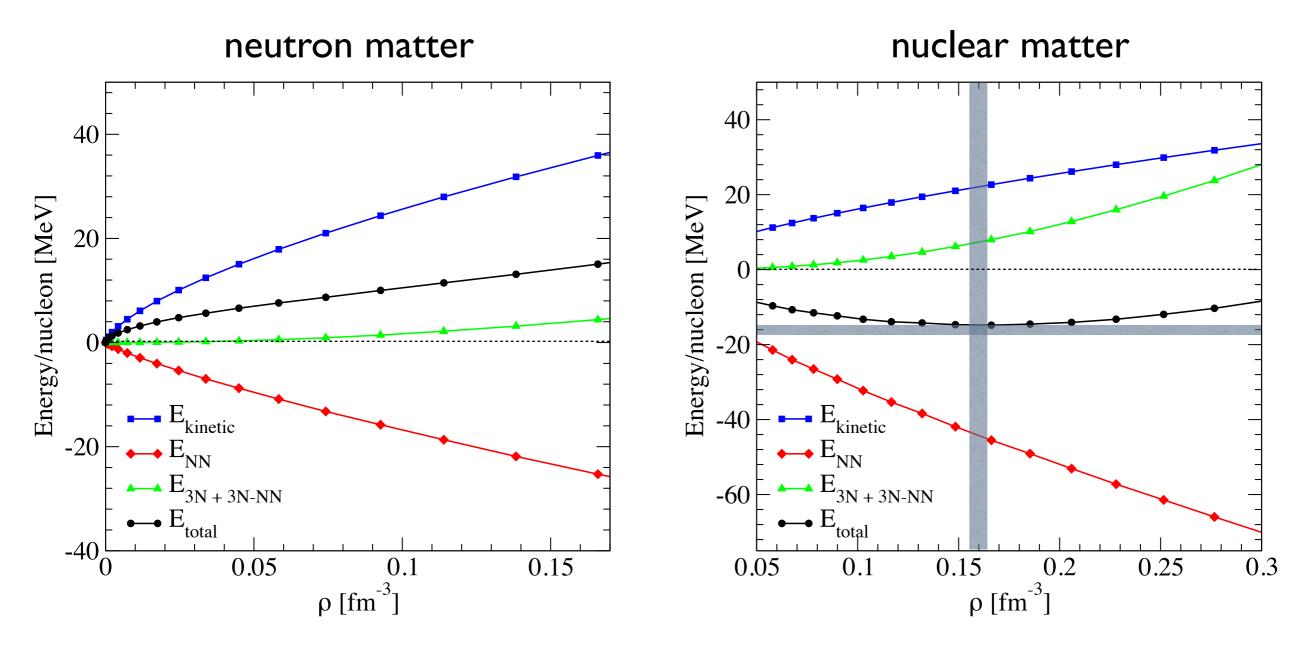
- binding energy results from cancellations of much larger kinetic and potential energy contributions
- chiral hierarchy of many-body terms preserved for considered density range
- ullet cutoff dependence of natural size, consistent with chiral exp. parameter $\sim 1/3$

Hierarchy of many-body contributions



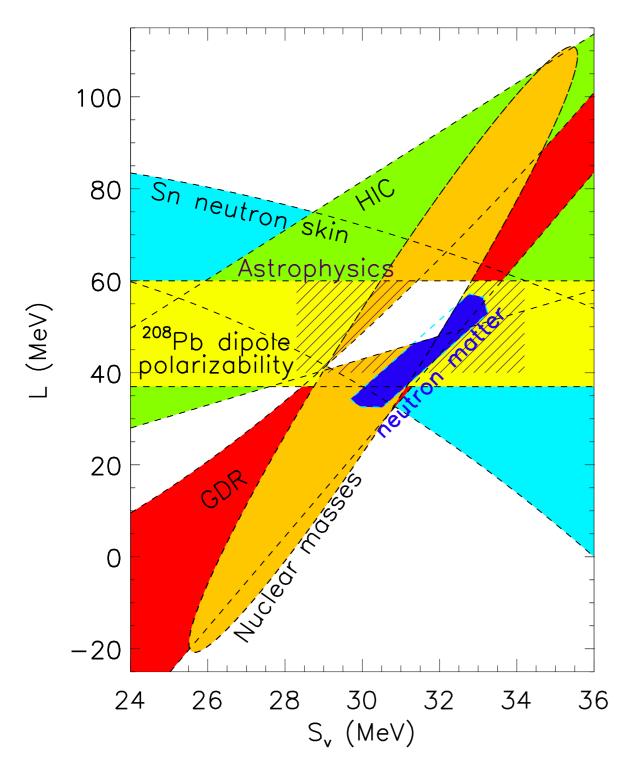
- binding energy results from cancellations of much larger kinetic and potential energy contributions
- chiral hierarchy of many-body terms preserved for considered density range
- ullet cutoff dependence of natural size, consistent with chiral exp. parameter $\sim 1/3$

Hierarchy of many-body contributions



- binding energy results from cancellations of much larger kinetic and potential energy contributions
- chiral hierarchy of many-body terms preserved for considered density range
- ullet cutoff dependence of natural size, consistent with chiral exp. parameter $\sim 1/3$

Symmetry energy constraints



extend EOS to finite proton fractions \boldsymbol{x}

and extract symmetry energy parameters

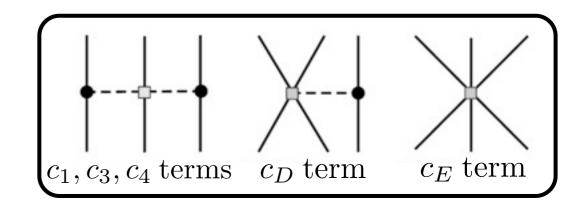
$$S_v = \frac{\partial^2 E/N}{\partial^2 x} \Big|_{\rho = \rho_0, x = 1/2}$$

$$L = \frac{3}{8} \left. \frac{\partial^3 E/N}{\partial \rho \partial^2 x} \right|_{\rho = \rho_0, x = 1/2}$$

KH, Lattimer, Pethick and Schwenk, in preparation

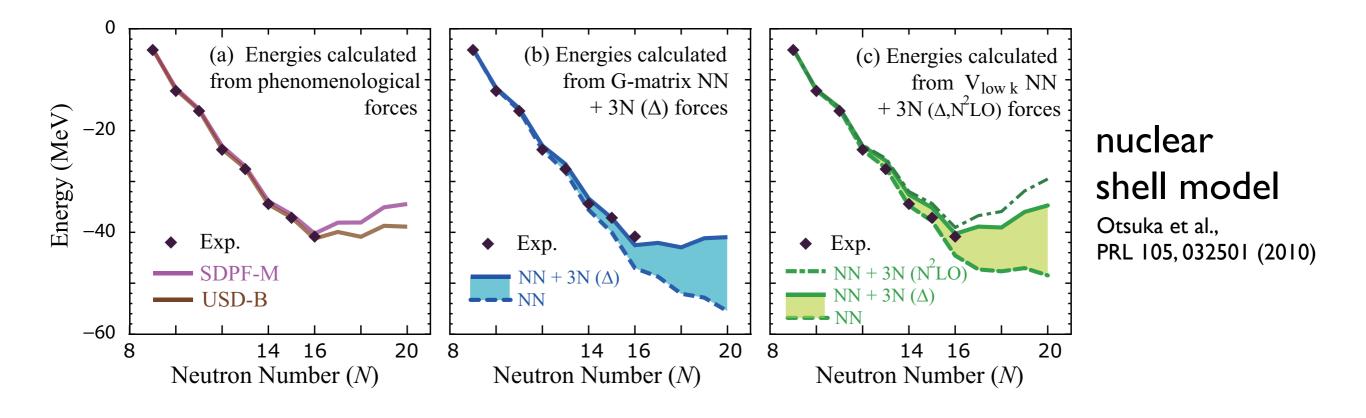
symmetry energy parameters consistent with other constraints

 So far (in momentum basis): intermediate (c_D) and short-range (c_E) 3NF couplings fitted to few-body systems at different resolution scales:

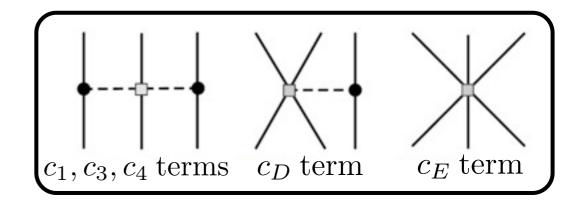


$$E_{^{3}\mathrm{H}} = -8.482\,\mathrm{MeV}$$
 and $r_{^{4}\mathrm{He}} = 1.464\,\mathrm{fm}$

- coupling constants of natural size
- in neutron matter contributions from c_D , c_E and c_4 terms vanish
- ullet long-range 2π contributions assumed to be invariant under RG evolution

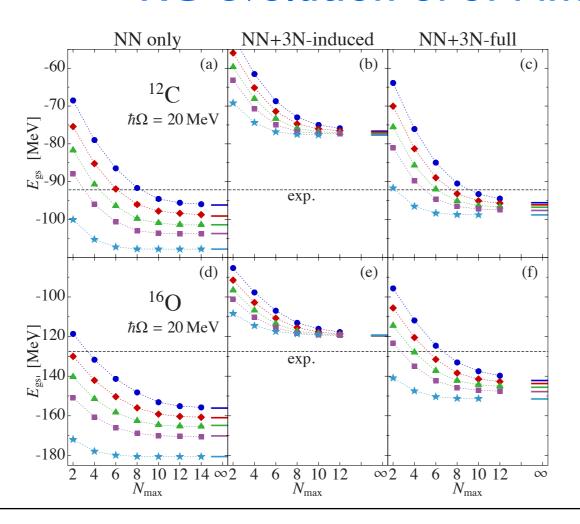


 So far (in momentum basis): intermediate (c_D) and short-range (c_E) 3NF couplings fitted to few-body systems at different resolution scales:



$$E_{^{3}\mathrm{H}} = -8.482\,\mathrm{MeV}$$
 and $r_{^{4}\mathrm{He}} = 1.464\,\mathrm{fm}$

- coupling constants of natural size
- in neutron matter contributions from c_D , c_E and c_4 terms vanish
- ullet long-range $\,2\pi$ contributions assumed to be invariant under RG evolution
- Ideal case: evolve 3NF consistently with NN interactions within the SRG
 - has been achieved in oscillator basis (Jurgenson, Roth)
 - promising results in very light nuclei
 - puzzling effects in heavier nuclei (higher-body forces?)
 - not immediately applicable to infinite systems
 - limitations on $\hbar\Omega$



Roth et al. PRL 107, 072501 (2011)

- Ideal case: evolve 3NF consistently with NN interactions within the SRG
 - has been achieved in oscillator basis (Jurgenson, Roth)
 - promising results in very light nuclei
 - puzzling effects in heavier nuclei (higher-body forces?)
 - not immediately applicable to infinite systems
 - limitations on $\hbar\Omega$

- application to infinite systems
 - equation of state (first results for neutron matter)
 - > systematic study of induced many-body contributions, scaling behavior
 - ▶ include initial N3LO 3N interactions

- application to infinite systems
 - equation of state (first results for neutron matter)
 - > systematic study of induced many-body contributions, scaling behavior
 - ▶ include initial N3LO 3N interactions

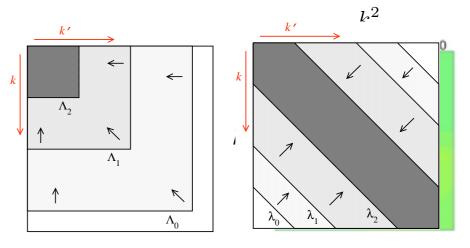
- application to infinite systems
 - equation of state (first results for neutron matter)
 - > systematic study of induced many-body contributions, scaling behavior
 - ▶ include initial N3LO 3N interactions
- transformation of evolved interactions to oscillator basis
 - ▶ application to finite nuclei, complimentary to HO evolution (no core shell model, coupled cluster, shell model)

with A. Calci, A. Ekstroem

- application to infinite systems
 - equation of state (first results for neutron matter)
 - > systematic study of induced many-body contributions, scaling behavior
 - ▶ include initial N3LO 3N interactions
- transformation of evolved interactions to oscillator basis
 - ▶ application to finite nuclei, complimentary to HO evolution (no core shell model, coupled cluster, shell model)

with A. Calci, A. Ekstroem

- study of various generators
 - ▶ different decoupling patterns (e.g. V_{low k})
 - improved efficiency of evolution
 - suppression of many-body forces?

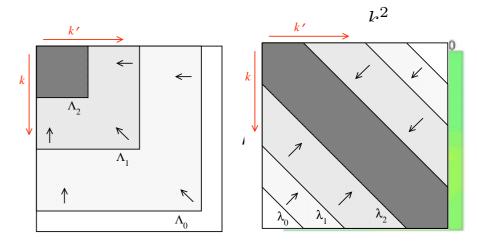


Anderson et al., PRC 77, 037001 (2008)

- application to infinite systems
 - equation of state (first results for neutron matter)
 - > systematic study of induced many-body contributions, scaling behavior
 - ▶ include initial N3LO 3N interactions
- transformation of evolved interactions to oscillator basis
 - ▶ application to finite nuclei, complimentary to HO evolution (no core shell model, coupled cluster, shell model)

with A. Calci, A. Ekstroem

- study of various generators
 - ▶ different decoupling patterns (e.g. V_{low k})
 - improved efficiency of evolution
 - suppression of many-body forces?



