



**TRIUMF**

Canada's national laboratory  
for particle and nuclear physics  
and accelerator-based science

# Towards triples inclusion in dipole excitations

Mirko Miorelli | TRIUMF – UBC

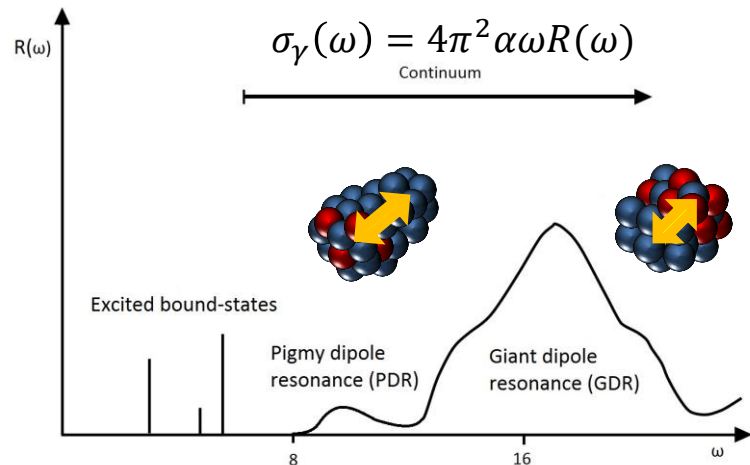
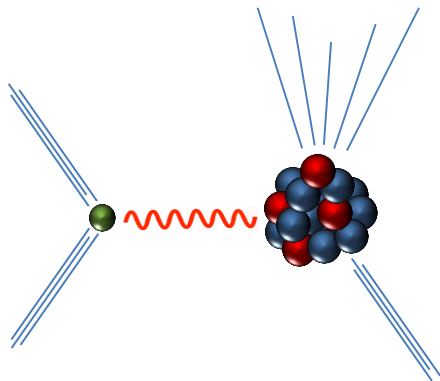
In collaboration with:

S.Bacca, N. Barnea, G. Hagen, G. R. Jansen  
G. Orlandini, T. Papenbrock

March 2, 2016



a place of mind  
THE UNIVERSITY OF  
BRITISH COLUMBIA



Experiments with EM probes allow for a direct connection with theory:

- Small coupling constant (perturbative treatment)
- Transition matrix elements and cross section are directly related

- Focus on the response of the nucleus to an external dipole excitation
- Continuum region: dipole collective modes (PDR & GDR)

- Nuclear matter:

$$E(\rho, \delta) = E(\rho, 0) + S(\rho)\delta^2 + o(\delta^4)$$

$$\rho = \rho_p + \rho_n$$

$$\delta = \frac{\rho_n - \rho_p}{\rho_p + \rho_n}$$

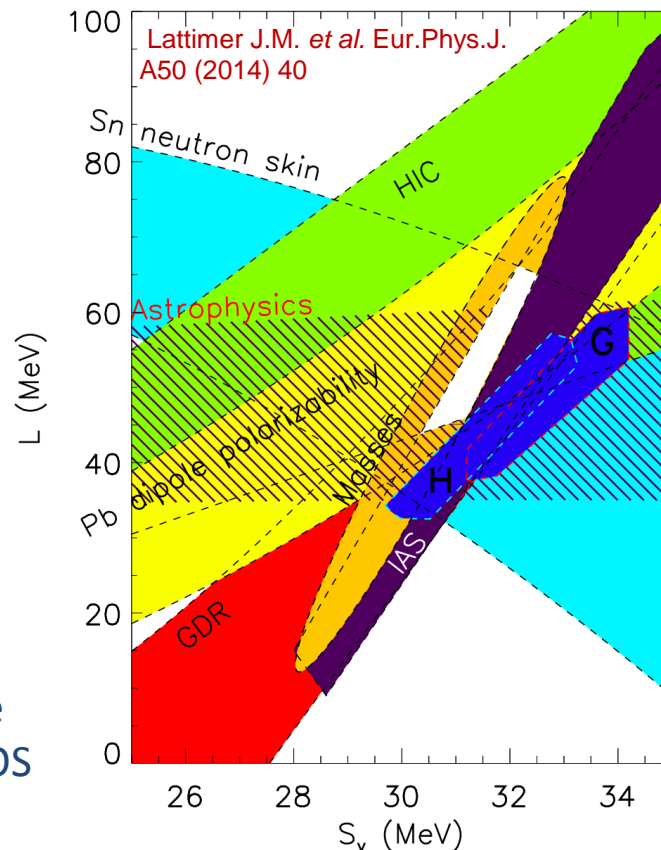
- Expand around saturation density:

$$S(\rho) = S_v + \frac{L}{3\rho_0}(\rho - \rho_0) + \frac{K_{sym}}{18\rho_0^2}(\rho - \rho_0)^2 + \dots$$

