

# Nonlocal translationally invariant nuclear density

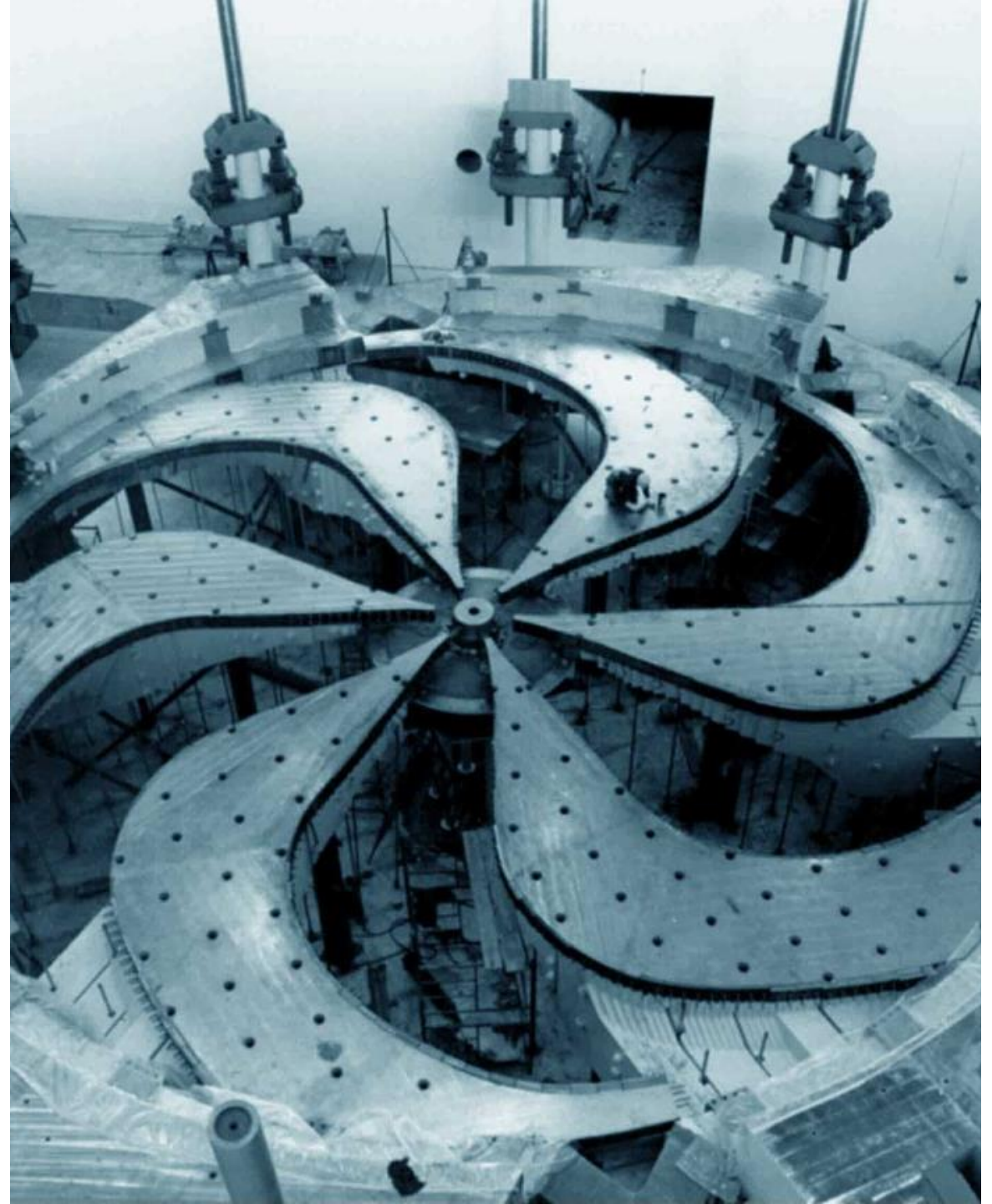
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In collaboration with


Petr Navrátil, Angelo Calci, and  
Matteo Vorabbi

2018-02-28



# No-core shell model (NCSM)

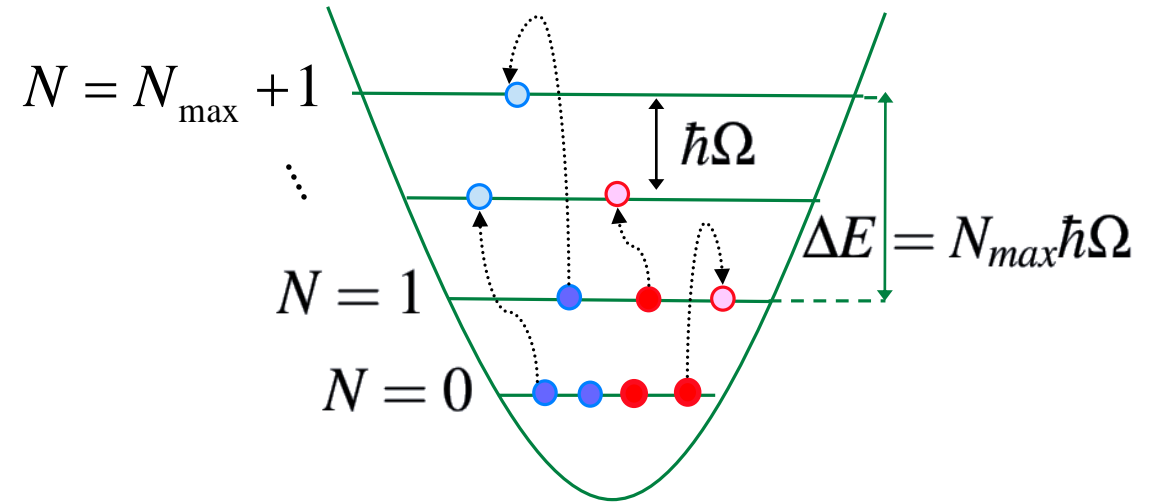
- NCSM is an *ab initio* approach to solve the many-body Schrödinger equation for bound states (narrow resonances) starting from [high-precision NN+NNN interactions](#)
- Uses large (but finite!) expansions in HO many-body basis states



