

# Improved Busch formula and look for a unified calculation of nuclear scattering and structure

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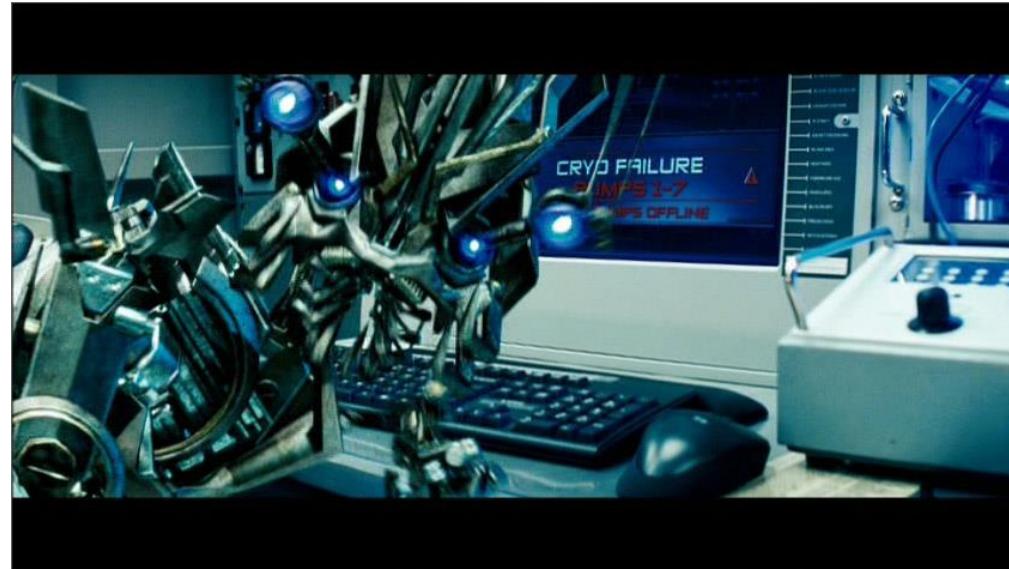
*Progress in Ab Initio Techniques in Nuclear Physics, TRIUMF, Vancouver BC, Feb 2018*

# What if aliens interfere ab initio calculation?

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# Outline

- Need to understand ab initio calculation on a macroscopic level: self-consistency, *global uncertainty analysis*, and **infrared extrapolation**
- Busch formula relates two-cluster spectrum in a harmonic trap to the two-cluster scattering
- Improve Busch formula: a toy model and effective field theory (EFT) generalization
- Test the formula and do a proof of principle calculation by studying He-5 system
- Summary and outlook

# Understand complexity on a macroscopic level

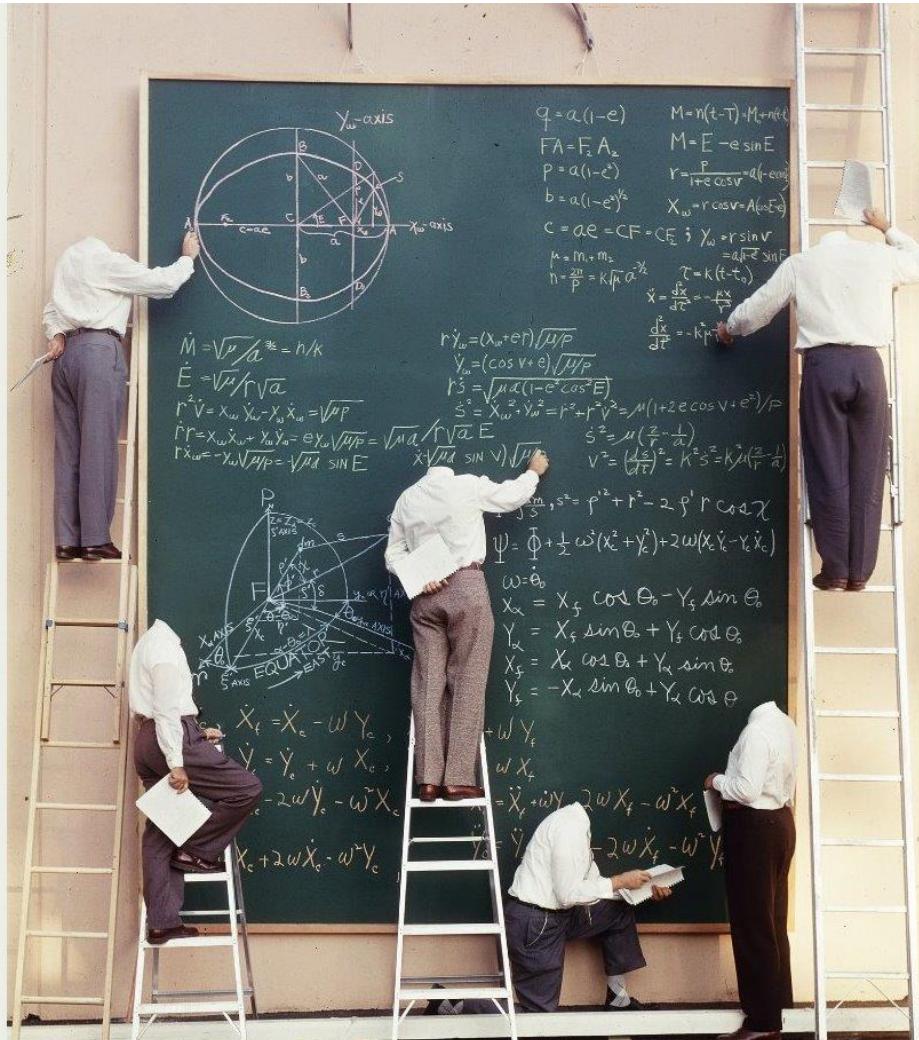
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## Can A.I. Be Taught to Explain Itself?

As machine learning becomes more powerful, the field's researchers increasingly find themselves unable to account for what their algorithms know — or how they know it.