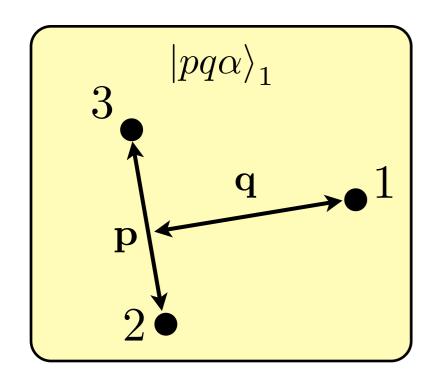
Anderson et al., PRC 77, 037001 (2008)

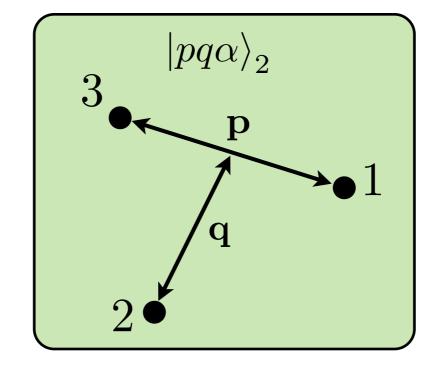
- explicit calculation of unitary 3N transformation
 - ▶ RG evolution of operators
 - > study of correlations in nuclear systems ----- factorization

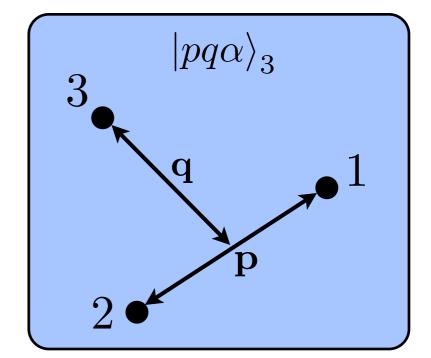
RG evolution of 3N interactions in momentum space

Three-body Faddeev basis:

$$|pq\alpha\rangle_i \equiv |p_iq_i; [(LS)J(ls_i)j] \mathcal{J}\mathcal{J}_z(Tt_i)\mathcal{T}\mathcal{T}_z\rangle$$







$$_{i}\langle pq\alpha|P|p'q'\alpha'\rangle_{i} =_{i}\langle pq\alpha|p'q'\alpha'\rangle_{i}$$

Faddeev bound-state equation:

$$|\psi_i\rangle = G_0 \left[2t_i P + (1 + t_i G_0) V_{3N}^i (1 + 2P) \right] |\psi_i\rangle$$

SRG flow equations of NN and 3N forces in momentum basis

$$\left(\begin{array}{c}
\frac{dH_s}{ds} = [\eta_s, H_s] & \eta_s = [T_{\text{rel}}, H_s]
\end{array}\right)$$

$$H = T + V_{12} + V_{13} + V_{23} + V_{123}$$

- ullet spectators correspond to delta functions, matrix representation of H_s ill-defined
- solution: explicit separation of NN and 3N flow equations

$$\frac{dV_{ij}}{ds} = [[T_{ij}, V_{ij}], T_{ij} + V_{ij}],$$

$$\frac{dV_{123}}{ds} = [[T_{12}, V_{12}], V_{13} + V_{23} + V_{123}]$$

$$+ [[T_{13}, V_{13}], V_{12} + V_{23} + V_{123}]$$

$$+ [[T_{23}, V_{23}], V_{12} + V_{13} + V_{123}]$$

$$+ [[T_{rel}, V_{123}], H_s]$$

 \bullet only connected terms remain in $\frac{dV_{123}}{ds}$,'dangerous' delta functions cancel

SRG evolution in momentum space

• evolve the antisymmetrized 3N interaction special thanks to J. Golak, R. Skibinski, K. Topolnicki

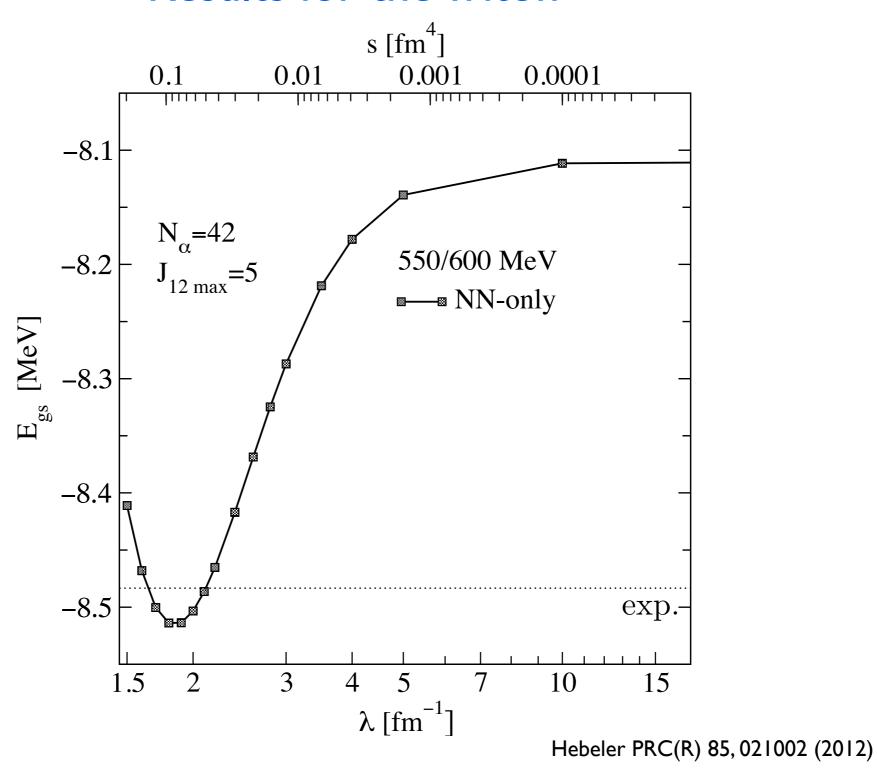
$$\overline{V}_{123} =_{i} \langle pq\alpha | (1 + P_{123} + P_{132}) V_{123}^{(i)} (1 + P_{123} + P_{132}) | p'q'\alpha' \rangle_{i}$$

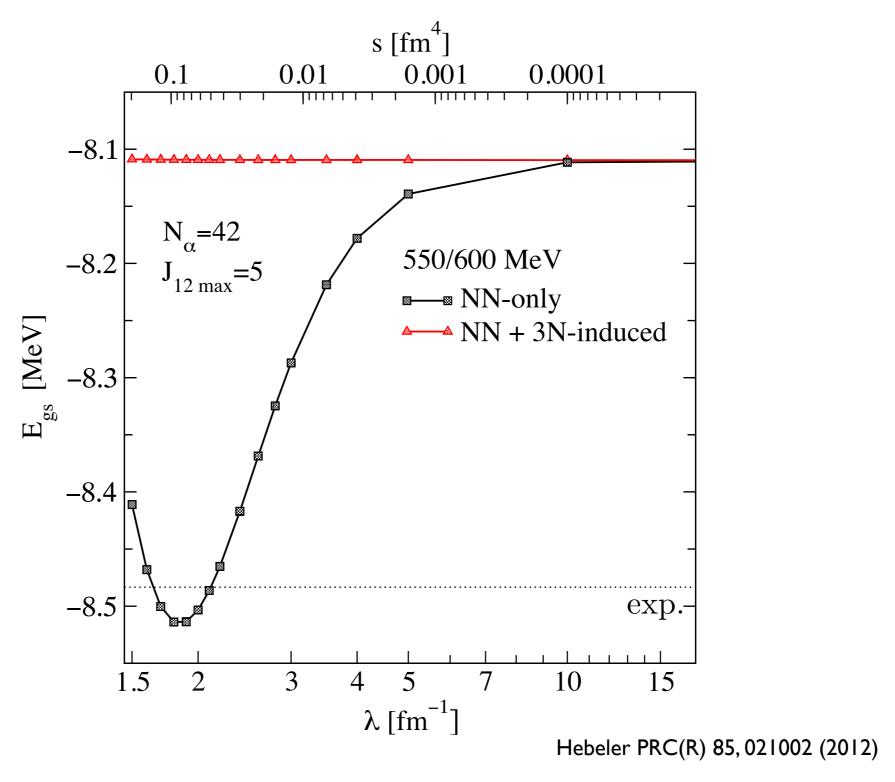
embed NN interaction in 3N basis:

$$V_{13}=P_{123}V_{12}P_{132},\quad V_{23}=P_{132}V_{12}P_{123}$$
 with $_3\langle pq\alpha|V_{12}|p'q'\alpha'\rangle_3=\langle p\tilde{\alpha}|V_{\rm NN}|p'\tilde{\alpha}'\rangle\,\delta(q-q')/q^2$

 \bullet use $\,P_{123}\overline{V}_{123}=P_{132}\overline{V}_{123}=\overline{V}_{123}$

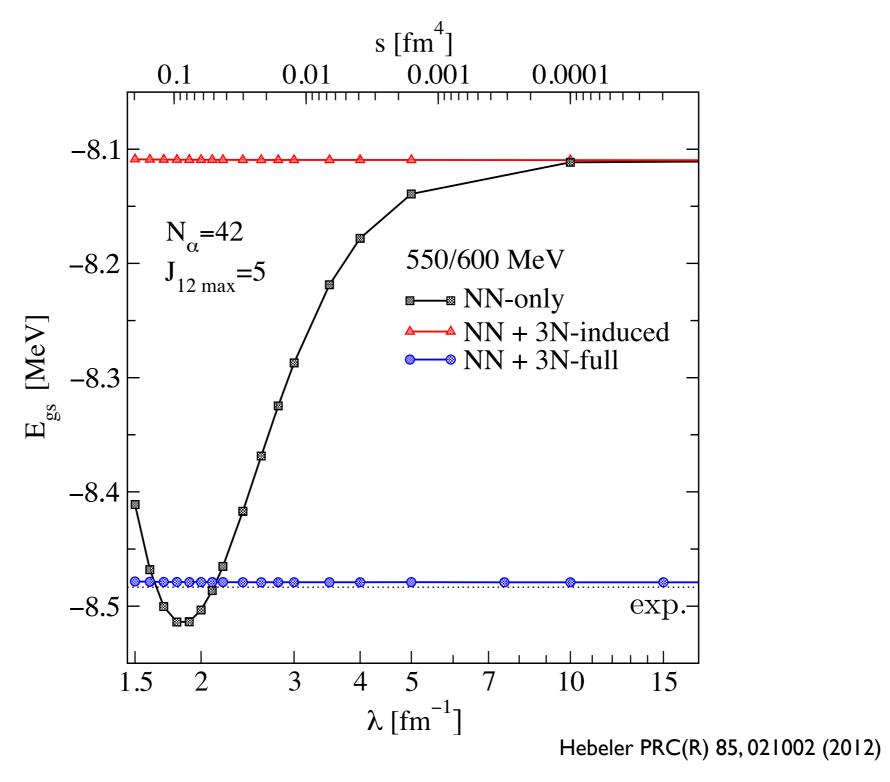
$$\Rightarrow d\overline{V}_{123}/ds = C_1(s, T, V_{NN}, P) + C_2(s, T, V_{NN}, \overline{V}_{123}, P) + C_3(s, T, \overline{V}_{123})$$





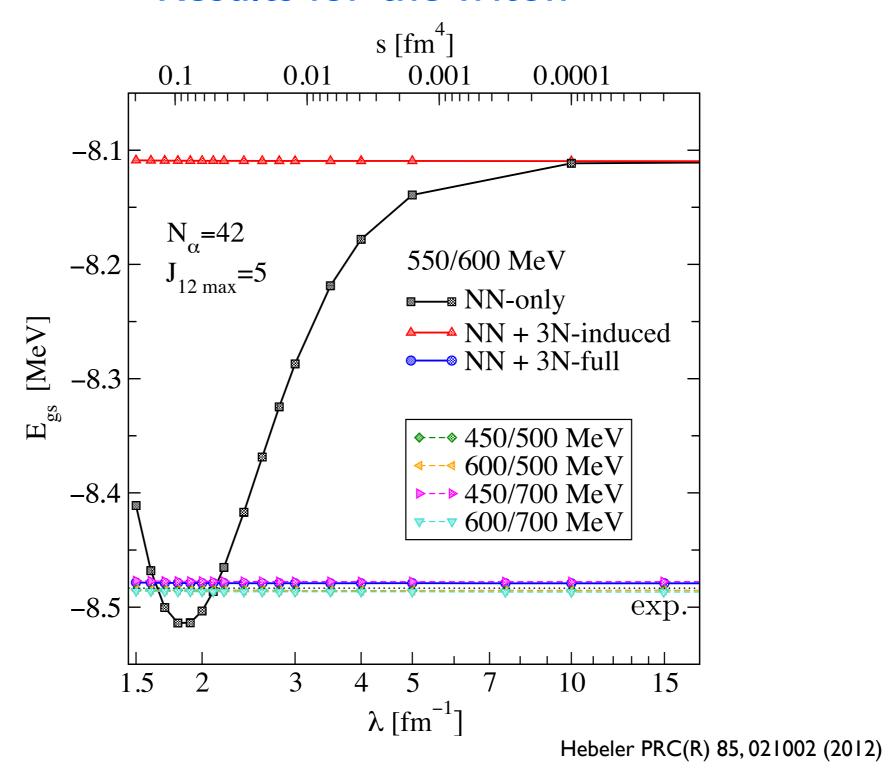
It works:

Invariance of $E_{\rm gs}^{^3\!H}$ within $\leq 1\,{\rm eV}$ for consistent chiral interactions at ${
m N}^2{
m LO}$



It works:

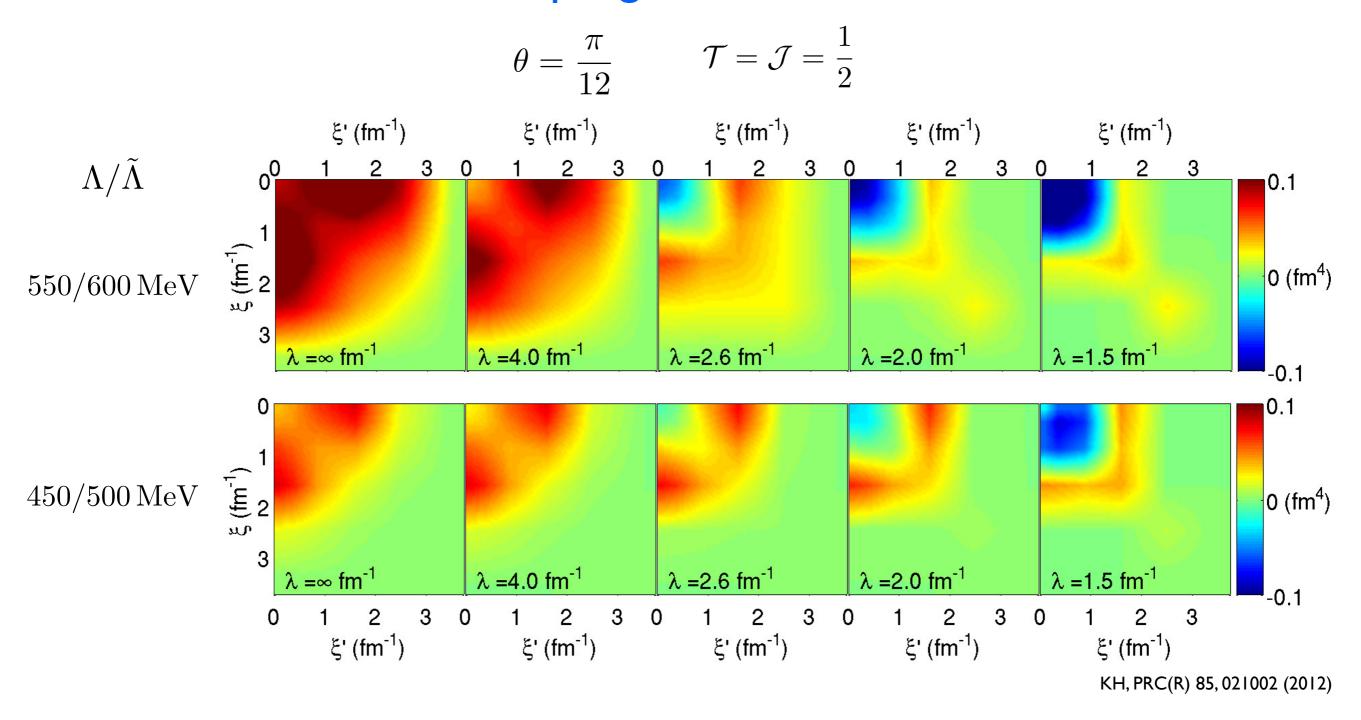
Invariance of $E_{\rm gs}^{^3\!H}$ within $\leq 1\,{\rm eV}$ for consistent chiral interactions at ${
m N}^2{
m LO}$