$$\Psi^A = \sum_{N=0}^{N_{max}} \sum_i c_{Ni} \Phi_{Ni}^A$$


- Translational invariance of the internal wave function is preserved when single-particle Slater Determinant (SD) basis is used with  $N_{max}$  truncation

$$\langle \vec{r}_1 \cdots \vec{r}_A \vec{\sigma}_1 \cdots \vec{\sigma}_A \vec{\tau}_1 \cdots \vec{\tau}_A | A\lambda JM \rangle_{SD} = \langle \vec{\xi}_1 \cdots \vec{\xi}_{A-1} \vec{\sigma}_1 \cdots \vec{\sigma}_A \vec{\tau}_1 \cdots \vec{\tau}_A | A\lambda JM \rangle \varphi_{000}(\vec{\xi}_0)$$



# No-core shell model (NCSM)

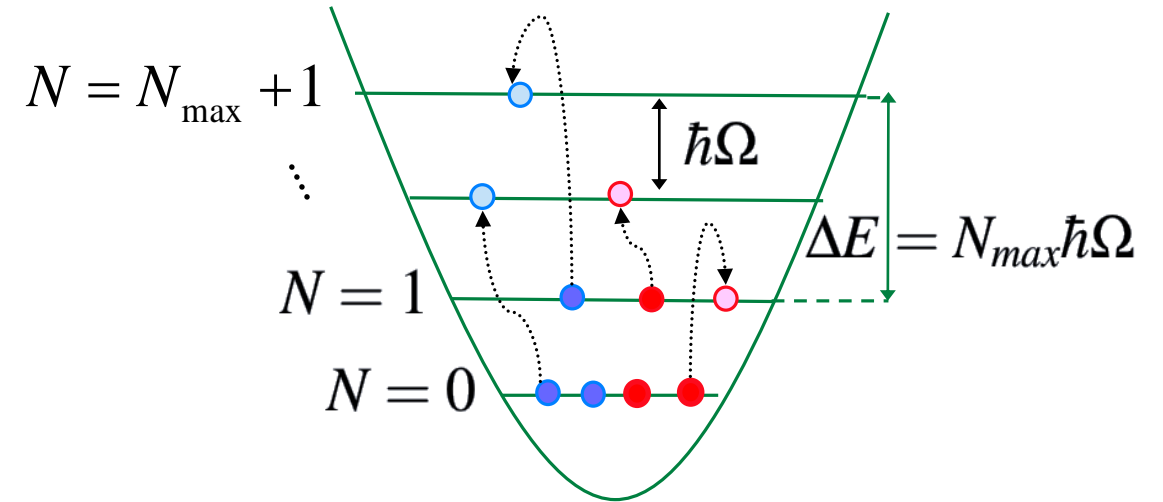
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Intrinsic wavefunction

COM wavefunction

# Coordinate form of the density

## Nonlocal translationally invariant density

- arXiv:1712.02879; Phys. Rev. C, in press.

$$\begin{aligned}
 & \langle A\lambda_f J_f M_f | \rho_{op}(\vec{r} - \vec{R}, \vec{r}' - \vec{R}) | A\lambda_i J_i M_i \rangle \\
 &= \left( \frac{A}{A-1} \right)^{\frac{3}{2}} \sum_{\hat{J}_f} \frac{1}{\hat{J}_f} (J_i M_i K k | J_f M_f) \left( Y_l^* \left( \widehat{\vec{r} - \vec{R}} \right) Y_{l'}^* \left( \widehat{\vec{r}' - \vec{R}} \right) \right)_k^{(K)} \\
 & \times R_{n,l} \left( \sqrt{\frac{A}{A-1}} |\vec{r} - \vec{R}| \right) R_{n',l'} \left( \sqrt{\frac{A}{A-1}} |\vec{r}' - \vec{R}| \right) \\
 & \times (M^K)_{n,l,n',l',n_1,l_1,n_2,l_2}^{-1} (-1)^{l_1+l_2+K+j_2-\frac{1}{2}} \hat{J}_1 \hat{J}_2 \hat{K} \left\{ \begin{matrix} j_1 & j_2 & K \\ l_2 & l_1 & 1/2 \end{matrix} \right\} \\
 & \times \frac{(-1)}{\hat{K}} {}_{SD} \langle A\lambda_f J_f \parallel (a_{n_1 l_1 j_1}^\dagger \tilde{a}_{n_2 l_2 j_2})^{(K)} \parallel A\lambda_i J_i \rangle_{SD}
 \end{aligned}$$