By CLIFF KUANG NOV. 21, 2017



A historical photograph of a chalkboard filled with complex mathematical equations and diagrams, with several people on ladders working on it. The chalkboard contains various formulas related to celestial mechanics or similar fields, including:  
$$q = a(1-e) \quad M = n(t-T) - M_{\text{ref}}$$
$$FA = F_z A_z \quad M = E - e \sin E$$
$$P = a(1-e^2) \quad r = \frac{P}{1+e \cos v} = a(1-\cos e)$$
$$b = a(1-e^2)^{1/2} \quad X_w = r \cos v = a(1-e^2)$$
$$C = ae - CF = CE; \quad Y_w = r \sin v$$
$$\mu = m_1 m_2 \quad = a^2 R \sin E$$
$$n = \frac{2\pi}{P} = \sqrt{\mu/a} \quad T = k(t-t_0)$$
$$\dot{x} = \frac{dx}{dt} = -\frac{v_x}{r}$$
$$\frac{d^2x}{dt^2} = -\frac{k^2 x}{r^3}$$
  
The chalkboard also features a large circle with points labeled A through S, and several smaller circles and lines representing geometric relationships between them. Several people are standing on wooden ladders, writing on different parts of the chalkboard.

2/28/2018

3

# Understand complexity on a macroscopic level

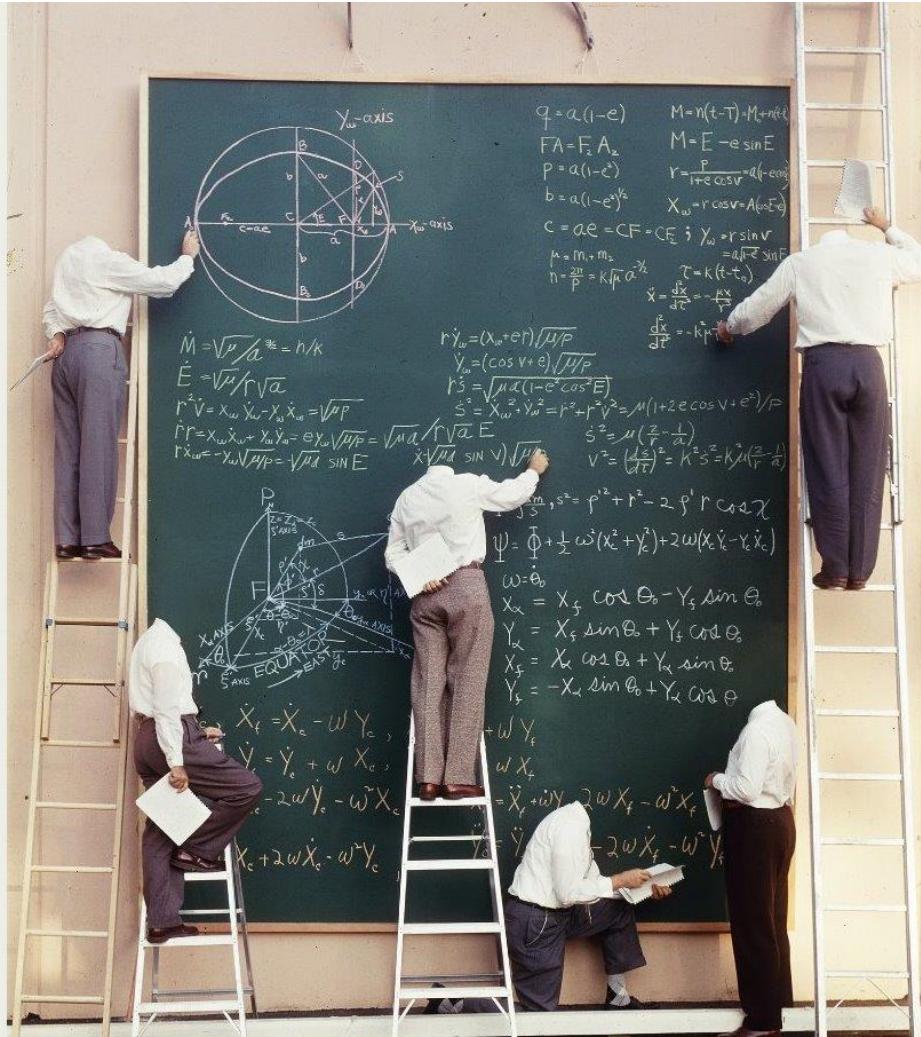
Home

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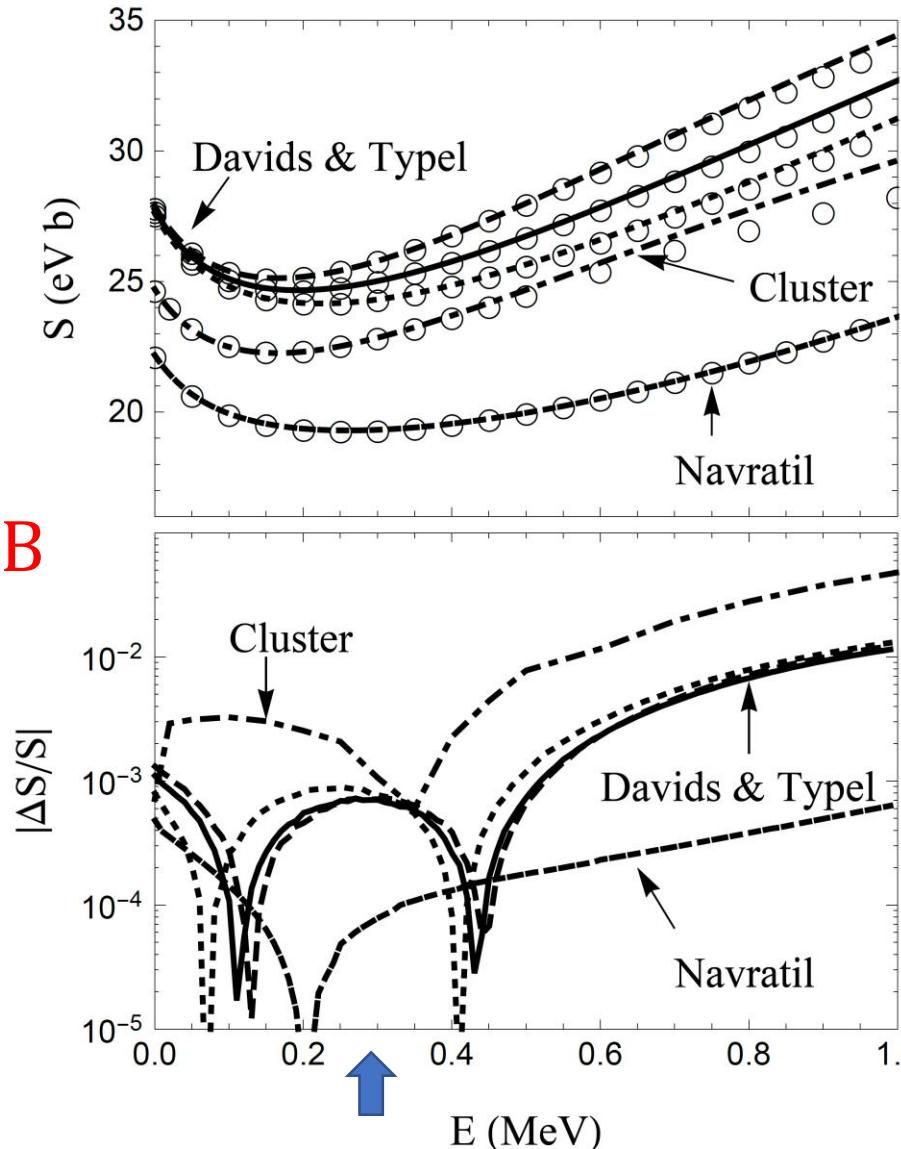
As machine learning becomes more powerful, the field's researchers increasingly find themselves unable to account for what their algorithms know — or how they know it.

By CLIFF KUANG NOV. 21, 2017