It works:

Invariance of $E_{\rm gs}^{^3\!H}$ within $\leq 1\,{\rm eV}$ for consistent chiral interactions at ${
m N}^2{
m LO}$

Decoupling of matrix elements



hyperradius:
$$\xi^{2} = p^{2} + \frac{3}{4}q^{2}$$

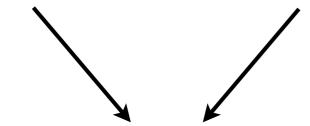
hyperangle:
$$\tan \theta = \frac{2 p}{\sqrt{3} q}$$

same decoupling patterns like in NN interactions

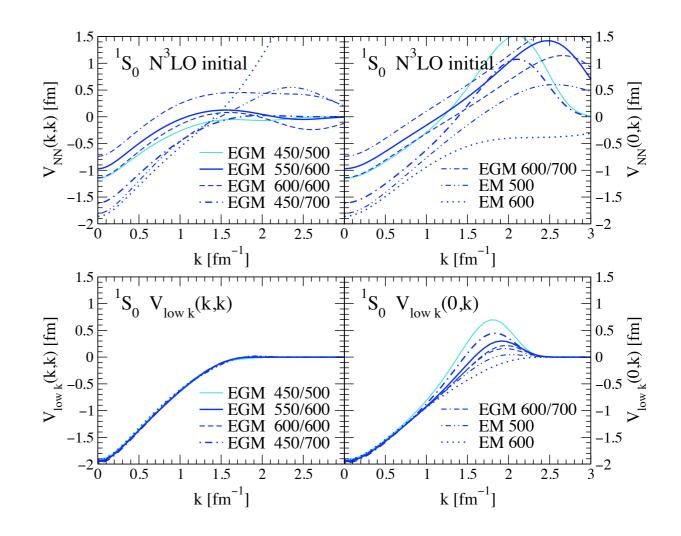
$k \text{ [fm}^{-1}]$ $k \text{ [fm}^{-1}]$

Universality in 3N into





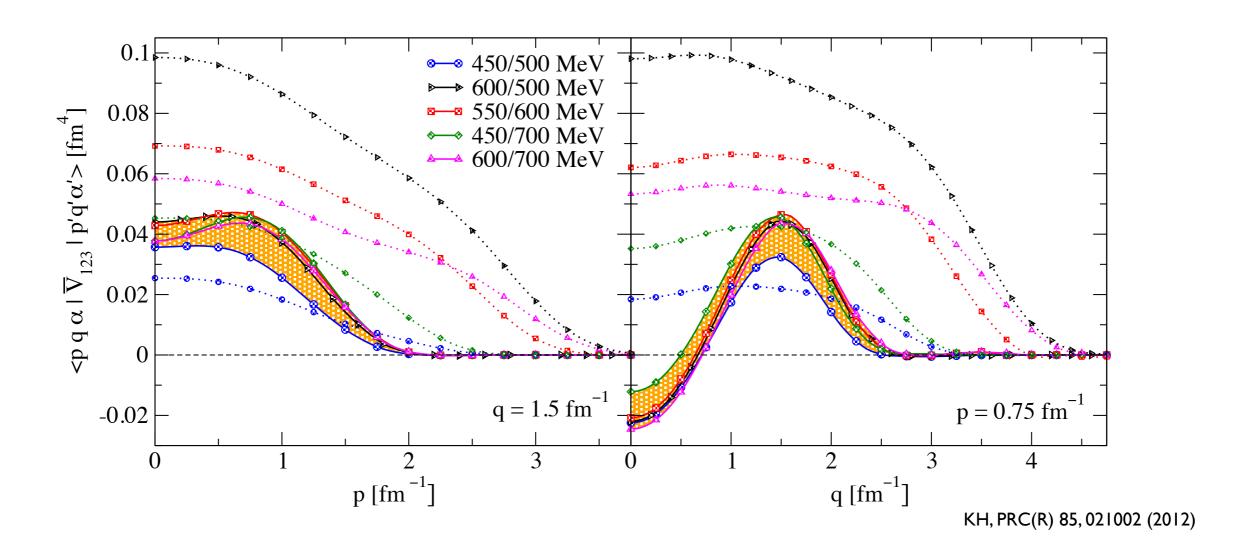
(approximate) universality of low-resolution NN interactions



To what extent are 3N interactions constrained at low resolution?

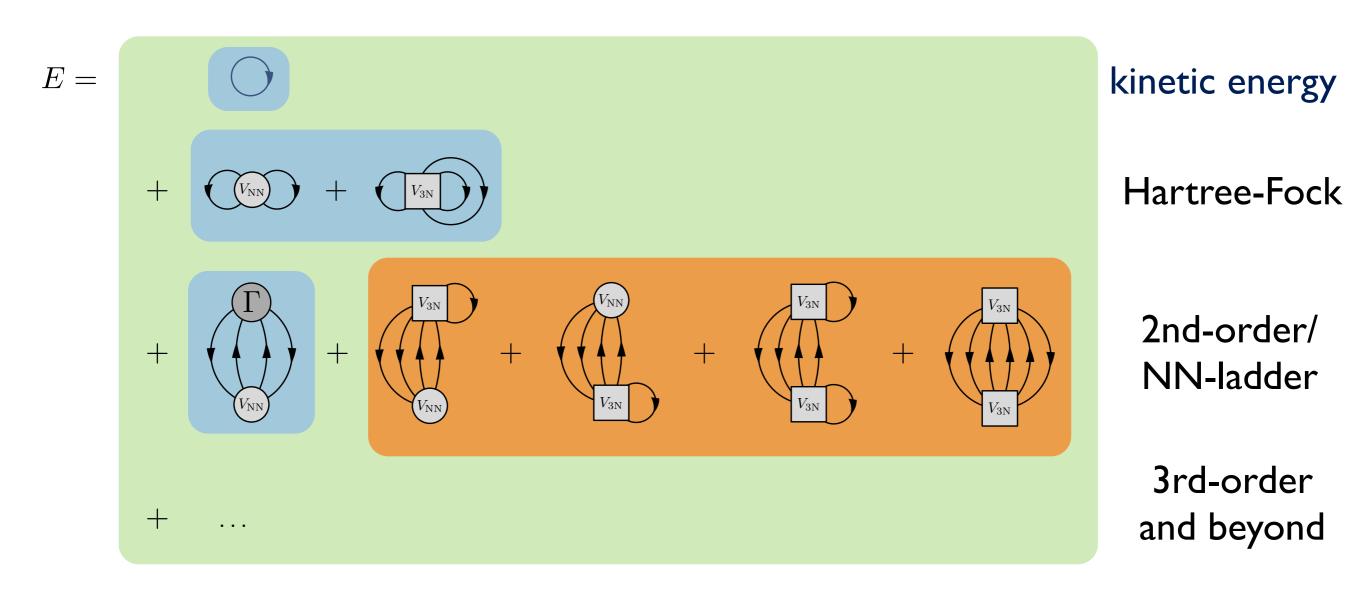
- ullet only two low-energy constants c_D and c_E
- 3N interactions give only subleading contributions to observables

Universality in 3N interactions at low resolution

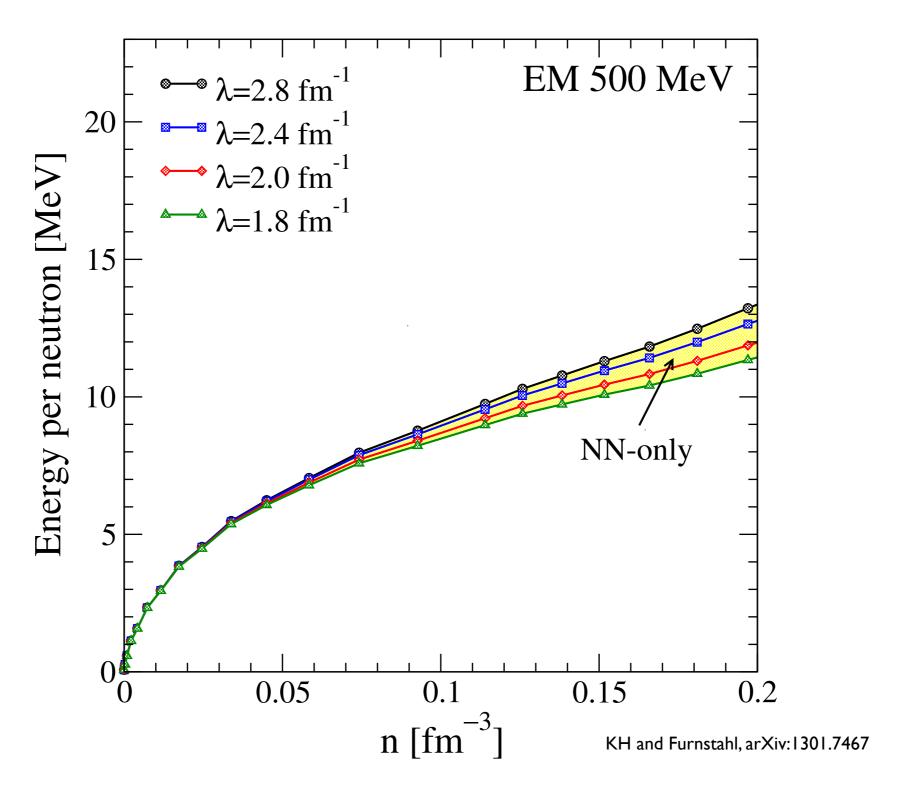


- remarkably reduced scheme dependence for typical momenta $\sim 1\,{\rm fm}^{-1}$, matrix elements with significant phase space well constrained at low resolution
- new momentum structures induced at low resolution
- ullet study based on $\mathrm{N^2LO}$ chiral interactions, improved universality at $\mathrm{N^3LO}$?

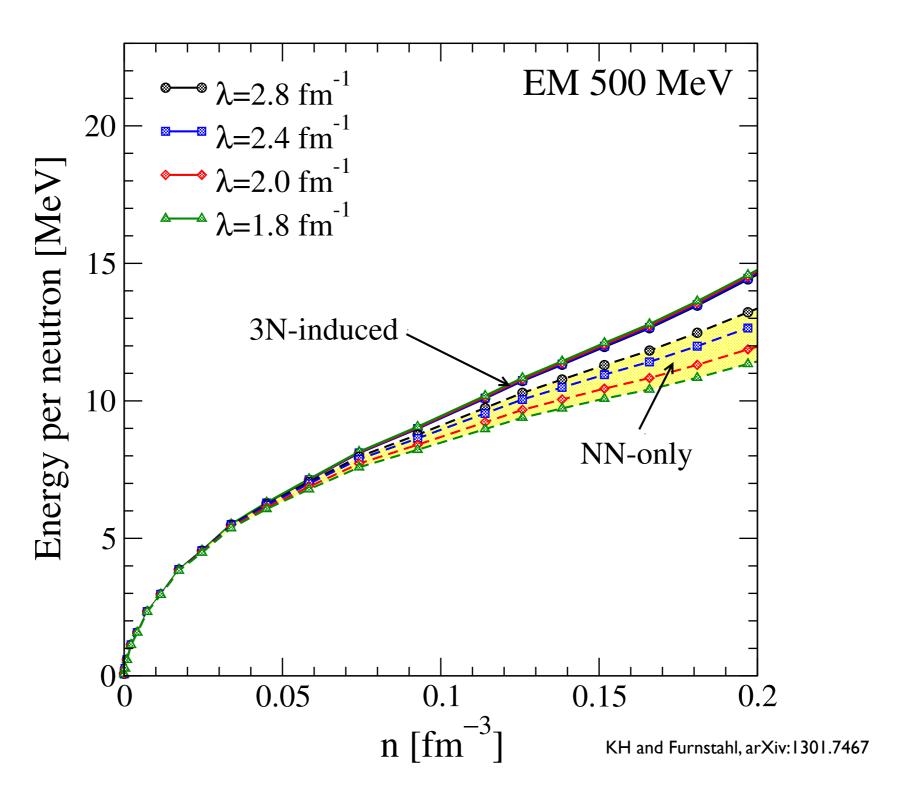
First application to neutron matter: Equation of state