# Coordinate form of the density

Normalization

5

Nonlocal translationally invariant density

- arXiv:1712.02879; Phys. Rev. C, in press.

$$\int d\vec{x} \langle A\lambda JM | \rho_{op}^{phys}(\vec{x}) | A\lambda JM \rangle = A$$

$$\begin{aligned} & \langle A\lambda_f J_f M_f | \rho_{op}(\vec{r} - \vec{R}, \vec{r}' - \vec{R}) | A\lambda_i J_i M_i \rangle \\ &= \left( \frac{A}{A-1} \right)^{\frac{3}{2}} \sum_{\hat{J}_f} \frac{1}{\hat{J}_f} (J_i M_i K k | J_f M_f) \left( Y_{l'}^* \left( \widehat{\vec{r} - \vec{R}} \right) Y_{l'}^* \left( \widehat{\vec{r}' - \vec{R}} \right) \right)_k^{(K)} \\ & \times R_{n,l} \left( \sqrt{\frac{A}{A-1}} |\vec{r} - \vec{R}| \right) R_{n',l'} \left( \sqrt{\frac{A}{A-1}} |\vec{r}' - \vec{R}| \right) \\ & \times (M^K)_{n,l,n',l',n_1,l_1,n_2,l_2}^{-1} (-1)^{l_1+l_2+K+j_2-\frac{1}{2}} \hat{J}_1 \hat{J}_2 \hat{K} \begin{Bmatrix} j_1 & j_2 & K \\ l_2 & l_1 & 1/2 \end{Bmatrix} \\ & \times \frac{(-1)}{\hat{K}} {}_{SD} \langle A\lambda_f J_f \parallel (a_{n_1 l_1 j_1}^\dagger \tilde{a}_{n_2 l_2 j_2})^{(K)} \parallel A\lambda_i J_i \rangle_{SD} \end{aligned}$$

All angular dependence factorized out for plotting

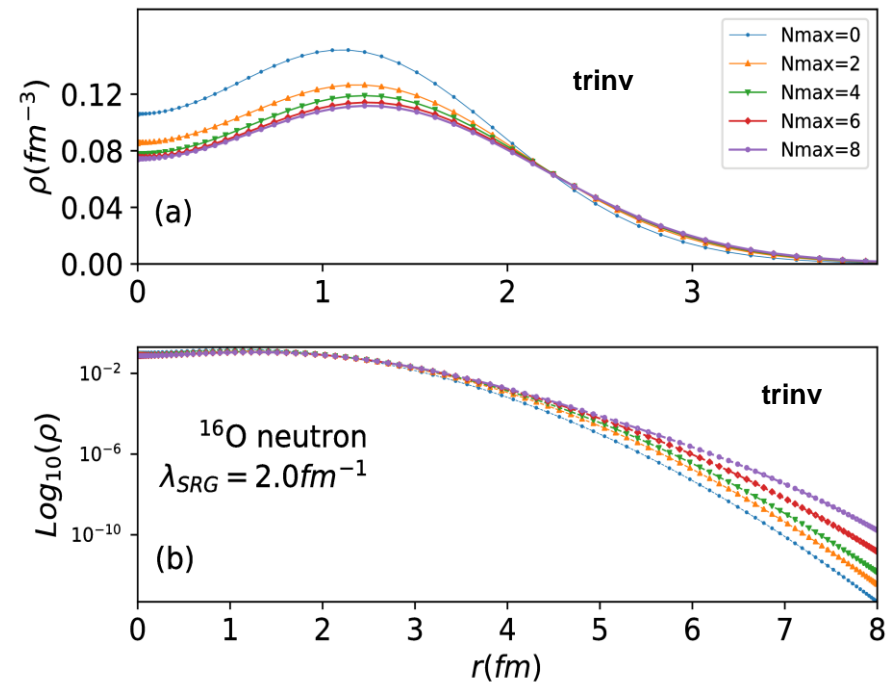
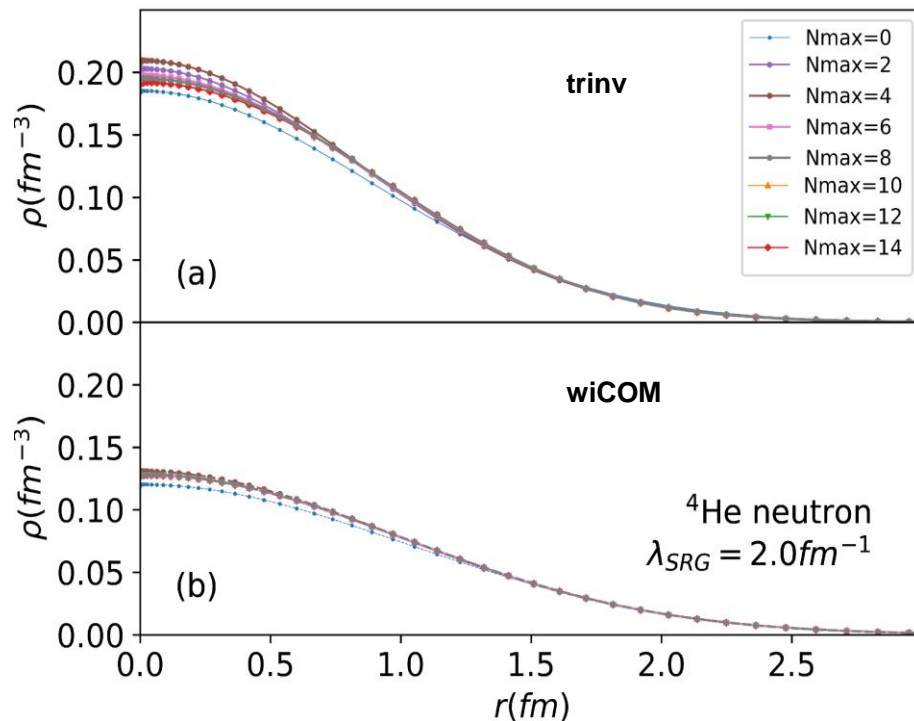
# NN and 3N interactions – N<sup>4</sup>LO(500)+3NInI