A composite image featuring a chalkboard covered in mathematical notation and diagrams. On the left, there is a diagram of a circle with points labeled A through S and various radii and angles. To the right of the diagram, several equations are written, including:  
$$q = a(1-e) \quad M = n(t-T) \cdot M_{initial}$$
$$FA = F_z A_z \quad M = E - e \sin E$$
$$P = a(1-e^2) \quad r = \frac{P}{1+e \cos v} = a(1-e)$$
$$b = a(1-e^2)^{1/2} \quad X_\omega = r \cos v = A(\omega t)$$
$$C = \alpha e - CF = CE; Y_\omega = r \sin v$$
$$M = m_1 + m_2 \quad \mu = aR \sin F$$
$$\eta = \frac{2\pi}{P} = k\sqrt{a} \quad T = k(t-t_0)$$
$$\dot{x} = \frac{dx}{dt} = -\frac{rX}{P} \quad \frac{dX}{dt} = -kP \sin F$$
$$\dot{Y}_\omega = (X_\omega + e r) / \sqrt{P}$$
$$\dot{Y}_\omega = (\cos v + e) / \sqrt{P}$$
$$\dot{r}^2 \dot{v} = X_\omega \dot{Y}_\omega - Y_\omega \dot{X}_\omega = \sqrt{P} P$$
$$\dot{r}^2 \dot{v} = \dot{X}_\omega^2 + \dot{Y}_\omega^2 = r^2 + r^2 v^2 = M(1 + 2e \cos v + e^2) / P$$
$$\dot{r} \dot{X}_\omega = X_\omega \dot{X}_\omega + Y_\omega \dot{Y}_\omega = e Y_\omega / \sqrt{P} = \sqrt{a} a / \sqrt{a} E$$
$$\dot{r} \dot{X}_\omega = r^2 \dot{v} \sin v / \sqrt{P} = r^2 \sqrt{a} \sin v / \sqrt{P}$$
$$\dot{r}^2 = \dot{X}_\omega^2 + \dot{Y}_\omega^2 = r^2 + r^2 v^2 = M(1 + 2e \cos v + e^2) / P$$
$$v^2 = \left(\frac{dr}{dt}\right)^2 = k^2 s^2 = k^2 / (r^2 - a^2)$$
$$m, s^2 = r^2 + r^2 - 2r^2 \cos v$$