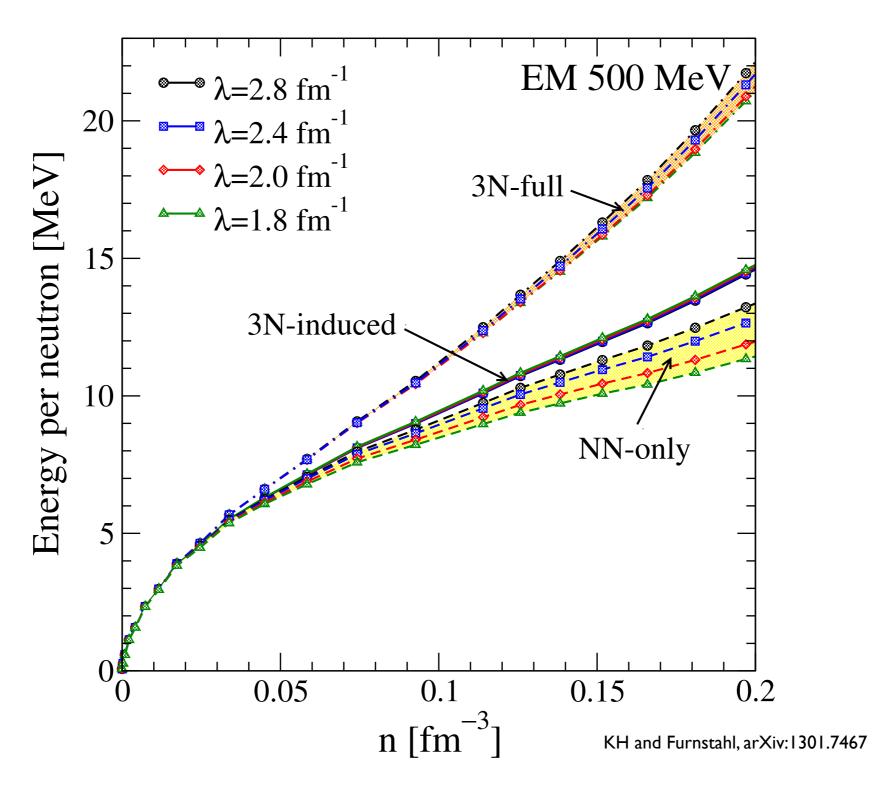
- ullet evolve consistently NN + 3NF in the isospin $\mathcal{T}=3/2$ channel
- calculate EOS by taking all blue-boxed contributions into account
- in this approximation NN and 3NF contributions factorize



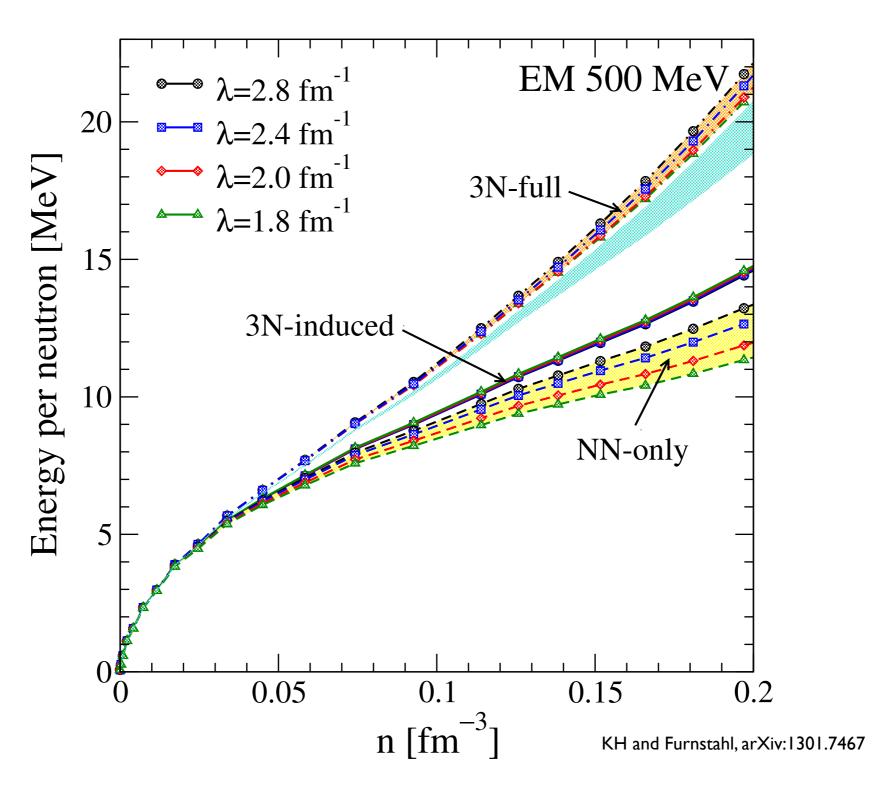
- ullet all partial waves included up to $\mathcal{J}=9/2$ in SRG evolution and EOS calculation
- \bullet consistent 3NF with $c_1=-0.81~{
 m GeV}^{-1}~{
 m and}~c_3=-3.2~{
 m GeV}^{-1}$



- \bullet all partial waves included up to $\mathcal{J}=9/2\,$ in SRG evolution and EOS calculation
- \bullet consistent 3NF with $c_1=-0.81~{
 m GeV}^{-1}~{
 m and}~c_3=-3.2~{
 m GeV}^{-1}$

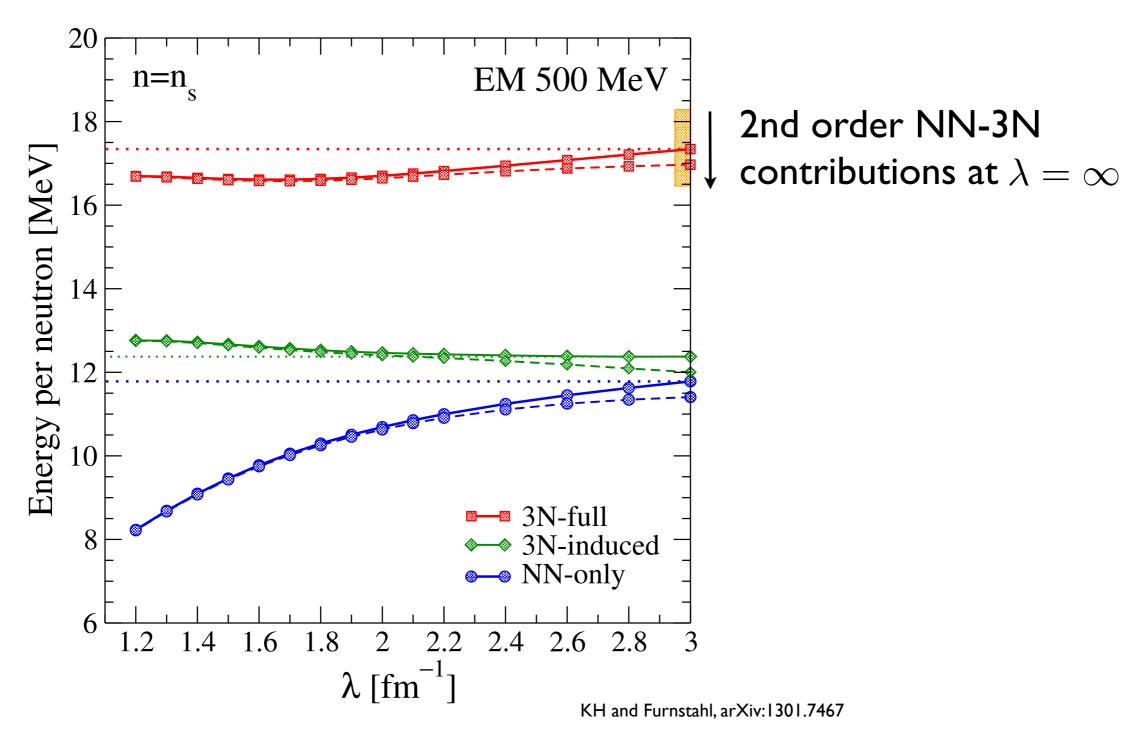


- ullet all partial waves included up to $\mathcal{J}=9/2$ in SRG evolution and EOS calculation
- \bullet consistent 3NF with $c_1=-0.81~{
 m GeV}^{-1}$ and $c_3=-3.2~{
 m GeV}^{-1}$



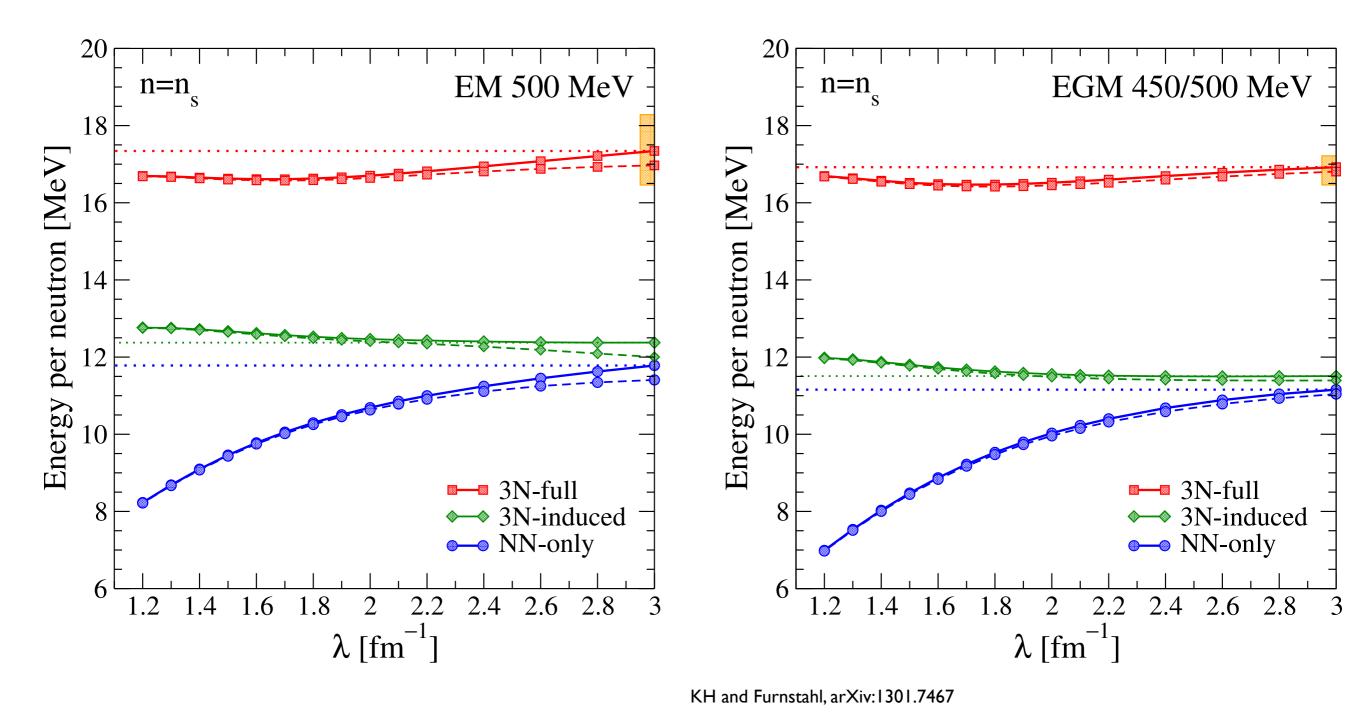
- ullet all partial waves included up to $\mathcal{J}=9/2$ in SRG evolution and EOS calculation
- \bullet consistent 3NF with $c_1=-0.81~{
 m GeV}^{-1}$ and $c_3=-3.2~{
 m GeV}^{-1}$

Resolution-scale dependence at saturation density



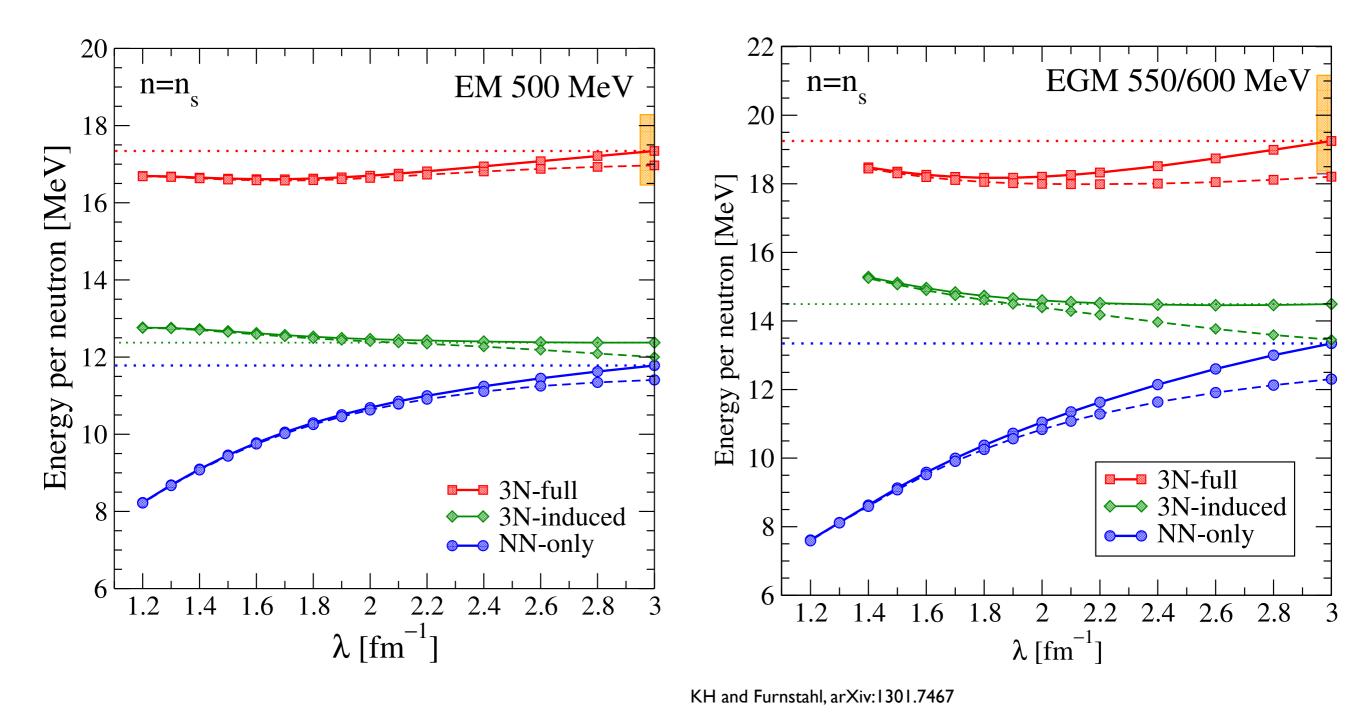
- solid lines: NN resummed, dashed lines: 2nd order
- variations: NN-only 3.6 MeV, 3N-induced: 390 keV, 3N-full: 650 keV
- indications for 4N forces at small λ ?