## NN systematic from LO to N<sup>4</sup>LO

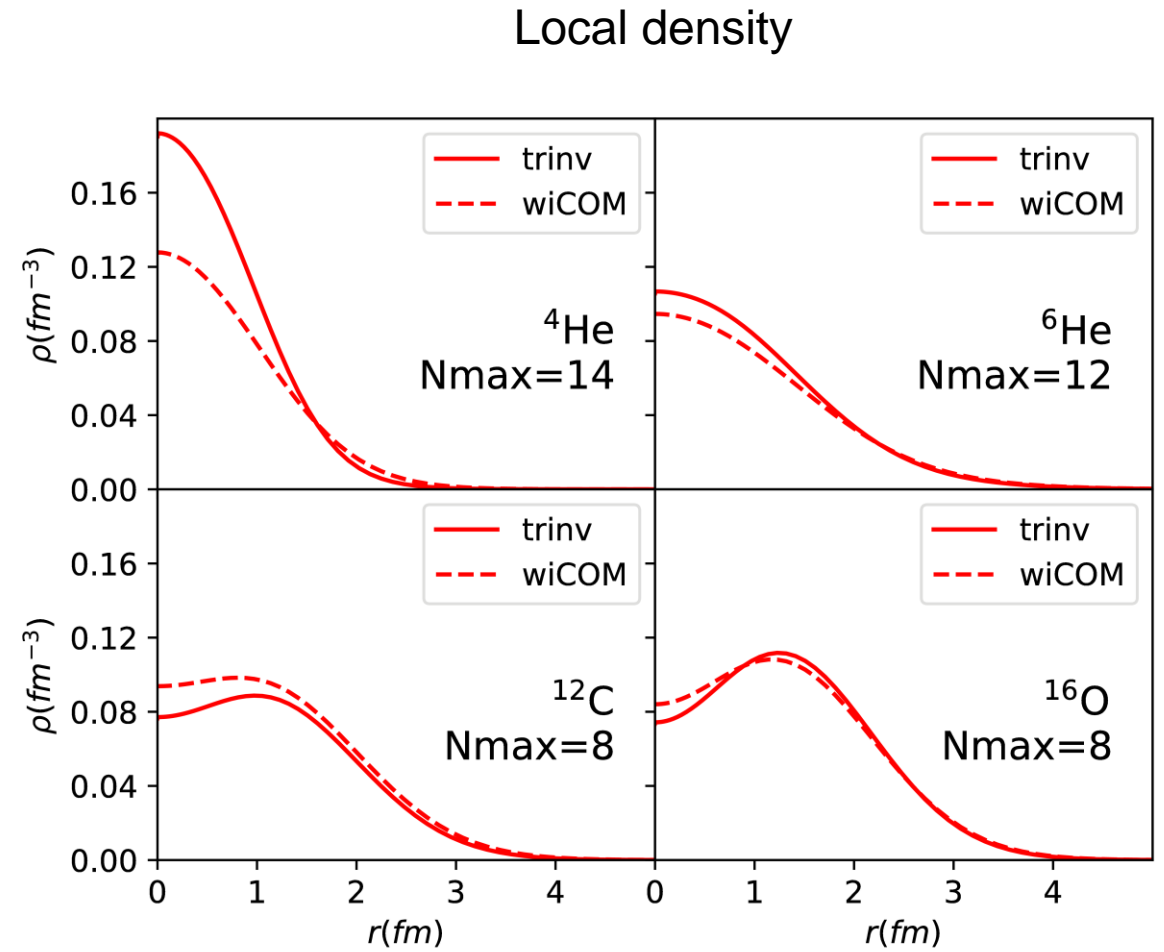
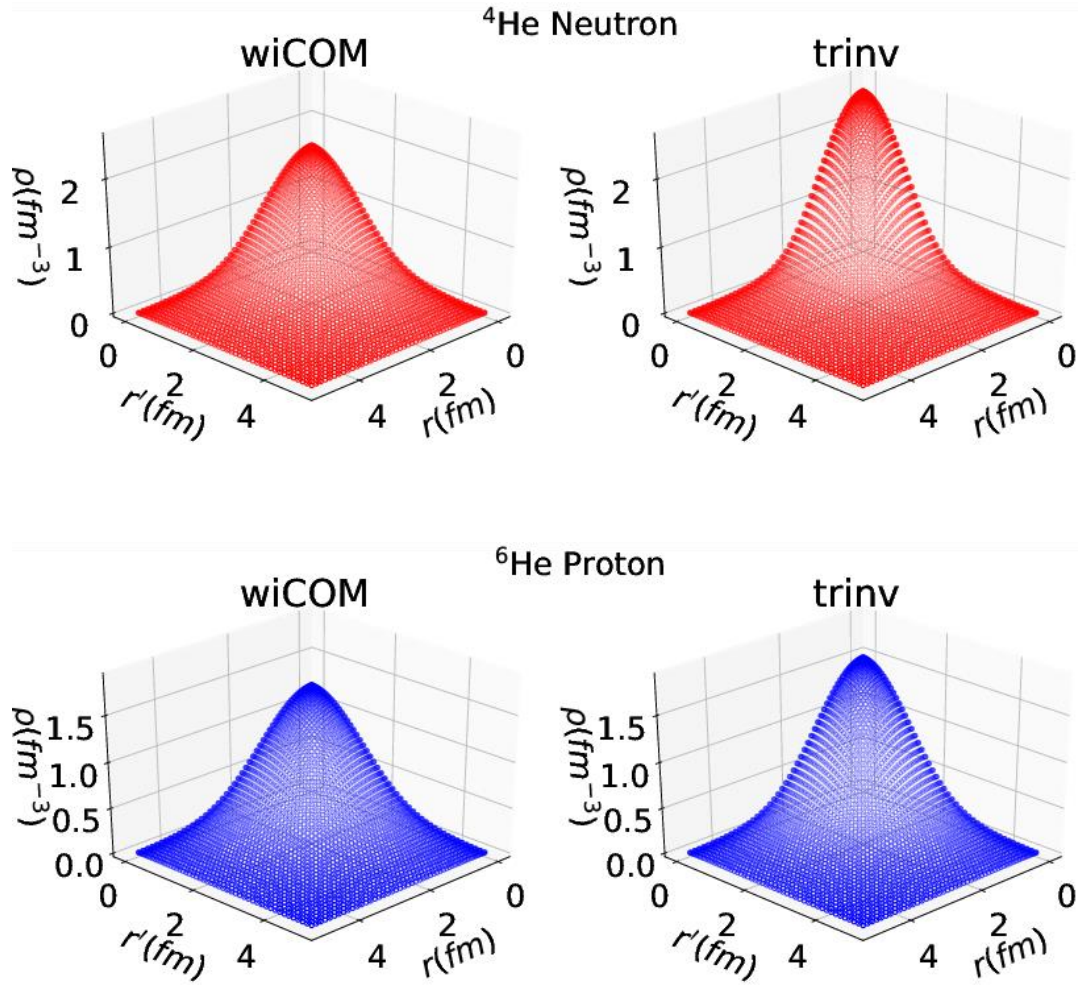
- D. R. Entem, N. Kaiser, R. Machleidt, and Y. Nosyk, Phys. Rev. C 91, 014002 (2015)
- D. R. Entem, R. Machleidt, and Y. Nosyk, arXiv:1703.05454