- $S_v, L$  and  $K_{sym}$  can be related to finite-nuclei properties



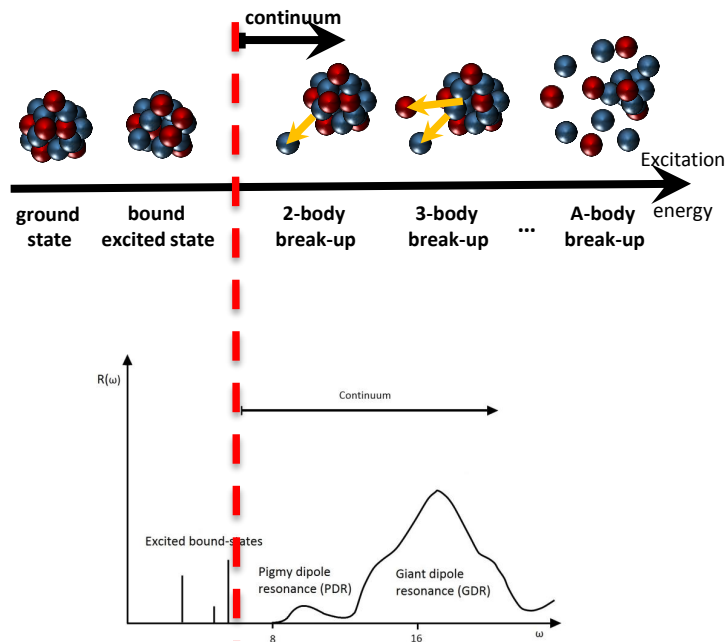
We cannot create neutron stars in a lab, but we can use nuclear observables to constrain the EOS



$$\alpha_D = 2\alpha \int \frac{R(\omega)}{\omega} d\omega$$

$$R(\omega) = \sum_f |\langle \Psi_f | D | \Psi_0 \rangle|^2 \delta(\omega - E_0 - E_f)$$

$|\Psi_f\rangle$  is a **many-body scattering state** and cannot be calculated explicitly  $\rightarrow$  we use **integral transforms**



$$\alpha_D = 2\alpha \int \frac{R(\omega)}{\omega} d\omega$$

$$R(\omega) = \sum_f |\langle \Psi_f | D | \Psi_0 \rangle|^2 \delta(\omega - E_0 - E_f)$$

$|\Psi_f\rangle$  is a **many-body scattering state** and cannot be calculated explicitly  $\rightarrow$  we use **integral transforms**

$$(\bar{H} - E_0 - z) |\tilde{\Psi}_R(z)\rangle = \bar{D} |0_R\rangle$$

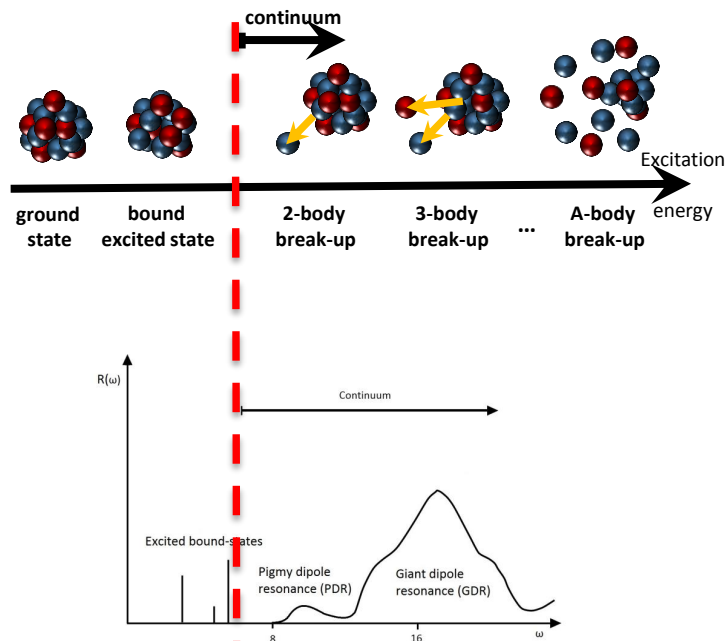
Bacca et al., Phys. Rev. Lett. 111, 122502 (2013)