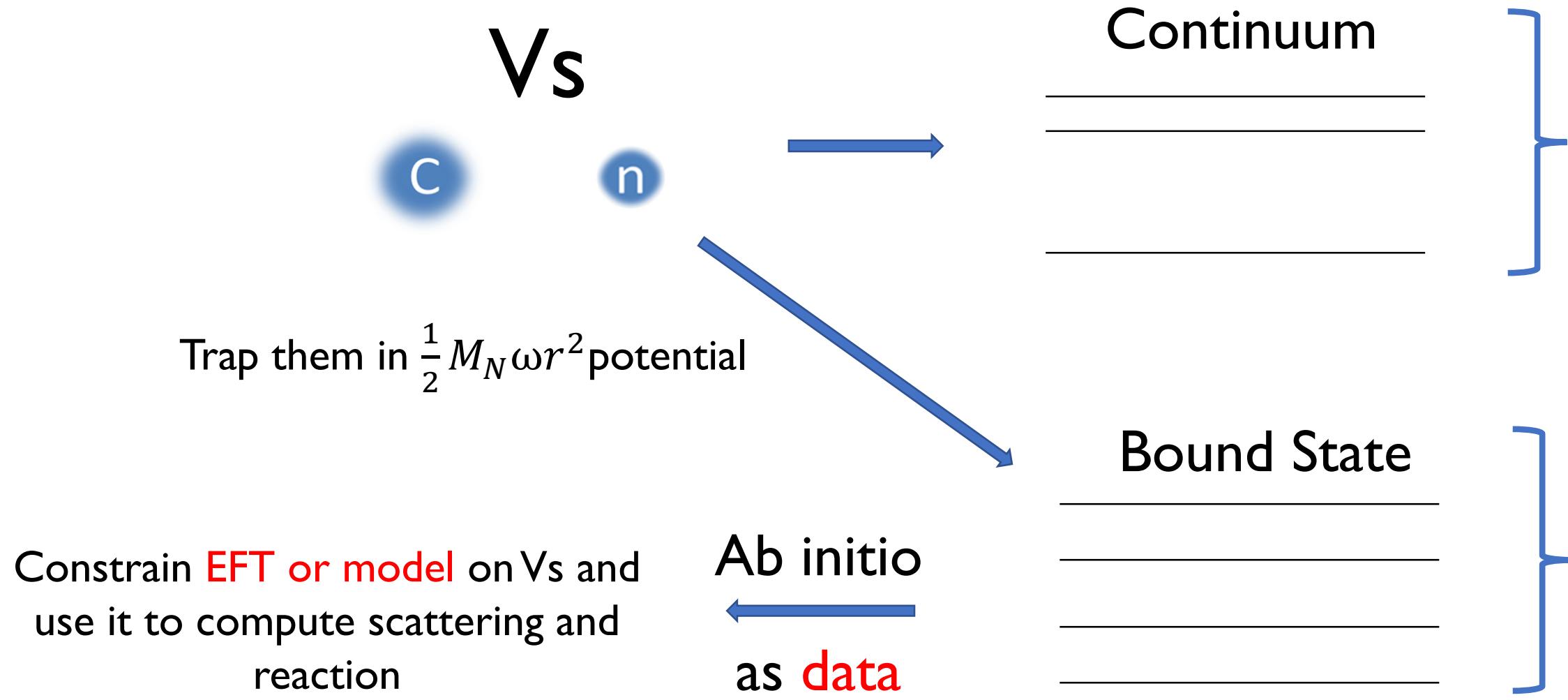
# Understand ab initio calculation on a macroscopic level: consistency and global uncertainty analysis

${}^7_4\text{Be}(\text{p}, \gamma){}^8_5\text{B}$



- 1) Halo-EFT: scattering lengths, effective ranges, ANC,,,
- 2) Computational “data” on different observables → macroscopic picture: a **systematical approach is needed**
- 3) Various information sources: experimental, computational, and empirical. The EFT approach can be used to combine these information and produce global uncertainty propagation

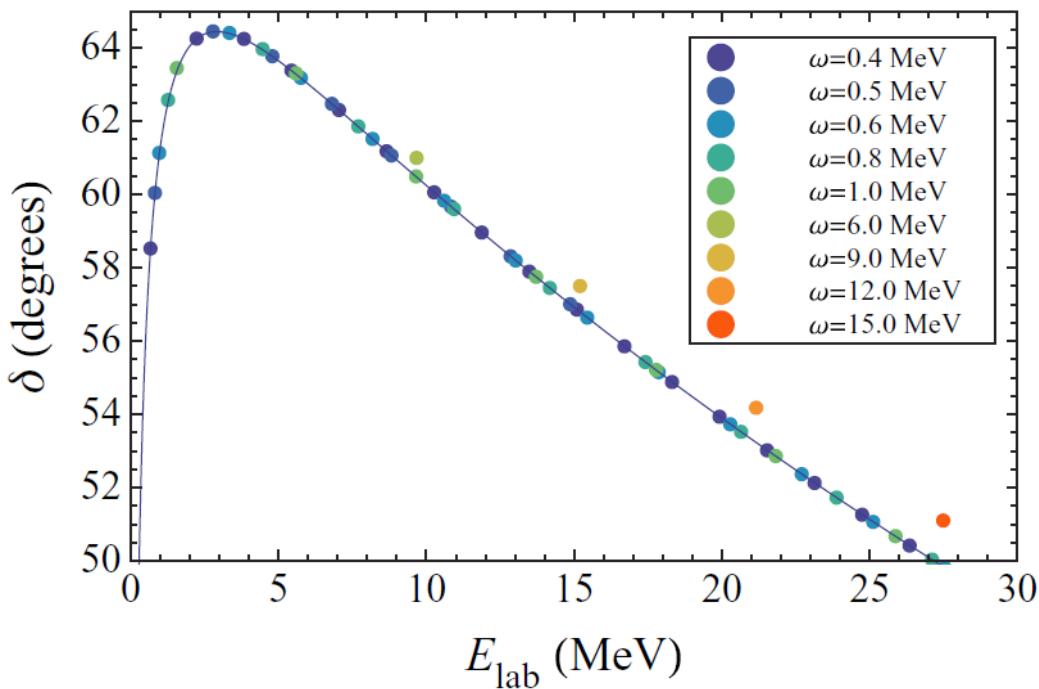
# Busch formula (infrared extrapolation)



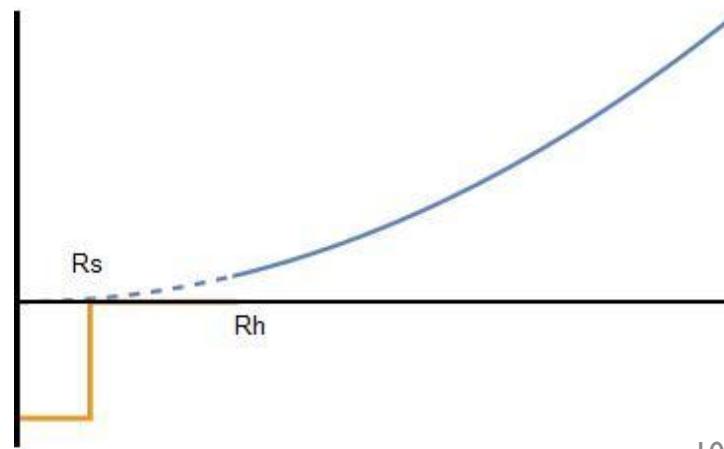
# Busch formula

$$p^{2l+1} \cot \delta_l(p) = (-1)^{l+1} (4M_R \omega)^{l+\frac{1}{2}} \frac{\Gamma\left(\frac{3}{4} + \frac{l}{2} - \frac{\epsilon}{2}\right)}{\Gamma\left(\frac{1}{4} - \frac{l}{2} - \frac{\epsilon}{2}\right)}$$