## 3N at N<sup>2</sup>LO

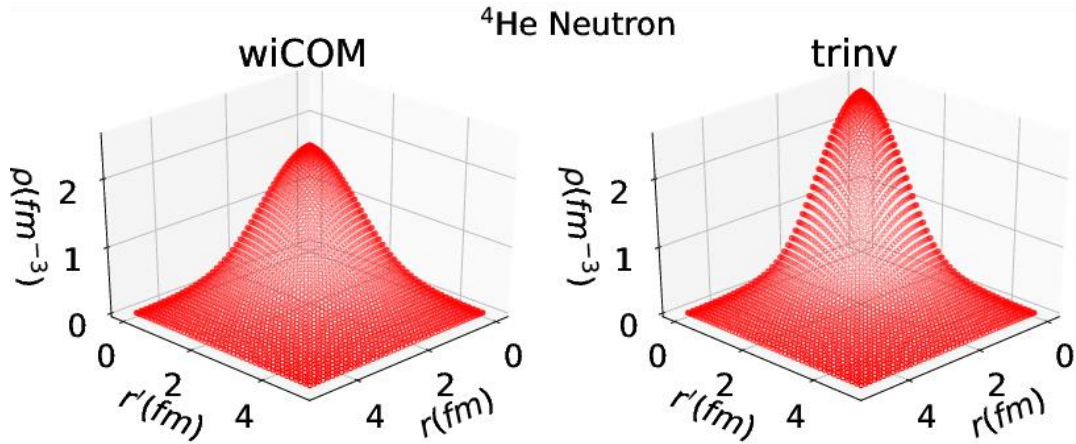
- Navrátil, 650 MeV local cut-off and 500 MeV non-local cut-off



# Density of ground state ${}^4\text{He}$ , ${}^6\text{He}$ , ${}^{12}\text{C}$ , ${}^{16}\text{O}$ with NN- $\text{N}^4\text{LO}(500)+3\text{Nlnl}$