Bacca S., M.M., et al., Phys. Rev. C 90, 064619 (2014)

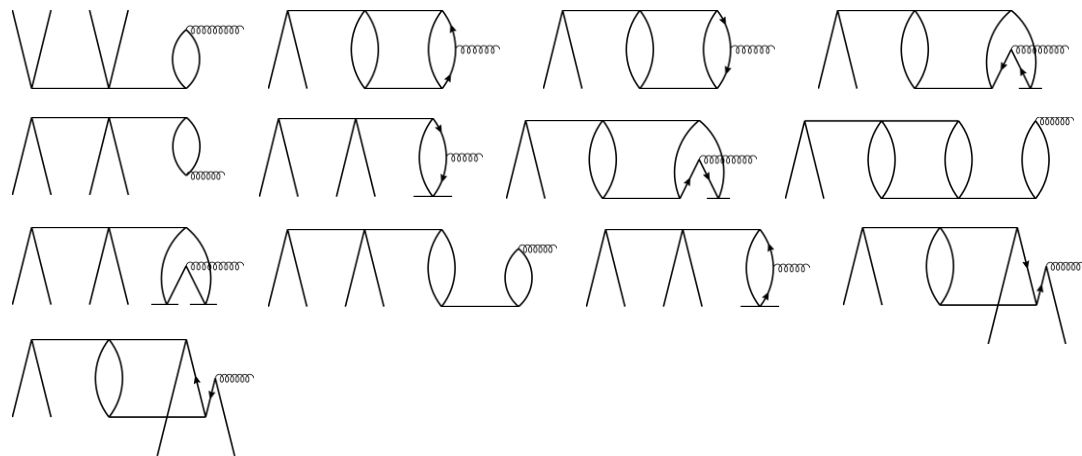
$$|\tilde{\Psi}_R(z)\rangle = [R_0(z) + R_1(z) + R_2(z)] |0_R\rangle$$

$$\left. \begin{aligned} \bar{D} &= e^{-T} D e^T \\ \bar{H} &= e^{-T} H e^T \end{aligned} \right\} T = T_1 + T_2$$

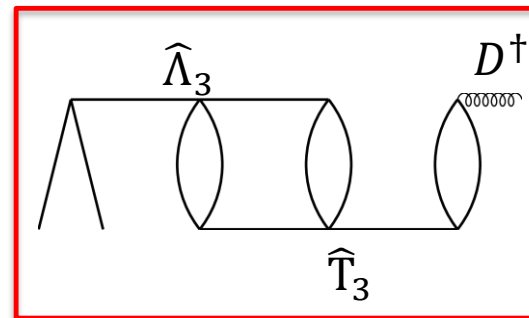
(Equivalently for left states:  $\Lambda = \Lambda_1 + \Lambda_2$  and  $L(z) = L_0(z) + L_1(z) + L_2(z)$ )