Resolution-scale dependence at saturation density

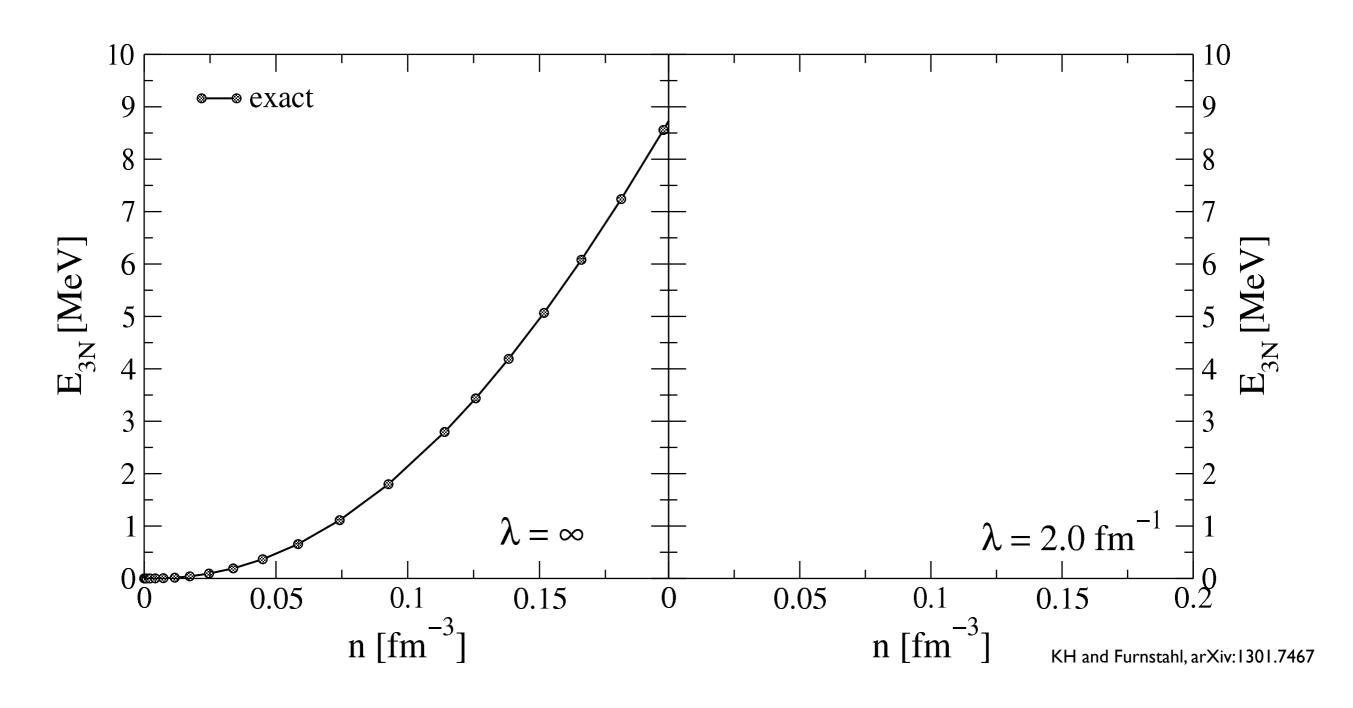


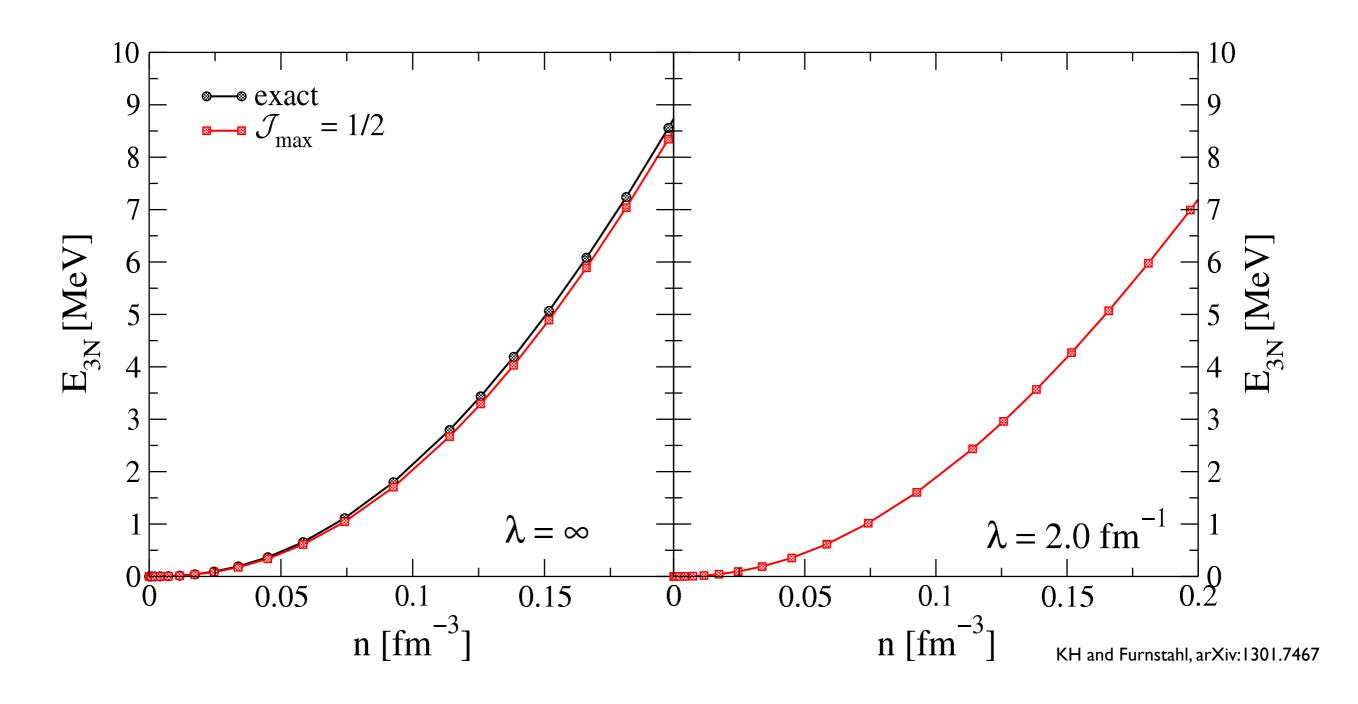
- solid lines: NN resummed, dashed lines: NN 2nd order
- indications for 4N forces at small λ ?

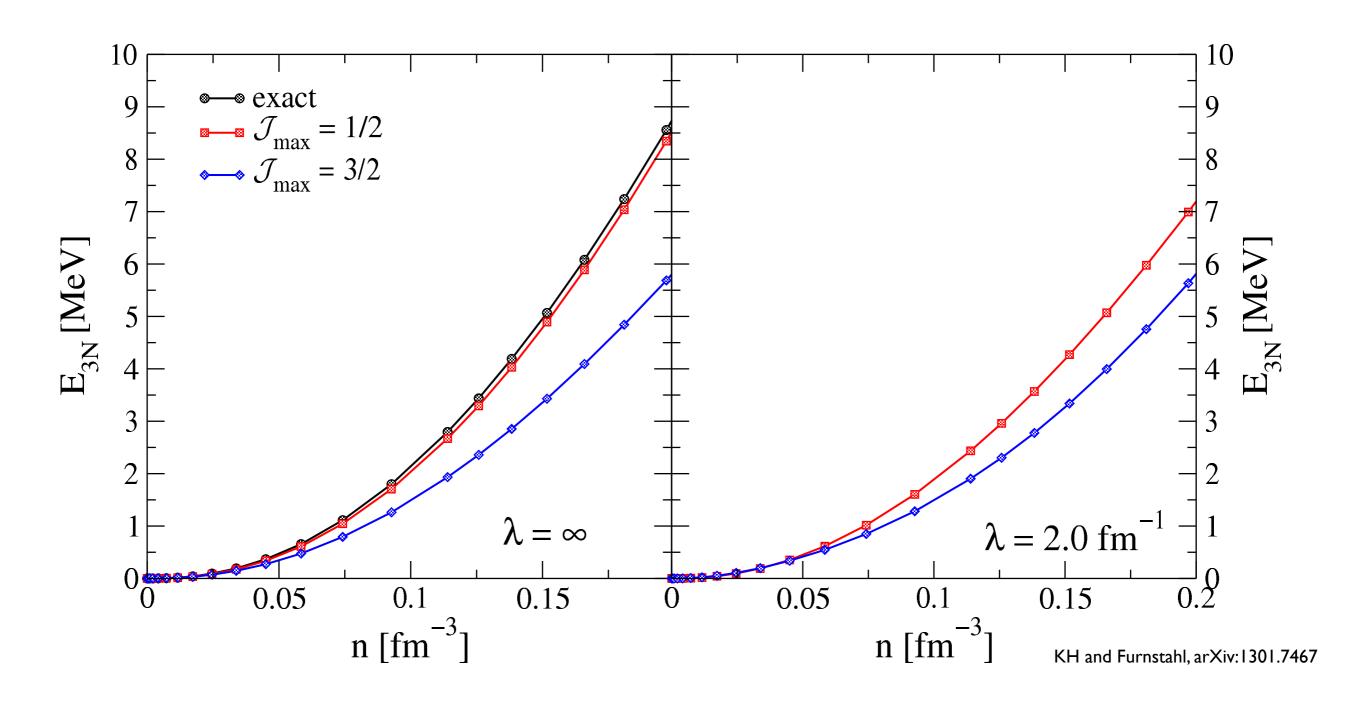
Resolution-scale dependence at saturation density

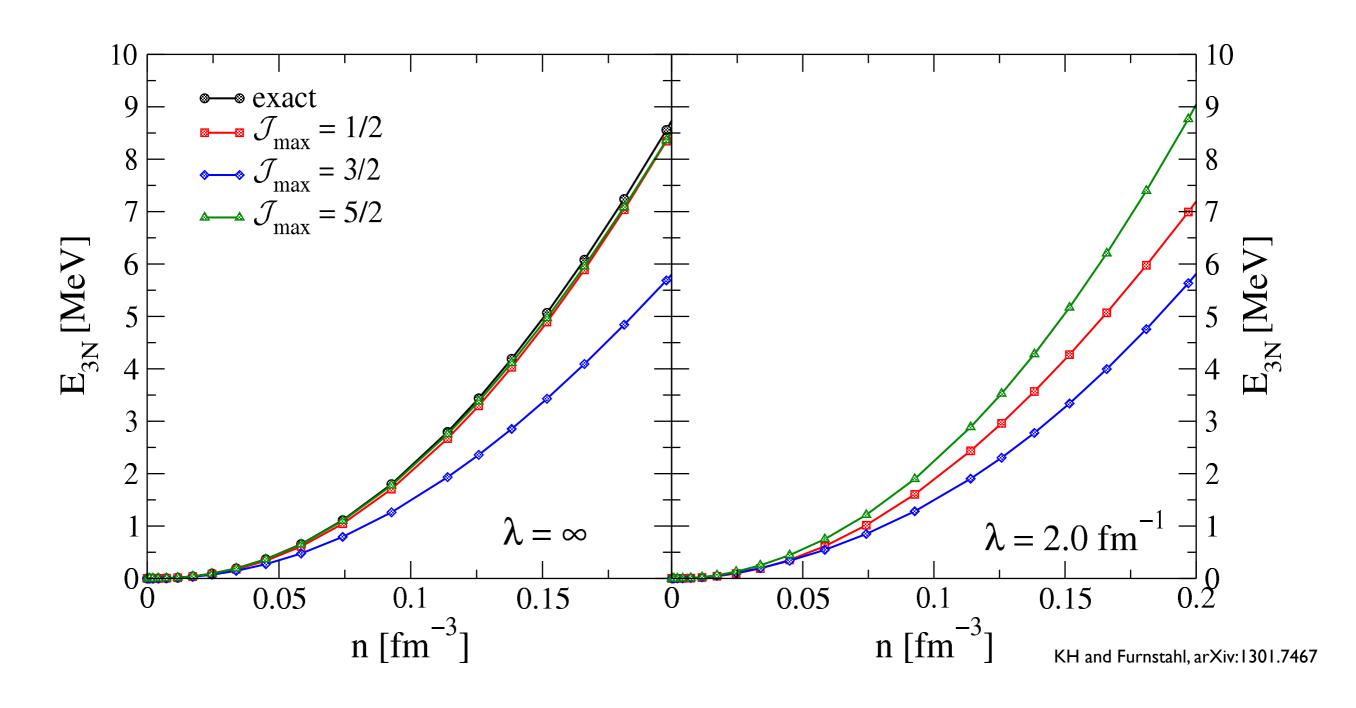


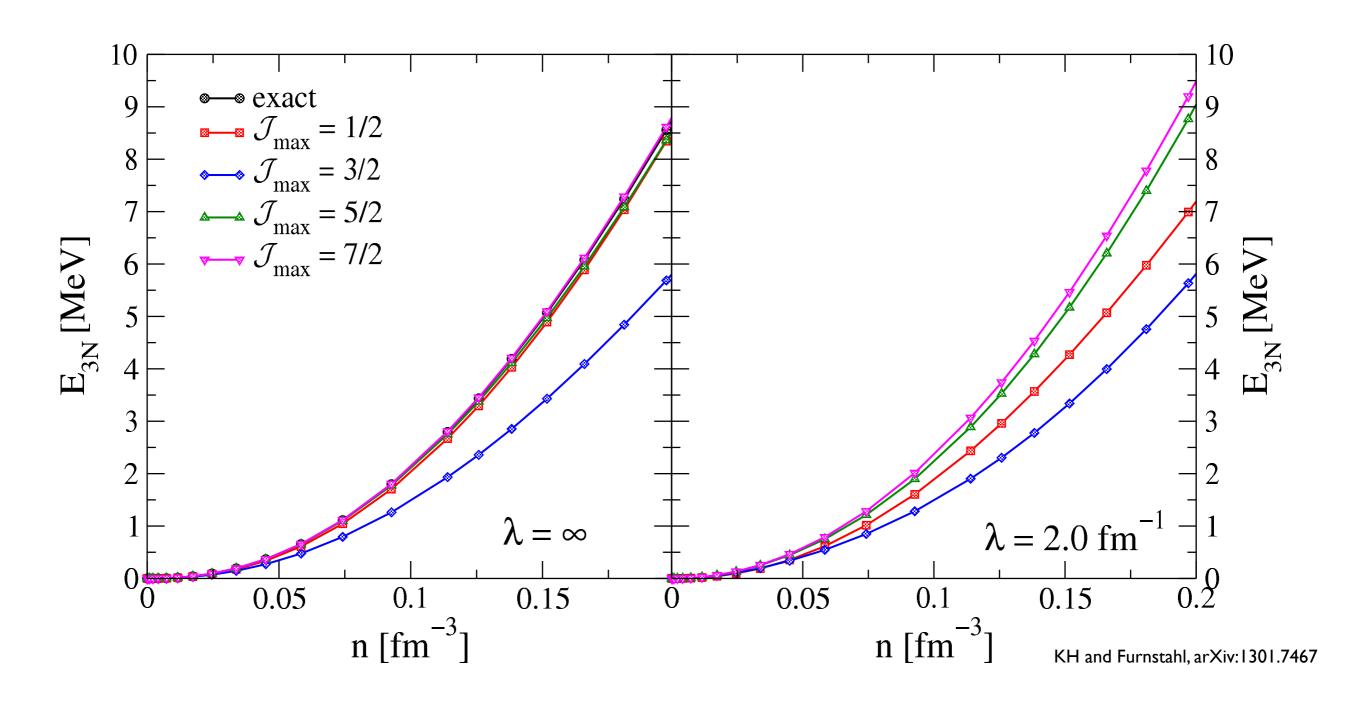
- solid lines: NN resummed, dashed lines: NN 2nd order
- indications for 4N forces at small λ ?

