

# Towards multi-major- shell effective hamiltonian from IMSRG

Takayuki Miyagi

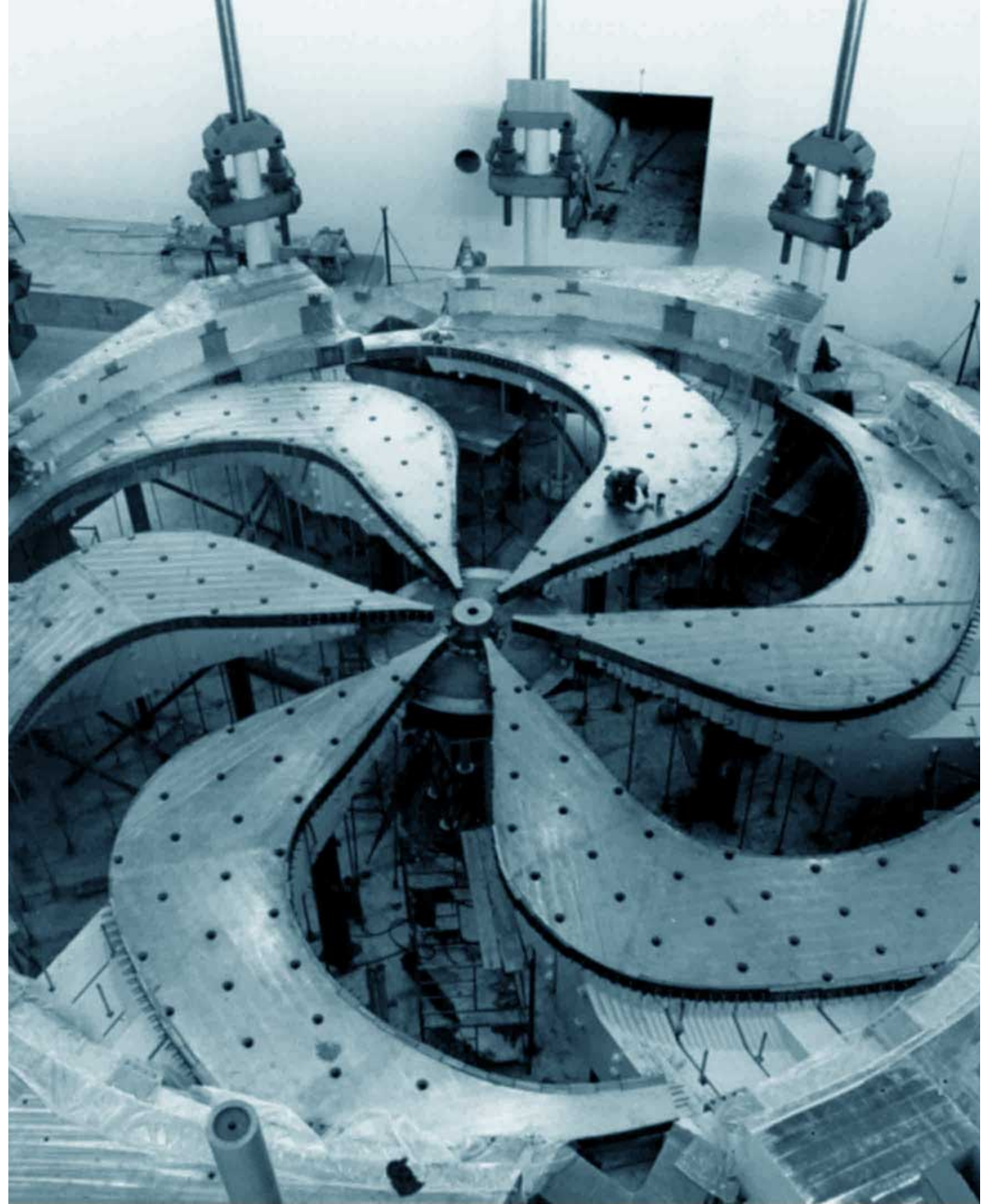
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Progress in Ab Initio Techniques in Nuclear Physics

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# Motivation

- Ab initio extension to shell-model calculation is a powerful tool.
- CCM+SM
  - ◆ G. R. Jansen, M. D. Schuster, A. Signoracci, G. Hagen, and P. Navrátil, Phys. Rev. C **94**, 011301 (2016).
  - ◆ Z. H. Sun, T. D. Morris, G. Hagen, G. R. Jansen, and T. Papenbrock, Phys. Rev. C **98**, 054320 (2018).
- IM-SRG+SM
  - ◆ S. R. Stroberg, H. Hergert, J. D. Holt, S. K. Bogner, and A. Schwenk, Phys. Rev. C **93**, 051301 (2016).
  - ◆ S. R. Stroberg, A. Calci, H. Hergert, J. D. Holt, S. K. Bogner, R. Roth, and A. Schwenk, Phys. Rev. Lett. **118**, 032502 (2017).