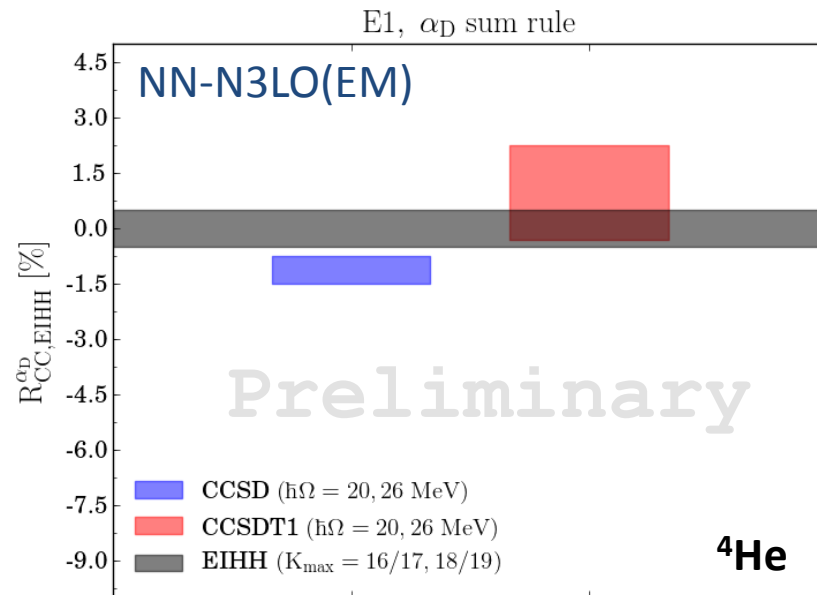
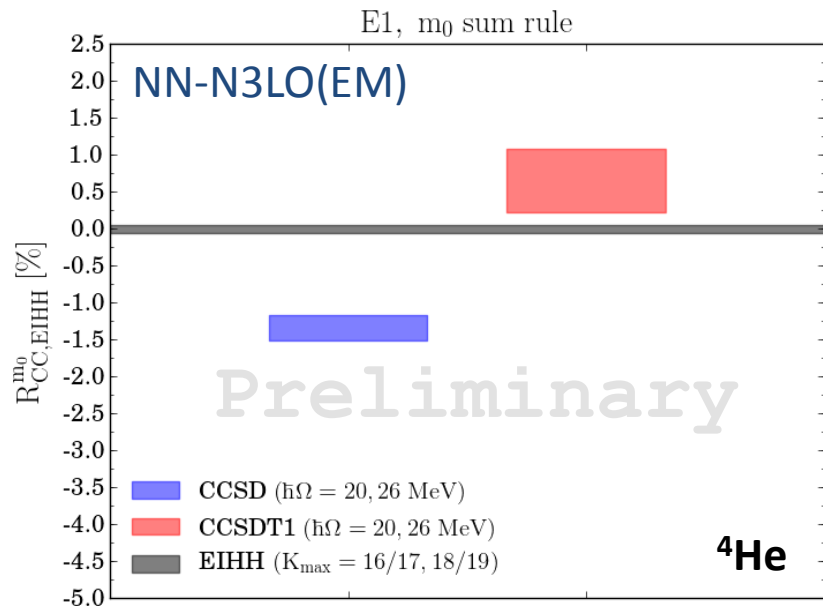
- Linearized triples for the ground state:  $e^{T_1+T_2} + T_3$  and  $\Lambda_1 + \Lambda_2 + \Lambda_3$
- Triples effects on excited states from linear triples ( $R, L$ )



$$m_0 = \langle \phi_0 | (1 + \hat{\Lambda}) \bar{D}^\dagger \bar{D} | \phi \rangle$$