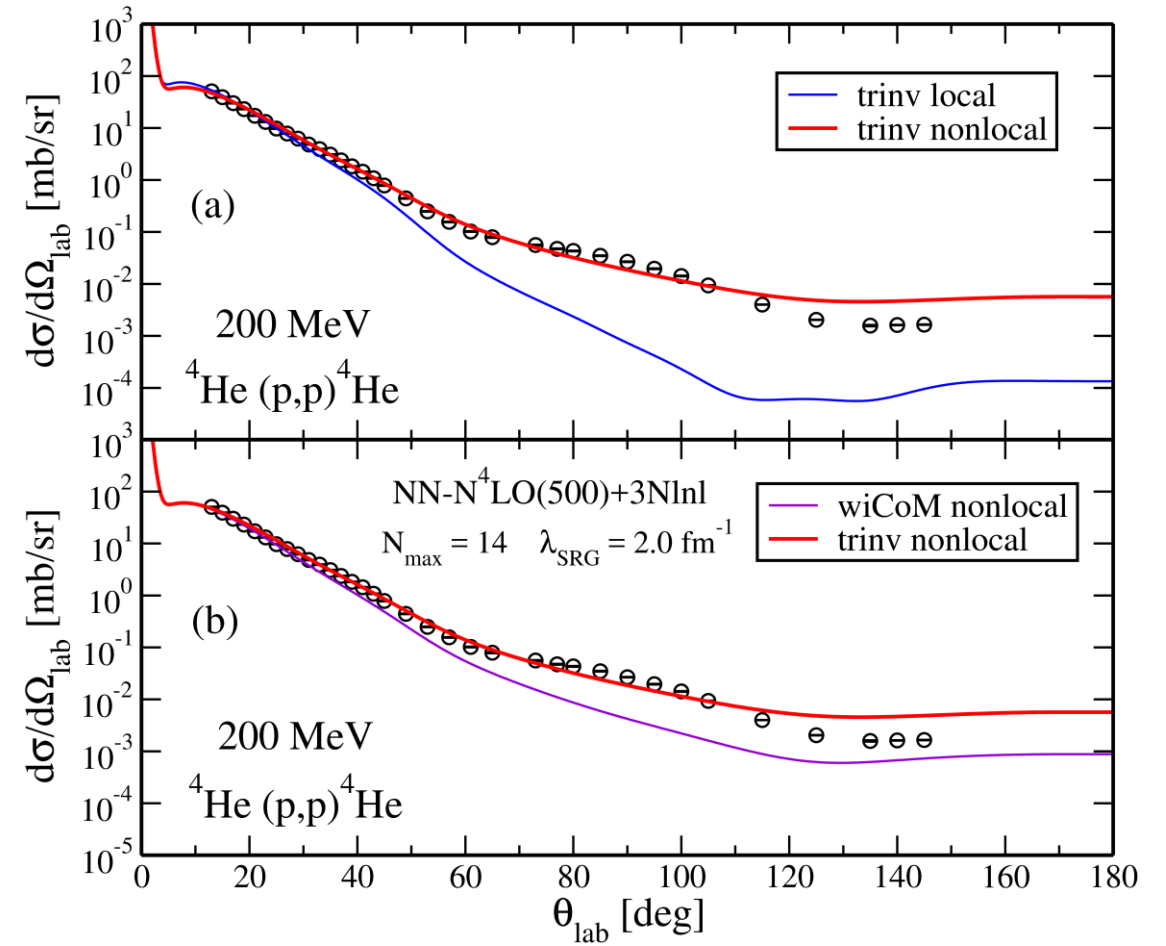
# Applications to optical potentials



- We can achieve an additional step towards consistent optical potential calculations by using the nonlocal density