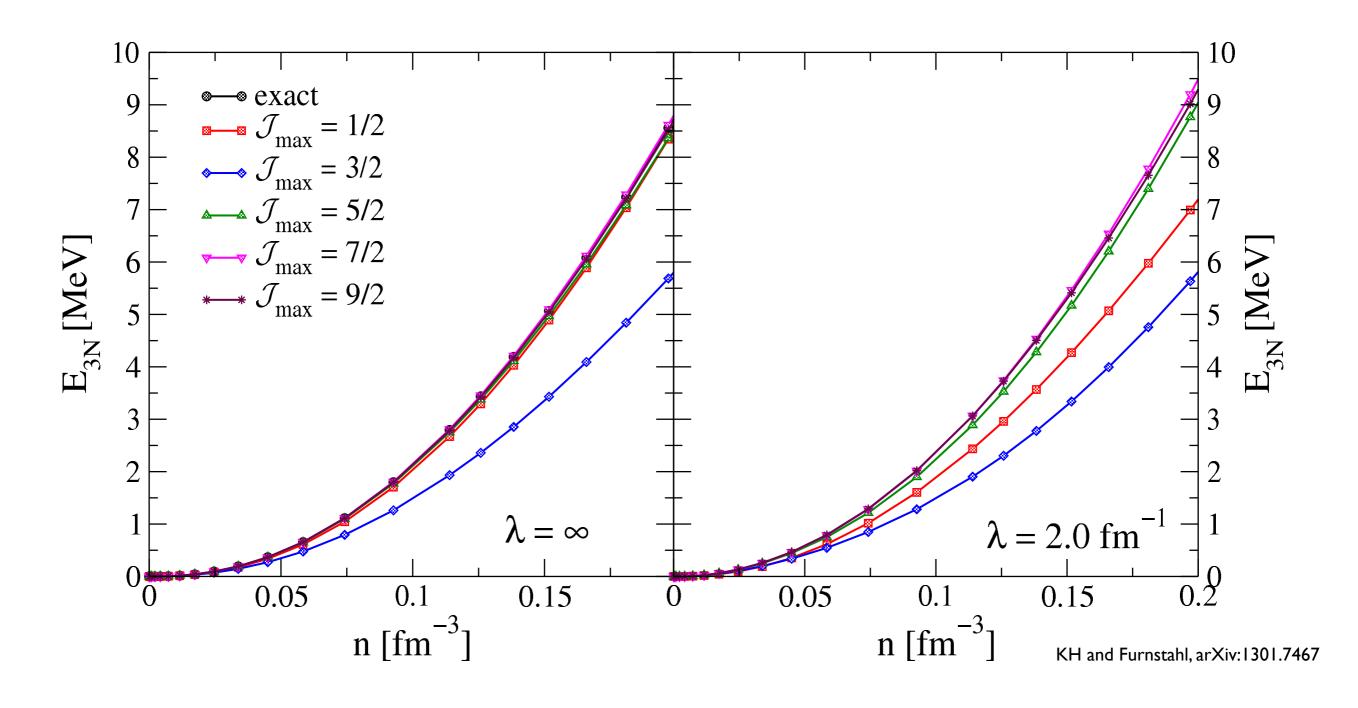


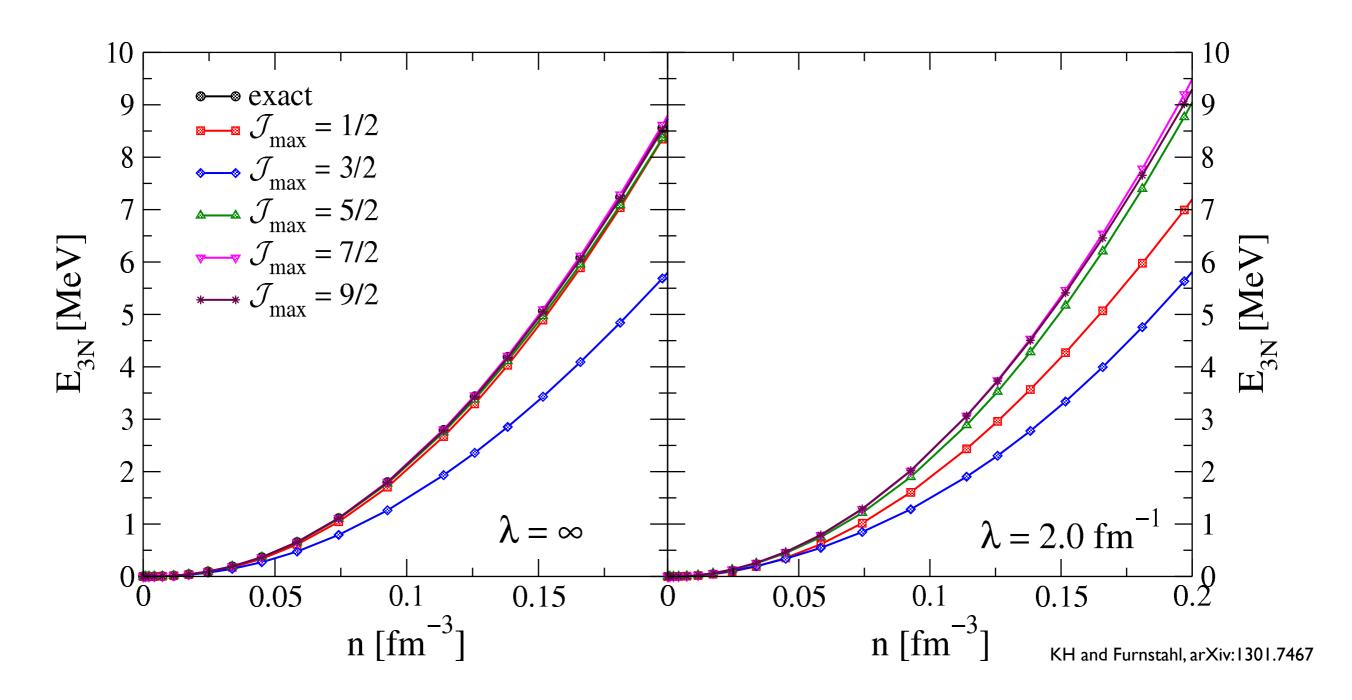






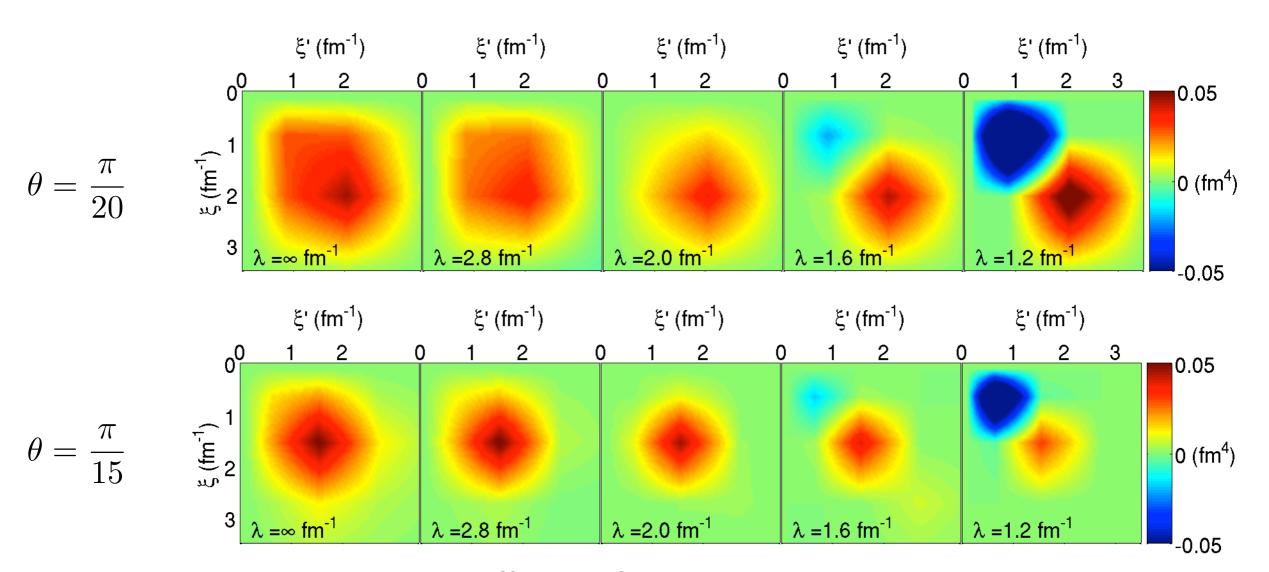






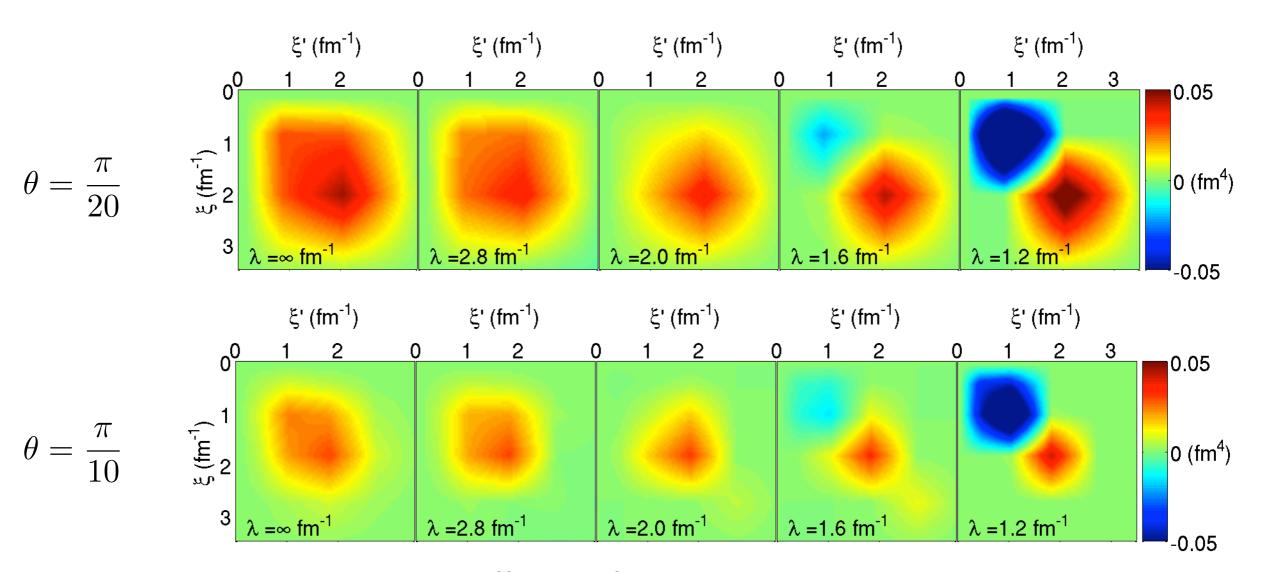
- E_{3N} agrees within 0.4 % with the exact result at saturation density
- E_{3N} converged in partial waves at both scales, $\lambda = \infty$ and $\lambda = 2.0~{\rm fm}^{-1}$

$$\xi^2 = p^2 + \frac{3}{4}q^2$$
 $\tan \theta = \frac{2p}{\sqrt{3}q}$



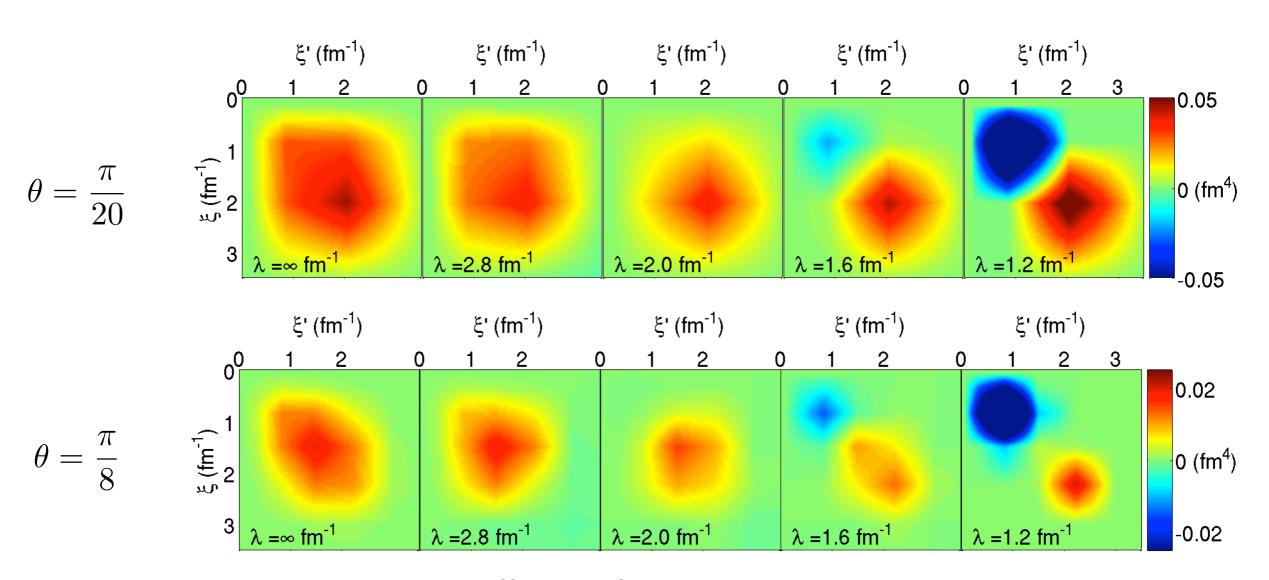
- strong renormalization effects of long-range two-pion exchange
- moderate effects in range $\lambda = \infty$ to $\lambda = 2.0 \, \mathrm{fm}^{-1}$