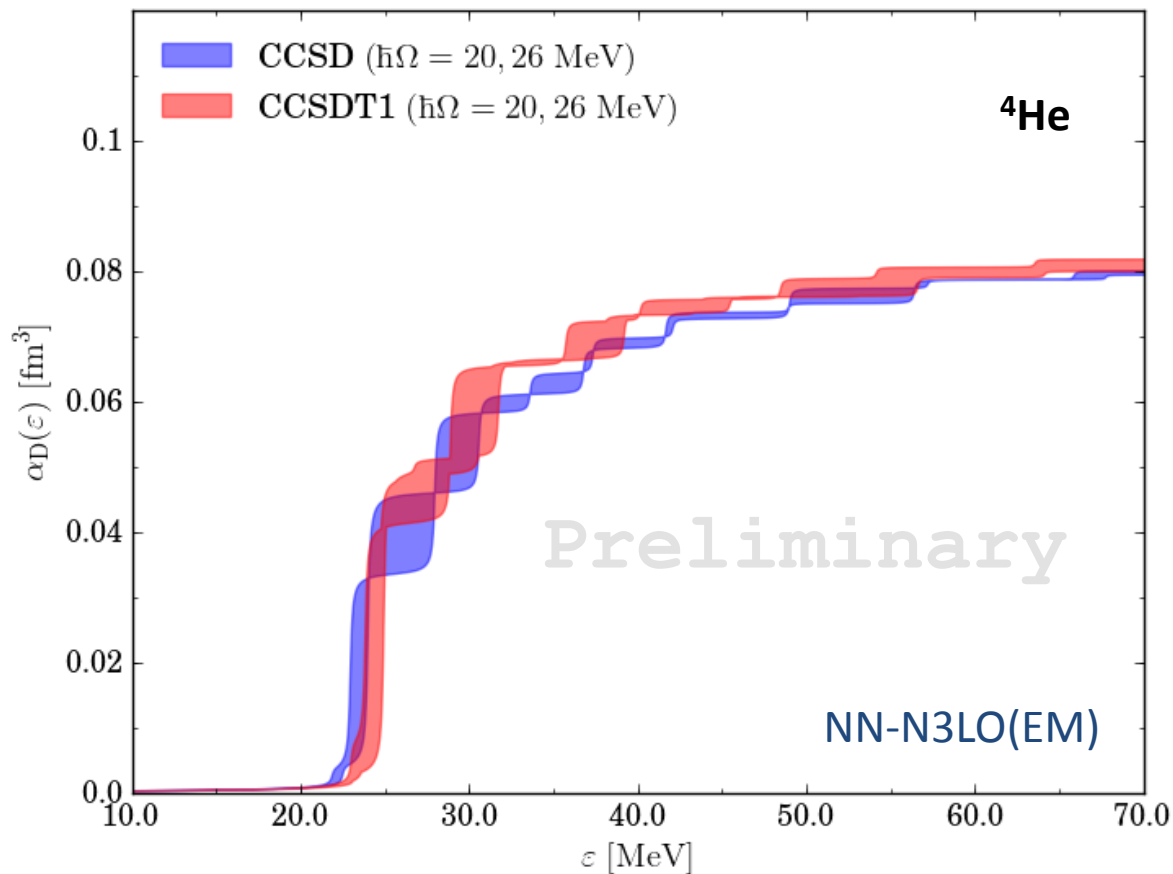


New  $1p1h$  and  $2p2h$  contributions to dipole sum rules from  $\hat{T}_3$  and  $\hat{\Lambda}_3$



- The dipole strength ( $m_0$ ) increases adding triples
- Polarizability also increases
- Coupled cluster results with triples are closer to EIHH

How about the strength distribution?



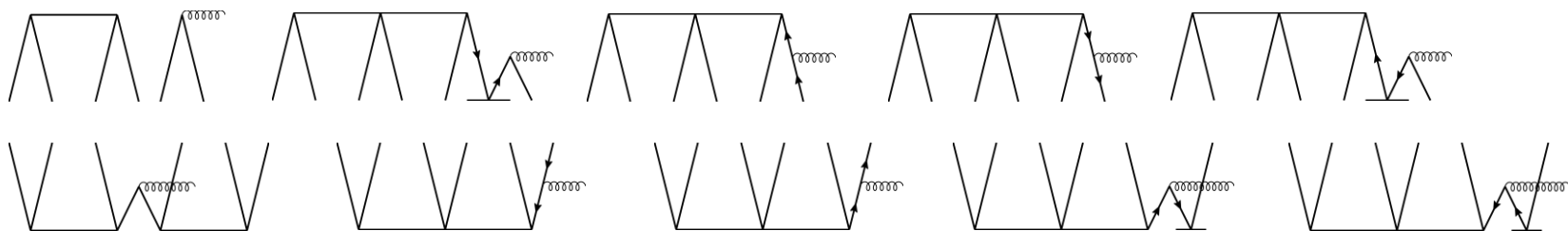
[M.M. et al., Phys. Rev. C 94, 034317 \(2016\)](#)

$$\alpha_D(\epsilon) = 2\alpha \lim_{\Gamma \rightarrow 0} \int^{\epsilon} \frac{L(\sigma, \Gamma)}{\sigma} d\sigma$$

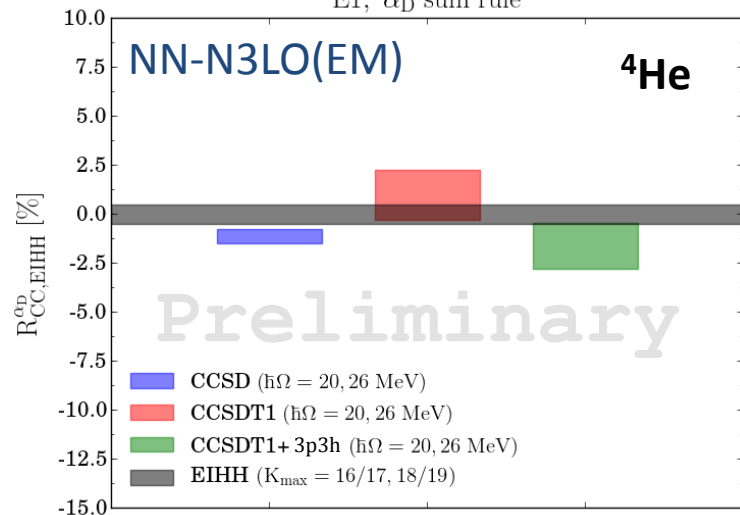
$$\alpha_D(\epsilon) = 2\alpha \int^{\epsilon} \frac{R(\omega)}{\omega} d\omega$$