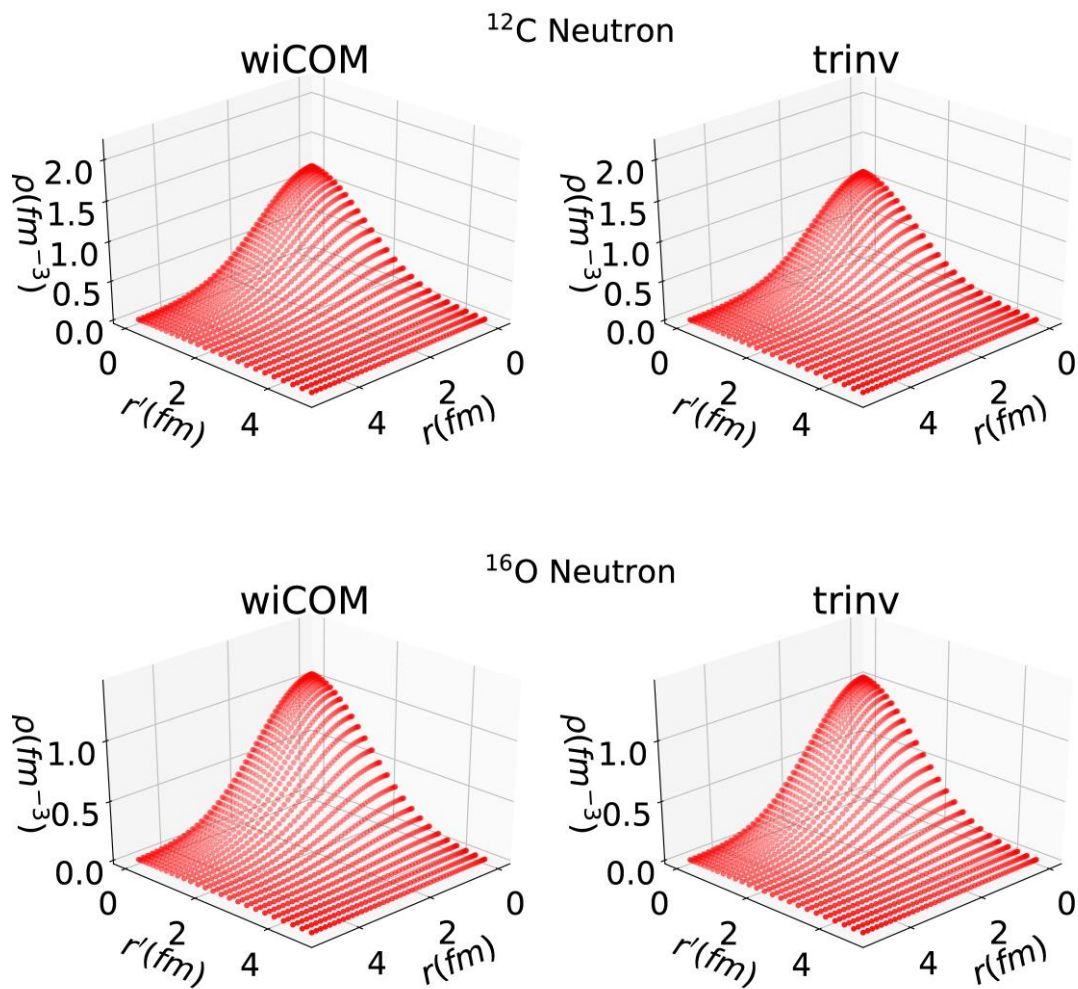
$$U(\vec{q}, \vec{K}) = \sum_{N=n,p} \int d\vec{P} \eta(\vec{q}, \vec{K}, \vec{P}) t_{pN}(\vec{q}, \vec{K}, \vec{P}) \rho_N(\vec{q}, \vec{P})$$