with  $\epsilon \equiv \frac{E}{\omega}$ ,  $E \equiv \frac{p^2}{2M_R}$



- 1) T. Luu, M. Savage, A. Schwenk, and J. Vary, PRC (2010): N-N phase shift
- 2) J. Rotureau, I. Stetcu, B.R. Barrett, and U. van Kolck, PRC (2012): N-D phase shift



# Improve Busch Formula: a model

$$V_s(r) = \begin{cases} +\infty & \text{when } r \leq r_c \\ 0 & \text{when } r > r_c , \end{cases}$$

$$\begin{aligned}
 & p^{2l+1} \cot \delta_l(p) - (-1)^{l+1} (4M_R \omega)^{l+\frac{1}{2}} \frac{\Gamma\left(\frac{3}{4} + \frac{l}{2} - \frac{\epsilon}{2}\right)}{\Gamma\left(\frac{1}{4} - \frac{l}{2} - \frac{\epsilon}{2}\right)} \\
 &= -\frac{(2l+1)!!(2l-1)!!}{r_c^{2l+1}} \left[ \frac{(2l+1)\left(\frac{r_c}{b}\right)^4}{2(2l-3)(2l+5)} + \frac{(2l+1)(6l+25)\frac{1}{2}(pr_c)^2\left(\frac{r_c}{b}\right)^4}{3(2l-5)(2l+3)(2l+5)(2l+7)} + O\left[\left(\frac{r_c}{b}\right)^8, (pr_c)^4\left(\frac{r_c}{b}\right)^4\right] \right] \\
 &\equiv -L_{al} \frac{1}{b^4 r_c^{2l-3}} - L_{rl} \frac{p^2}{b^4 r_c^{2l-1}} + \dots
 \end{aligned}$$

Note:  $b = \sqrt{\frac{1}{M_R \omega}}$

**Effective range expansion (ERE):**

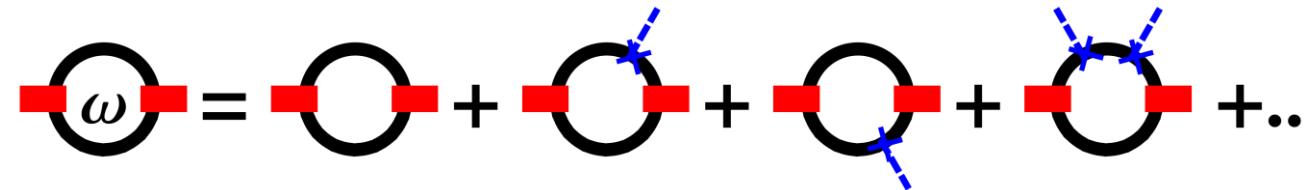
$$p^{2l+1} \cot \delta_l(p) = -\frac{\Lambda^{2l+1}}{a_l} + \frac{1}{2} r_l \Lambda^{2l-1} p^2 + \frac{1}{4} \tilde{r}_l^{(1)} \Lambda^{2l-3} p^4$$

# Improve Busch Formula: EFT

$$\mathcal{L}_0 = \begin{pmatrix} c^* & n^* & -\phi^* \end{pmatrix} \text{diag} \left( i\partial_t - \hat{m}_c \psi + \frac{\partial^2}{2M_c}, i\partial_t - \hat{m}_n \psi + \frac{\partial^2}{2M_n}, i\partial_t - \hat{m}_\phi \psi + \frac{\partial^2}{2M_{nc}} + \Delta_0 \right) \begin{pmatrix} c & n & \phi \end{pmatrix}^T$$

$$\mathcal{L}_{I0} = g_0 \phi^* c n - \phi^* \left[ \sum_{j=2} d_j^{(0)} \left( i\partial_t - \hat{m}_\phi \psi + \frac{\partial^2}{2M_{nc}} \right)^j \right] \phi + \text{C.C.} \quad \text{Note: } \psi = \frac{1}{2} m_N \omega^2 r^2$$

Self-energy bubble:



Dimer-field propagator:



$$p_{\tilde{E}} \cot \delta_0(\tilde{E}) = -\frac{2\pi}{g_0^2 M_R} (\Sigma_\omega(\tilde{E}) - \Sigma(\tilde{E})) \text{ reproduces the Busch formula.}$$

Then what went wrong?

# Improve Busch Formula: EFT

$$\mathcal{L}_{I0} = g_0 \phi^* c n - \phi^* \left[ \sum_{j=2} d_j^{(0)} \left( i\partial_t - \hat{m}_\phi \psi + \frac{\partial^2}{2M_{nc}} \right)^j \right] \phi + \text{C.C.}$$

$$- \phi^* \underbrace{\left[ \sum_{j=0} \sum_{k=1} d_{j,k}^{(0)} \left( i\partial_t - \hat{m}_\phi \psi + \frac{\partial^2}{2M_{nc}} \right)^j \left( \frac{M_R^2}{3m} \partial^2 \psi \right)^k \right]}_{\text{---} + \text{---} \square \text{---} + \dots \text{---} \times \text{---} + \dots} \phi$$

The factorizability of CM motion severely constrains two-body current like couplings.