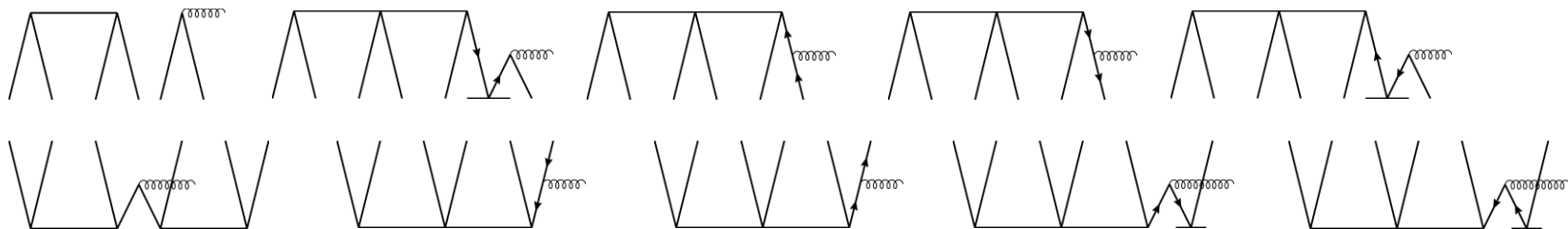
3p3h  
contr.



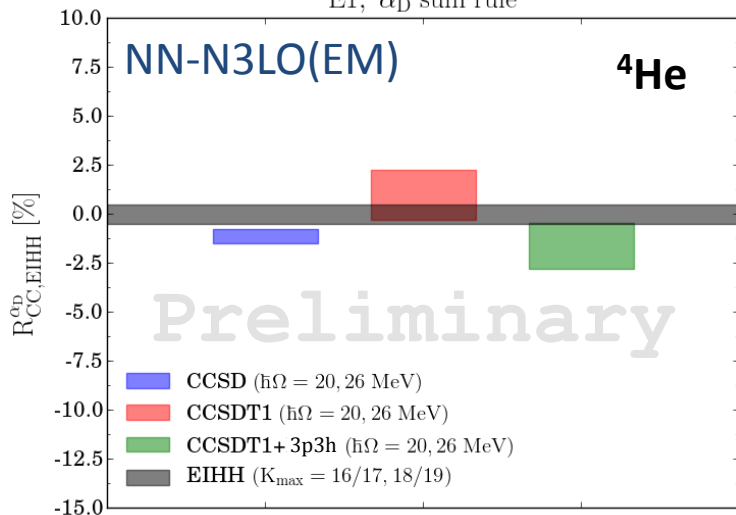
E1,  $\alpha_D$  sum rule



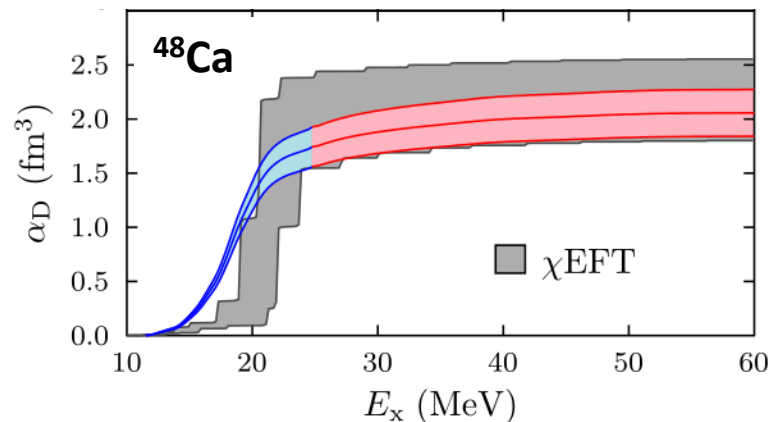
3p3h  
contr.



E1,  $\alpha_D$  sum rule



J. Birkhan, *M.M. et al.*, arXiv:1611.07072 (2016)



- Effect of triples on heavier nuclei?
- Response function with triples?



---

Canada's national laboratory  
for particle and nuclear  
physics  
and accelerator-based science

Thank you!

TRIUMF: Alberta | British Columbia | Calgary |  
Carleton | Guelph | Manitoba | McGill |  
McMaster | Montréal | Northern British  
Columbia | Queen's | Regina | Saint Mary's |  
Simon Fraser | Toronto | Victoria | Western |  
Winnipeg | York

SUMMARY...

... come to see my poster!!