## Differential cross sections

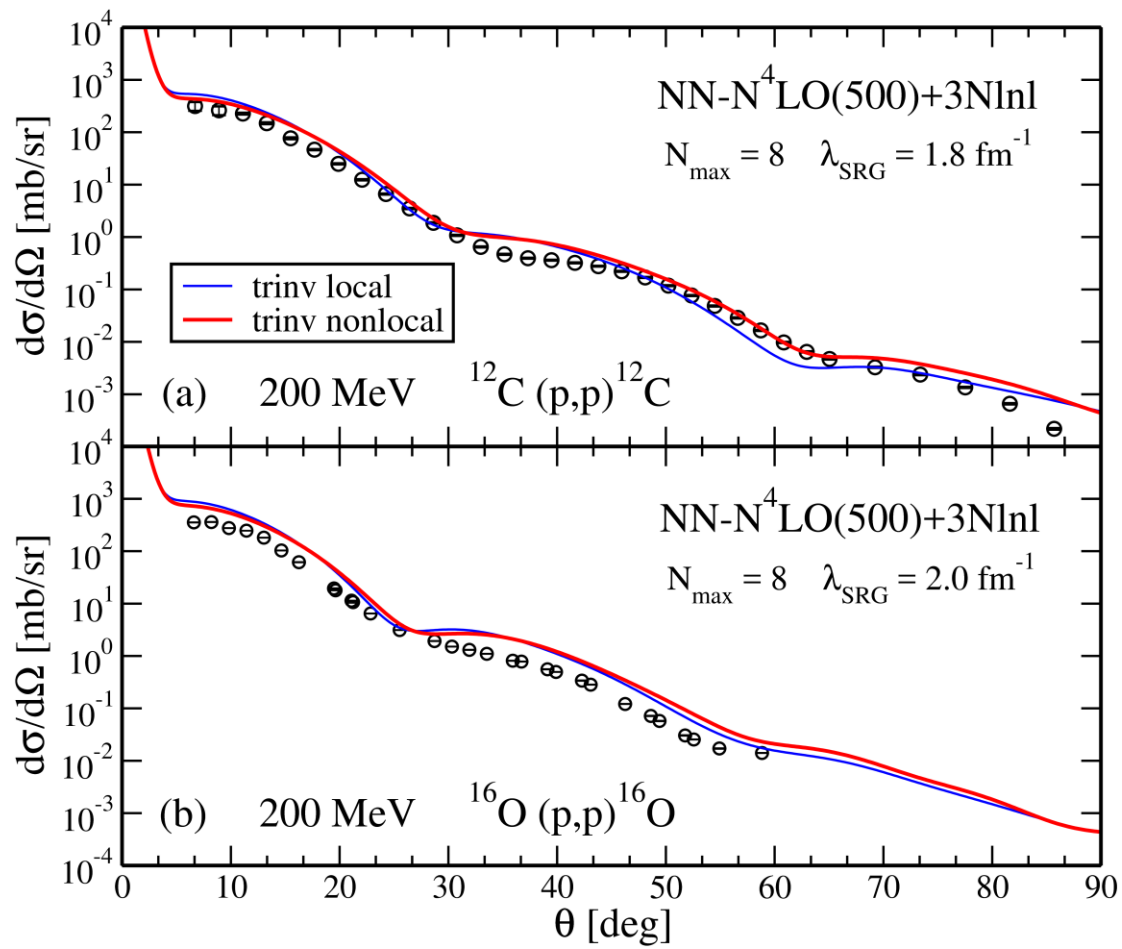




# Applications to optical potentials



## Differential cross sections



# Applications to density functional theory

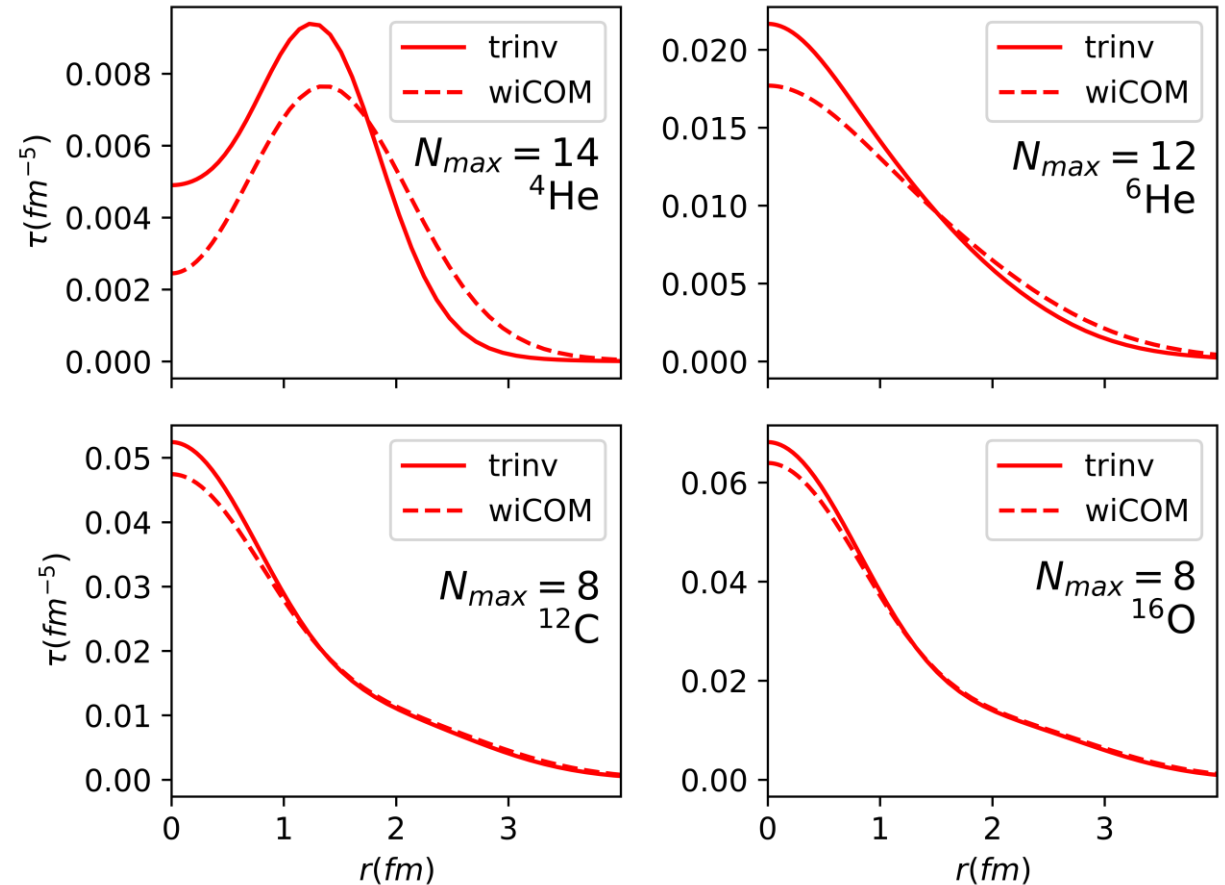
- Kinetic density (and other related densities) is a DFT quantity we are capable of calculating from *ab initio* wavefunctions according to

$$\tau_T(\vec{r}) = \left( \vec{\nabla} \cdot \vec{\nabla}' \rho_T(\vec{r}, \vec{r}') \right) \Big|_{\vec{r}=\vec{r}'}$$

$$\frac{d}{dr} R_{n,l}(r) = \frac{l}{r} R_{n,l}(r) - \frac{1}{b} \left[ \sqrt{n+l+\frac{3}{2}} \cdot R_{n,l+1}(r) + \sqrt{n} \cdot R_{n-1,l+1}(r) \right]$$

- Effects of COM removal should be amplified in DFT quantities such as the kinetic density, due to the application of gradients on the nonlocal density

Kinetic density



# Conclusions and outlook

- Conclusions
  - We observed significant differences in the nuclear density of light systems when the COM was removed
  - We can now use the more general, nonlocal density for optical potentials of nuclear reactions and *ab initio* calculations in DFT
  - More details on some of these results can be found in arXiv:1712.02879; Phys. Rev. C, in press.
- Outlook
  - We are now pursuing the use of these densities in calculations with natural orbitals (arxiv:1605.04976), reducing basis sizes and improving convergence
  - We will attempt to extend this to the translationally invariant one-body nuclear density matrix and further cut down basis sizes

Thank you  
Merci

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