$$\xi^2 = p^2 + \frac{3}{4}q^2$$
 $\tan \theta = \frac{2p}{\sqrt{3}q}$



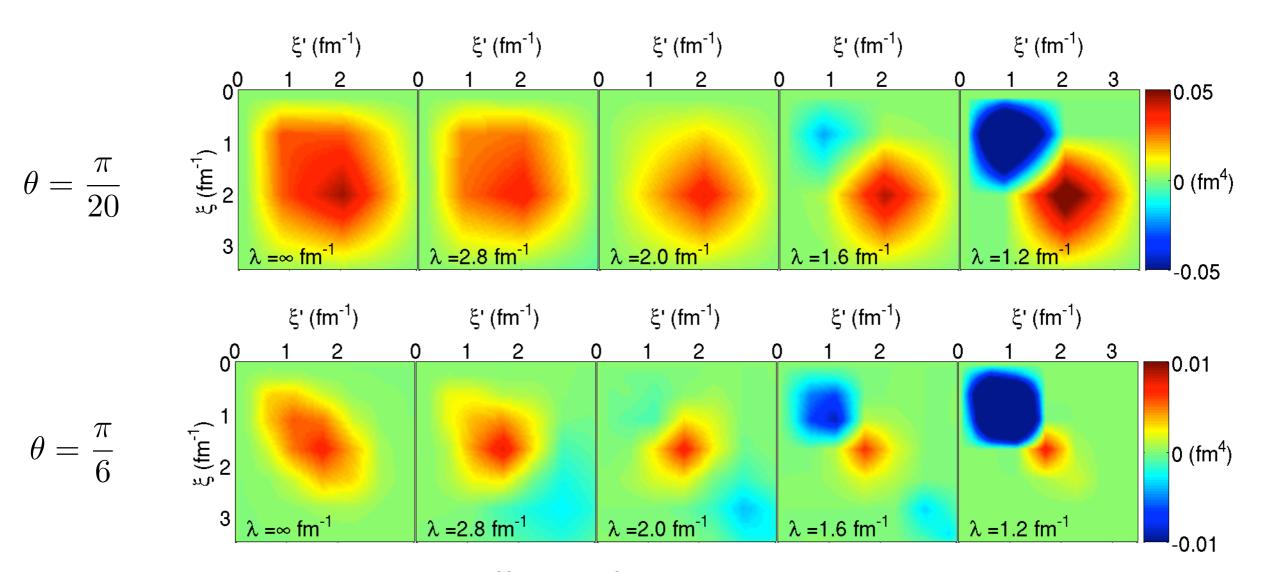
- strong renormalization effects of long-range two-pion exchange
- moderate effects in range $\lambda=\infty$ to $\lambda=2.0~{\rm fm}^{-1}$

$$\xi^2 = p^2 + \frac{3}{4}q^2$$
 $\tan \theta = \frac{2p}{\sqrt{3}q}$



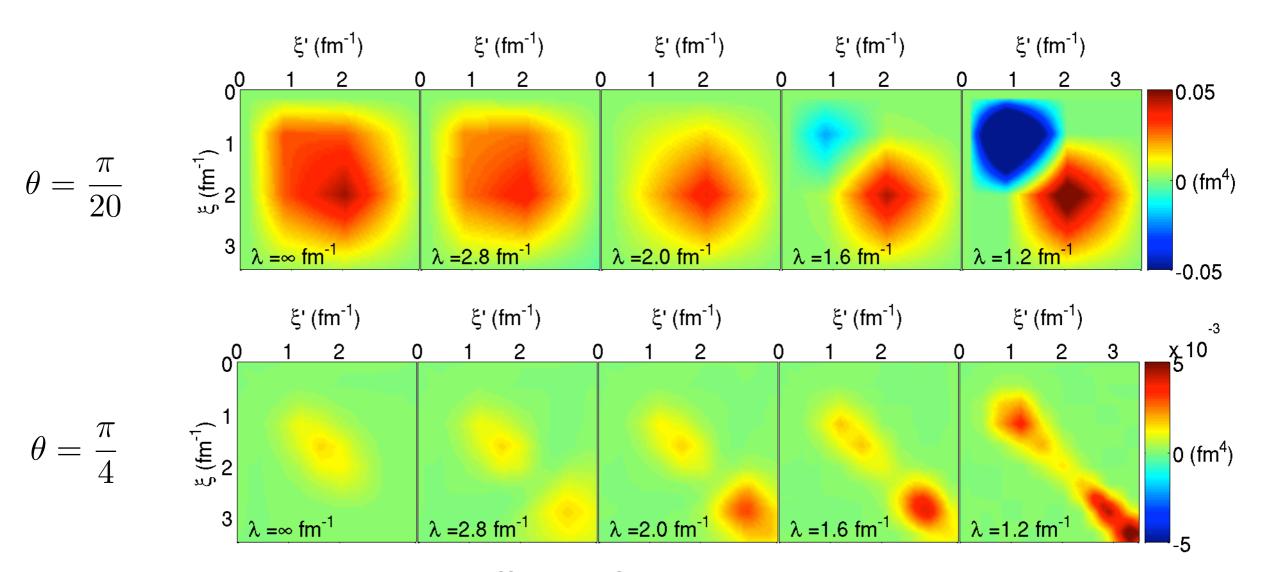
- strong renormalization effects of long-range two-pion exchange
- moderate effects in range $\lambda=\infty$ to $\lambda=2.0~{\rm fm}^{-1}$

$$\xi^2 = p^2 + \frac{3}{4}q^2$$
 $\tan \theta = \frac{2p}{\sqrt{3}q}$



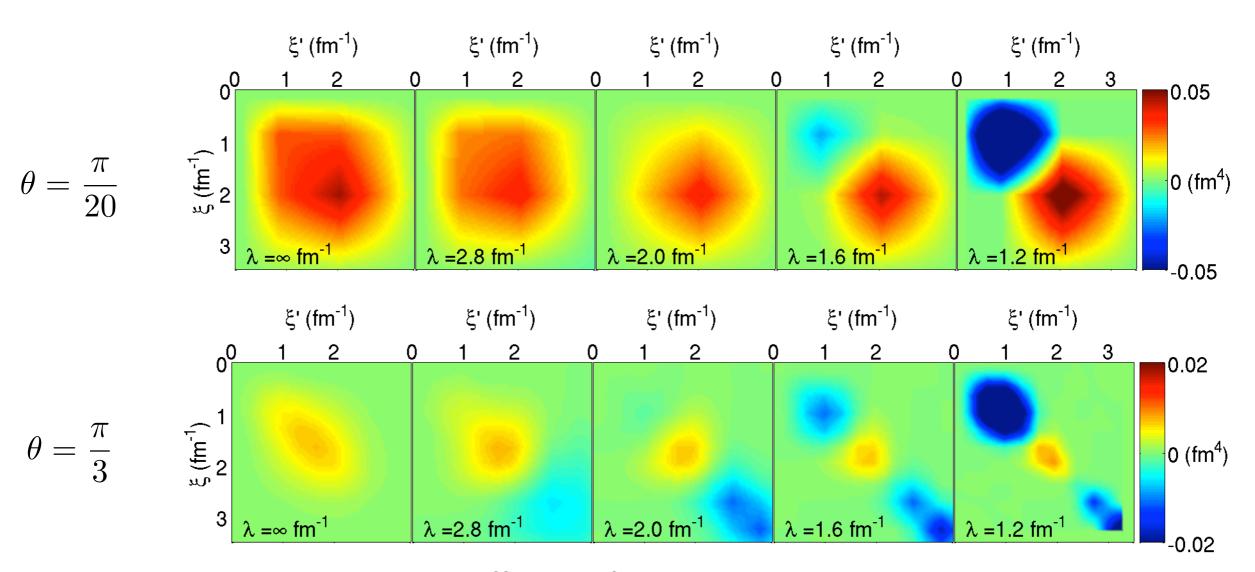
- strong renormalization effects of long-range two-pion exchange
- moderate effects in range $\lambda=\infty$ to $\lambda=2.0~{\rm fm}^{-1}$

$$\xi^2 = p^2 + \frac{3}{4}q^2$$
 $\tan \theta = \frac{2p}{\sqrt{3}q}$



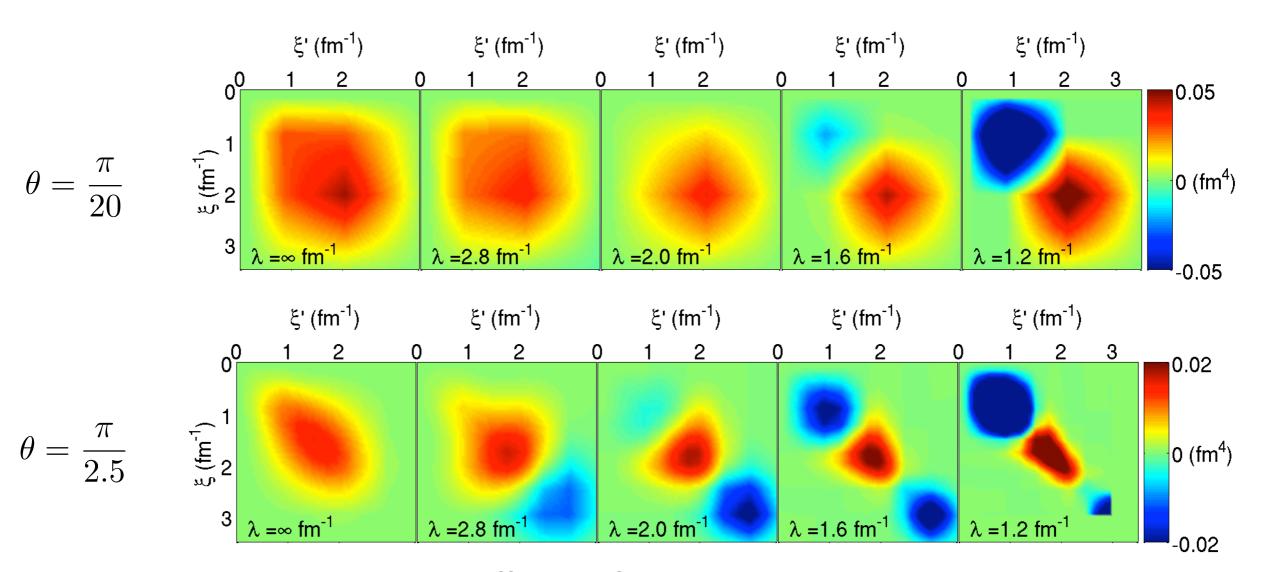
- strong renormalization effects of long-range two-pion exchange
- moderate effects in range $\lambda=\infty$ to $\lambda=2.0~{\rm fm}^{-1}$

$$\xi^2 = p^2 + \frac{3}{4}q^2$$
 $\tan \theta = \frac{2p}{\sqrt{3}q}$



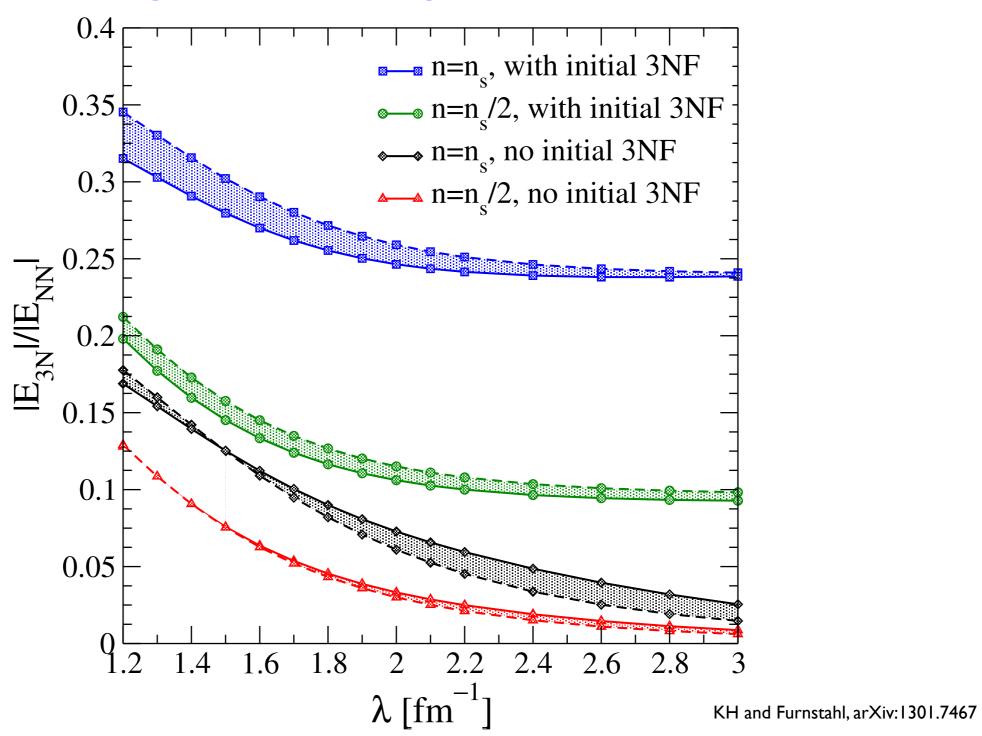
- strong renormalization effects of long-range two-pion exchange
- moderate effects in range $\lambda=\infty$ to $\lambda=2.0~{\rm fm}^{-1}$

$$\xi^2 = p^2 + \frac{3}{4}q^2$$
 $\tan \theta = \frac{2p}{\sqrt{3}q}$



- strong renormalization effects of long-range two-pion exchange
- moderate effects in range $\lambda=\infty$ to $\lambda=2.0~{\rm fm}^{-1}$