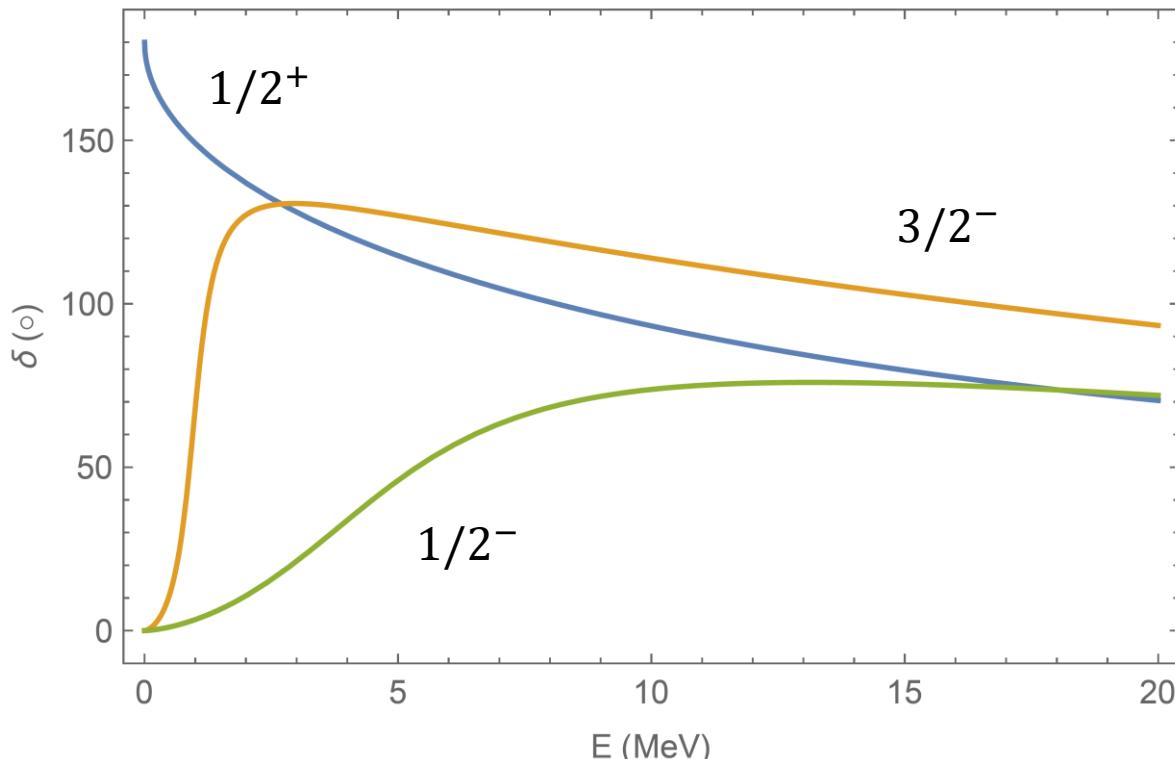
$$\Sigma_\omega \rightarrow \Sigma_\omega + \sum_{j=0} \sum_{k=1} d_{j,k}^{(0)} \tilde{E}^j b^{-4k}$$

$$p \cot \delta_0(p) \rightarrow p \cot \delta_0(p) + L_{a_0} \frac{1}{b^4 \Lambda^3} + L_{r_0} \frac{p^2}{b^4 \Lambda^5} + L_{\tilde{r}_0^{(1)}} \frac{p^4}{b^4 \Lambda^7} + \dots$$

$$= -\frac{\Lambda}{a_0} + \frac{1}{2} \frac{r_0}{\Lambda} p^2 + \frac{1}{4} \frac{\tilde{r}_0^{(1)}}{\Lambda^3} p^4 + L_{a_0} \frac{1}{b^4 \Lambda^3} + L_{r_0} \frac{p^2}{b^4 \Lambda^5} + L_{\tilde{r}_0^{(1)}} \frac{p^4}{b^4 \Lambda^7}$$

# Test: $n - \alpha$ system

$$V_s(r) = \begin{cases} V_0(1 + \beta \mathbf{L} \cdot \boldsymbol{\sigma}) & \text{when } r < r_c \\ 0 & \text{when } r > r_c \end{cases}$$

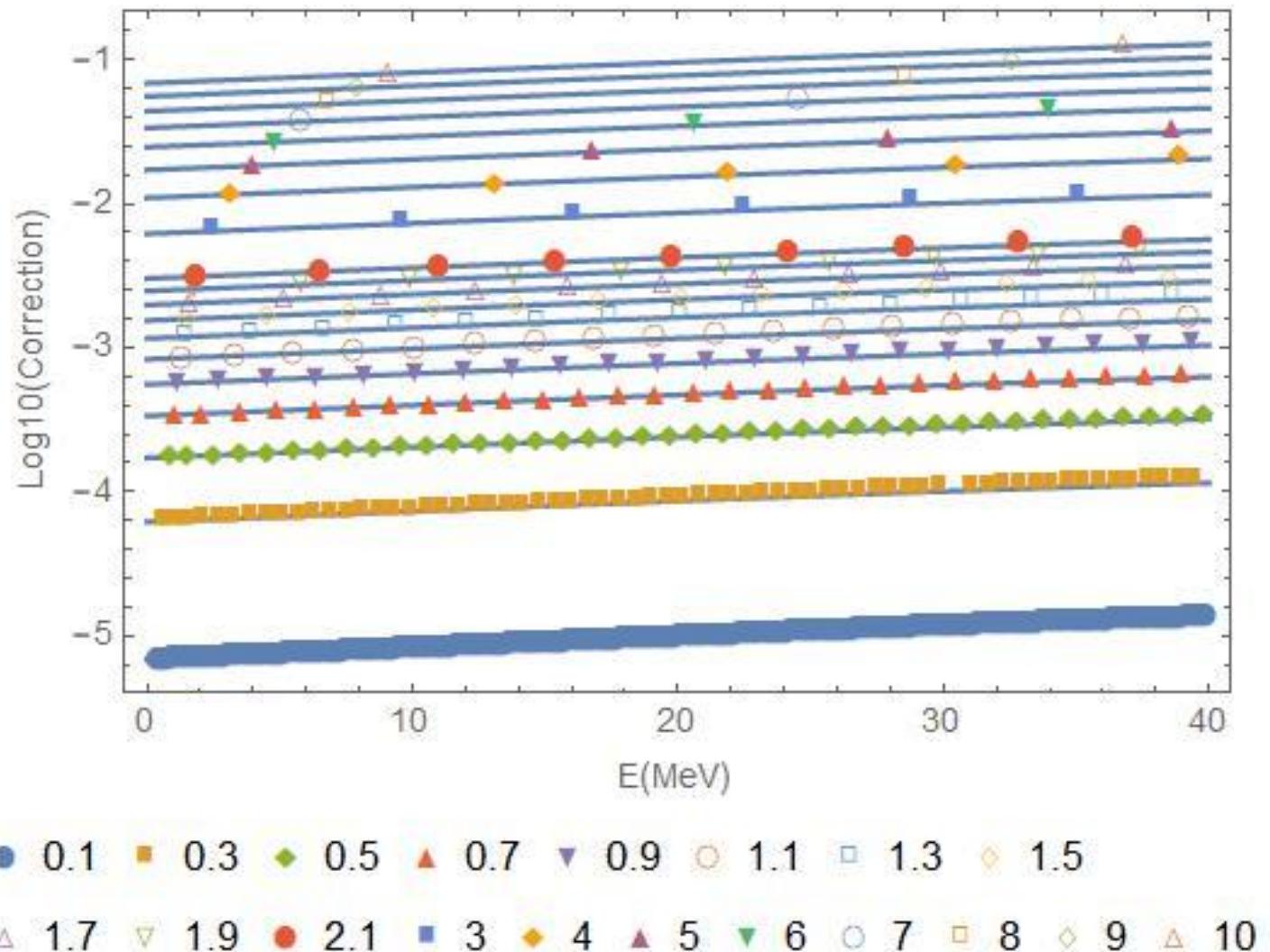


S.Ali et.al., RMP 57, 923 (1985)

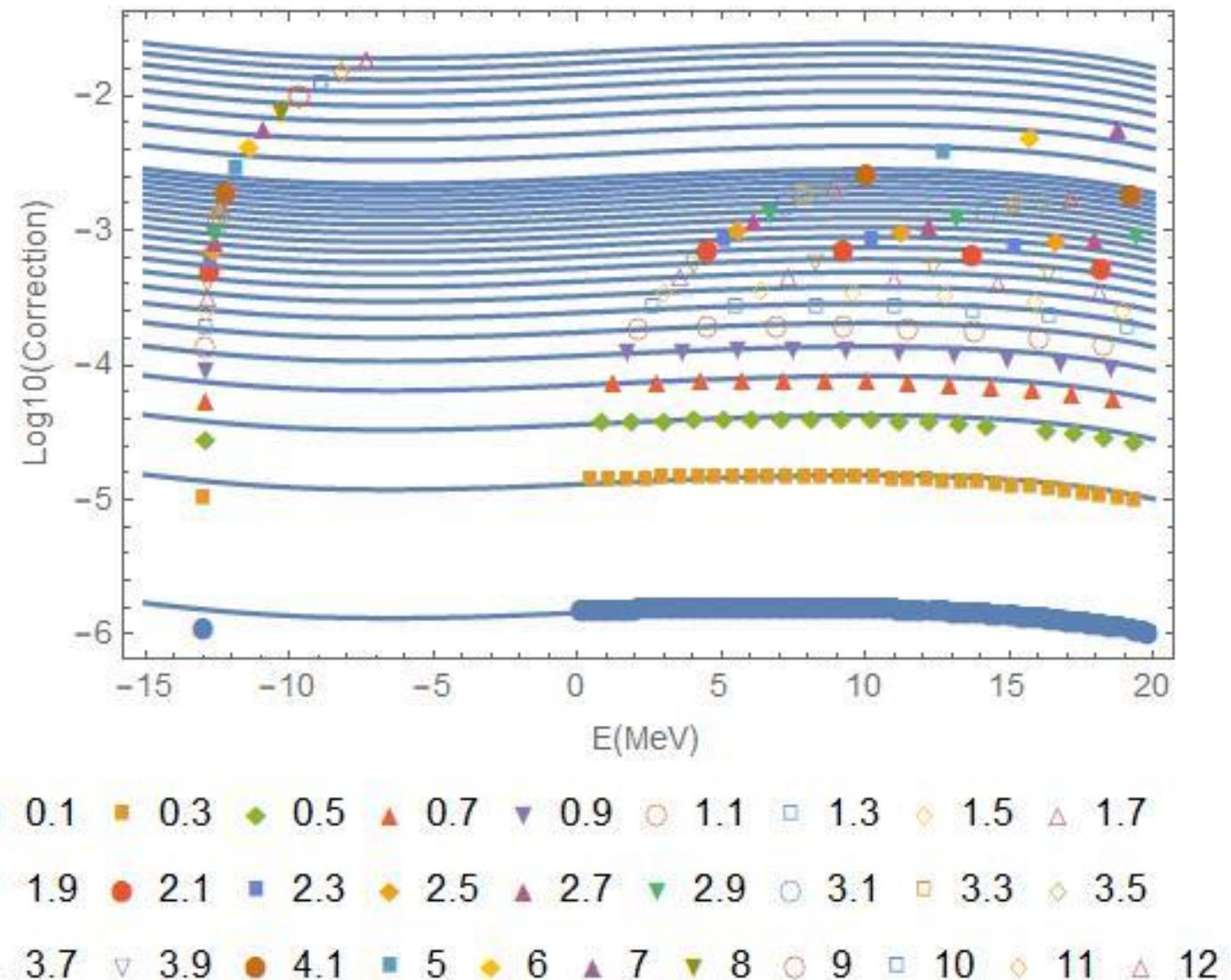
# Test: $n - \alpha$ system in p-wave

$$p^{2l+1} \cot \delta_l(p) - (-1)^{l+1} (4M_R\omega)^{l+\frac{1}{2}} \frac{\Gamma\left(\frac{3}{4} + \frac{l}{2} - \frac{\epsilon}{2}\right)}{\Gamma\left(\frac{1}{4} - \frac{l}{2} - \frac{\epsilon}{2}\right)}$$