Scaling of three-body contributions



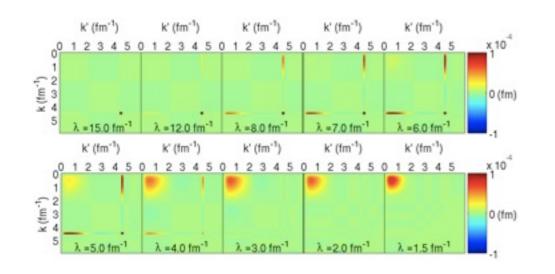
- ullet relative size of 3N contribution grows systematically towards smaller λ
- no obvious trend with density (may be obscured by cancellations among contributions)

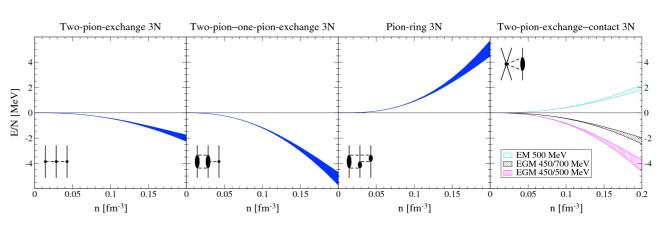
Summary

- demonstrated the feasibility of SRG evolution of NN+3NF in momentum space
- first results of neutron matter based on consistently evolved NN+3NF interactions
- strong renormalization effects of chiral two-pion exch. interaction in neutron matter
- no indications of significant contributions from 4N forces down to $\lambda=1.2~{\rm fm}^{-1}$ in neutron matter

Outlook

- inclusion of 3NF N3LO contributions in RG evolution
- ullet extend RG evolution to $\mathcal{T}=1/2$ channels, application to nuclear matter
- transformation to HO basis, application to finite nuclei (CC, NCSM)
- RG evolution of operators: nuclear scaling and correlations in nuclear systems



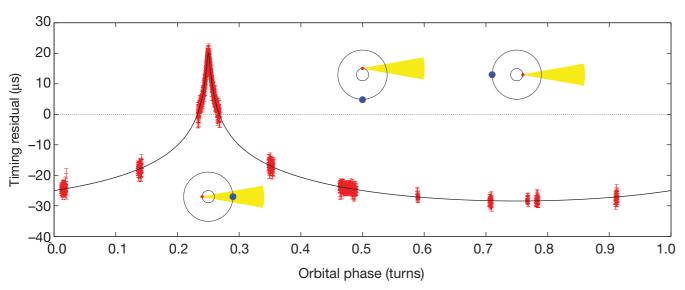


Constraints on the nuclear equation of state (EOS)

nature

A two-solar-mass neutron star measured using Shapiro delay

P. B. Demorest¹, T. Pennucci², S. M. Ransom¹, M. S. E. Roberts³ & J. W. T. Hessels^{4,5}

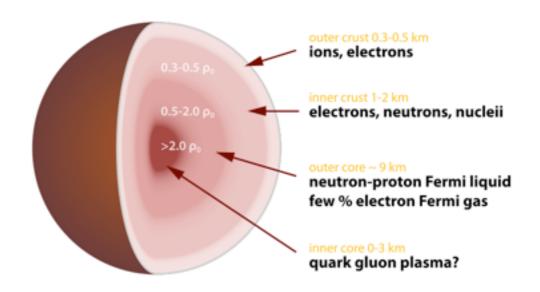


Demorest et al., Nature 467, 1081 (2010)

$$M_{\rm max} = 1.65 M_{\odot} \rightarrow 1.97 \pm 0.04 M_{\odot}$$

Calculation of neutron star properties requires EOS up to high densities.





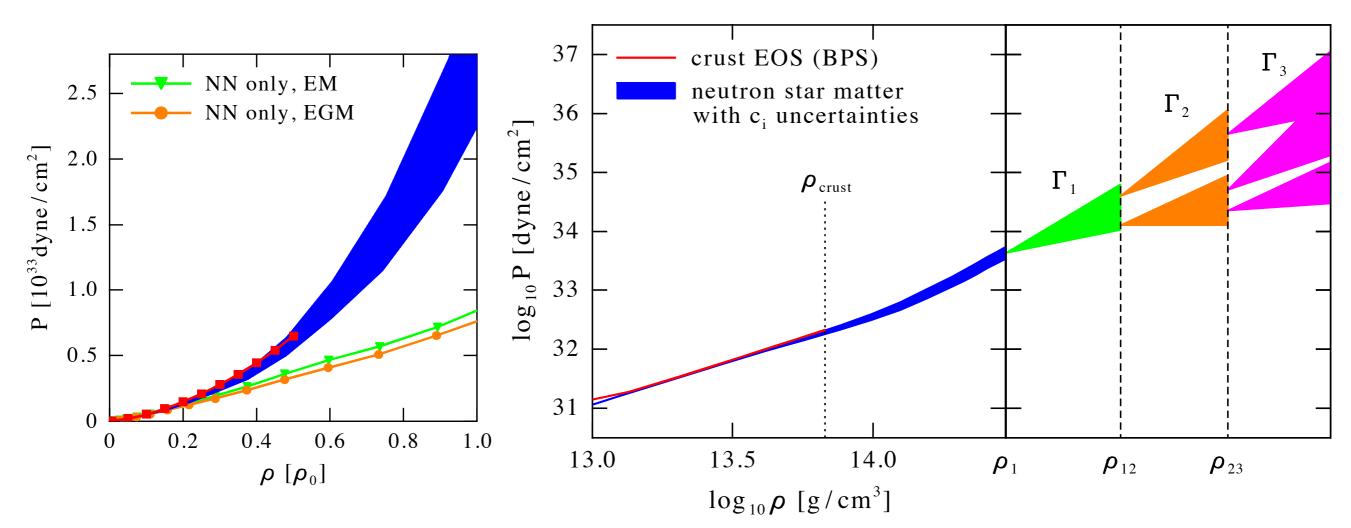
Strategy:

Use observations to constrain the high-density part of the nuclear EOS.

Neutron star radius constraints

incorporation of beta-equilibrium: neutron matter → neutron star matter parametrize piecewise high-density extensions of EOS:

- ullet use polytropic ansatz $\,p\sim
 ho^{\Gamma}$
- ullet range of parameters $\Gamma_1,
 ho_{12}, \Gamma_2,
 ho_{23}, \Gamma_3$ limited by physics!



KH, Lattimer, Pethick, Schwenk, in preparation KH, Lattimer, Pethick, Schwenk, PRL 105, 161102 (2010)

Constraints on the nuclear equation of state

use the constraints:

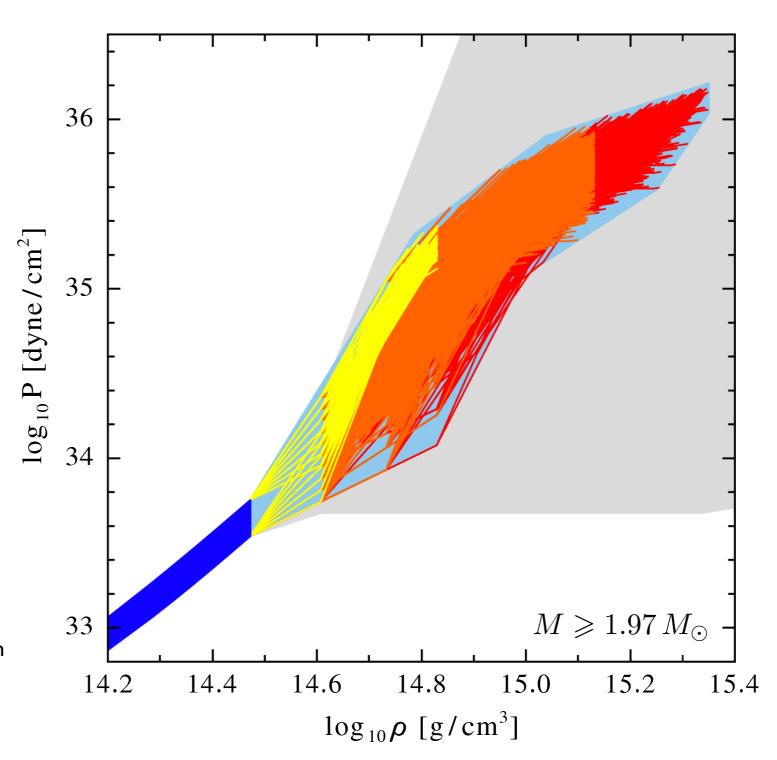
recent NS observation

$$M_{\rm max} > 1.97 \, M_{\odot}$$

causality

$$v_s(\rho) = \sqrt{dP/d\varepsilon} < c$$

KH, Lattimer, Pethick, Schwenk, in preparation



significant reduction of uncertainty band

Constraints on the nuclear equation of state



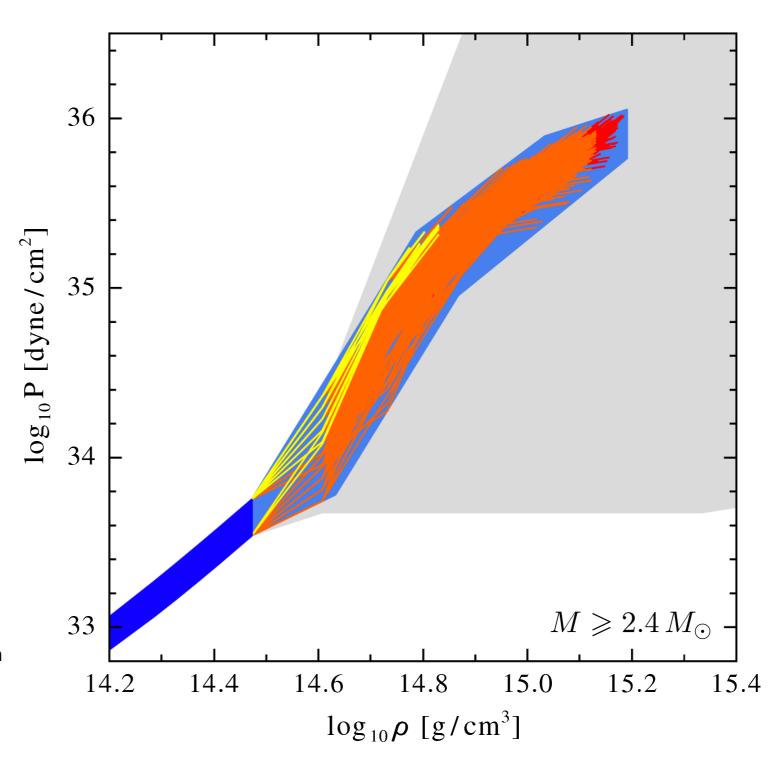
NS mass

$$M_{\rm max} > 2.4 \, M_{\odot}$$

causality

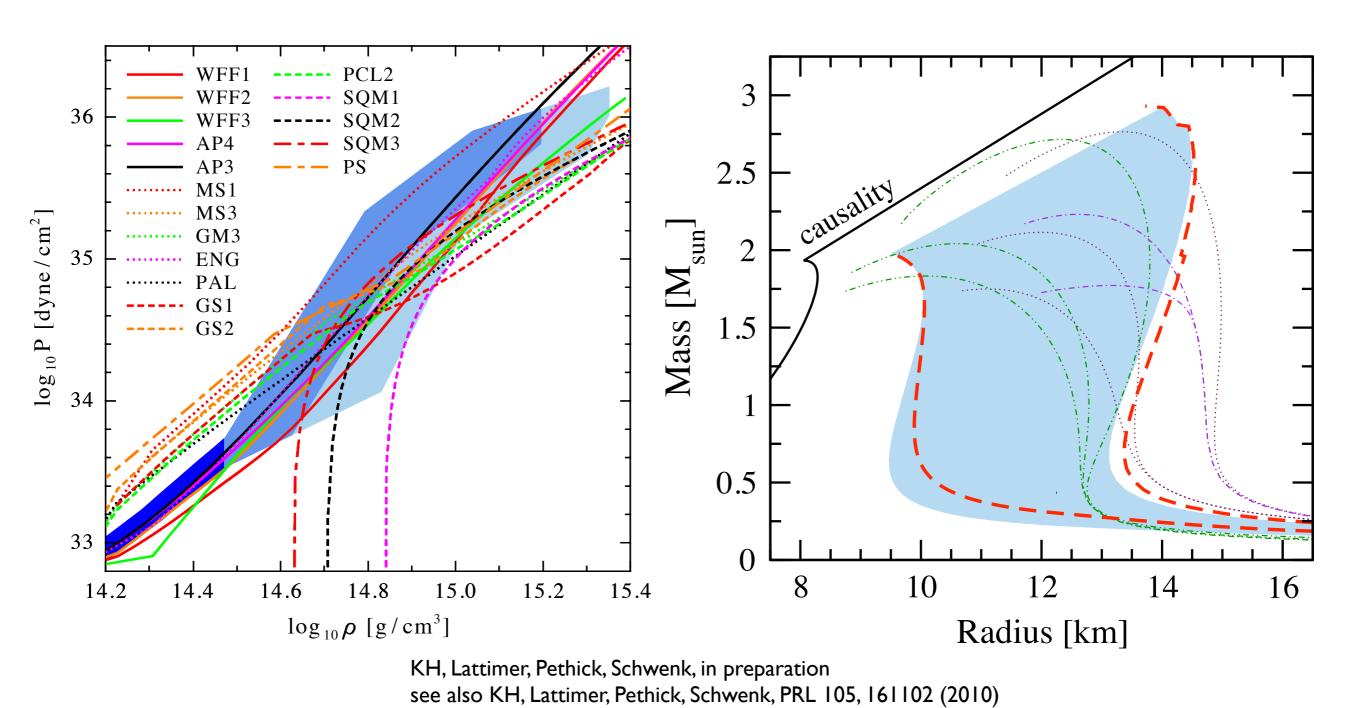
$$v_s(\rho) = \sqrt{dP/d\varepsilon} < c$$

KH, Lattimer, Pethick, Schwenk, in preparation



increased $M_{
m max}$ systematically reduces width of band

Constraints on neutron star radii



- low-density part of EOS sets scale for allowed high-density extensions
- radius constraint for typical $1.4\,M_\odot$ neutron star: $9.8-13.4\,\mathrm{km}$