- 1) Corrections  $\propto \omega^2$
- 2) At  $\omega = 10$  MeV,  
10% corrections



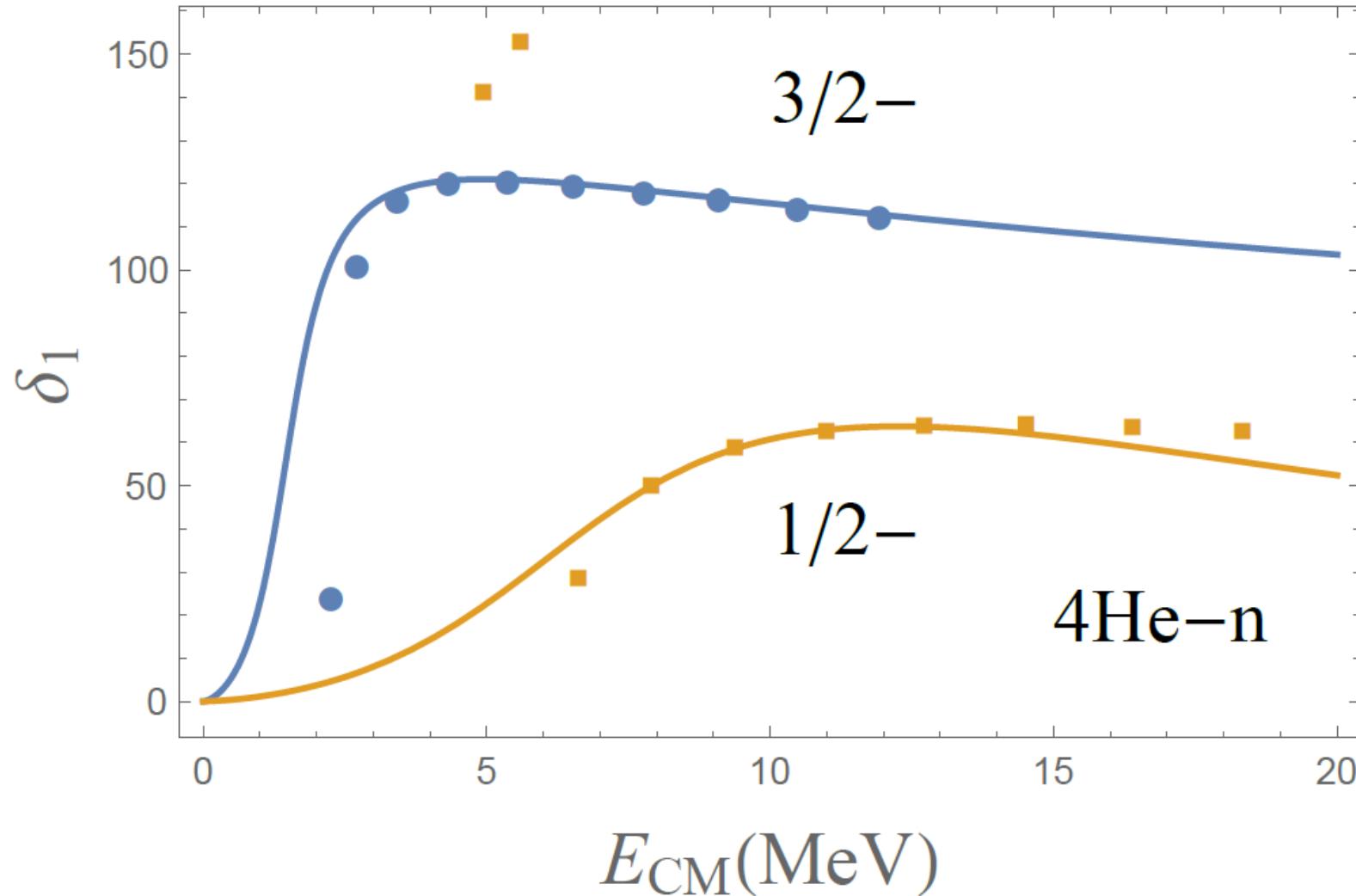
# Test: $n - \alpha$ system in s-wave



# Test: extract $n - \alpha$ ERE parameters

	$a$	$r$	$\tilde{r}^{(1)}$	$\tilde{r}^{(2)}$	$L_a$	$L_r$	$L_{\tilde{r}^{(1)}}$	$L_{\tilde{r}^{(2)}}$
P-wave fit J=3/2	-43.4101	-1.11405	2.23369	0.605012	-1.87099	-0.827113	-0.269323	0.0685459
P-wave J=3/2	-43.408	-1.11425	2.23965	0.5243				
P-wave fit J=1/2	-7.19466	-0.882978	2.44336	0.86573	-2.08059	-1.00758	-0.706915	0.257639
P-wave J=1/2	-7.19286	-0.885359	2.49157	0.668183				
S-wave fit J=1/2	2.82818	1.57966	0.700462	-0.22943	-0.405303	-0.240855	-0.172349	0.867527
S-wave fit J=1/2	2.82818	1.57967	0.700221	-0.076691				

# Trial results by analyzing IM-SRG “data” from G. Chan, R. Stroberg, and J. Holt



# Summary and outlook

- Need to understand ab initio calculation on a macroscopic level **in a systematical way**
- The improved Busch formula can be used to infer scattering (infrared extrapolation)
- Test on  $n - \alpha$  is encouraging
- Working with P. Narvatal on  $n - \alpha$
- Also applying it to study  $n - {}^{24}O$  with G. Chan, R. Stroberg, and J. Holt
- Consider generalizing it to study two-cluster reactions and three-cluster systems