# Motivation

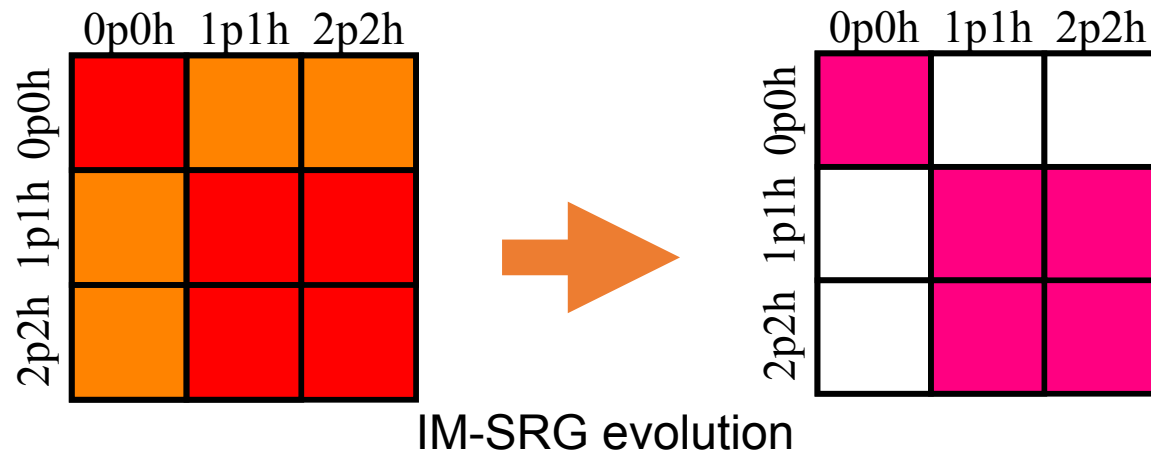
- Two-major-shell effective Hamiltonian is required:
  - ◆ For unnatural parity state
  - ◆ For excitation spectra for doubly closed ( $^{16}\text{O}$ ,  $^{40}\text{Ca}$ , ...)
  - ◆ For exotic nuclei far from stability (Island of inversion, ...)
  - ◆ ...

# Motivation

- Shell-model effective Hamiltonian across the major shells:
  - ◆ Phenomenological
    - ✧ *psd*-shell, *sdpf*-shell, ...
  - ◆ Ab initio
    - ✧ Many-body perturbation theory
    - ✧ IM-SRG (this work)

## IM-SRG (Quick reminder)

- In-medium similarity renormalization group:



$$H \longrightarrow H(s) \approx E(s) + \sum_{12} f_{12}(s) a_1^\dagger a_2 + \frac{1}{4} \sum_{1234} \Gamma_{1234}(s) a_1^\dagger a_2^\dagger a_4 a_3$$

s: flow parameter

$$\frac{dE}{ds} = \sum_{ab} (n_a - n_b) \eta_{ab} f_{ba} + \frac{1}{2} \sum_{abcd} \eta_{abcd} \Gamma_{cdab} n_a n_b \bar{n}_c \bar{n}_d$$

$$\begin{aligned} \frac{df_{12}}{ds} &= \sum_a (1 + P_{12}) \eta_{1a} f_{a2} + \frac{1}{2} \sum_{ab} (n_a - n_b) (\eta_{ab} \Gamma_{b1a2} - f_{ab} \eta_{b1a2}) \\ &+ \frac{1}{2} \sum_{abc} (n_a n_b \bar{n}_c + \bar{n}_a \bar{n}_b n_c) (1 + P_{12}) \eta_{c1ab} \Gamma_{abc2} \end{aligned}$$

$$\begin{aligned} \frac{d\Gamma_{1234}}{ds} &= \sum_a [(1 - P_{12})(\eta_{1a} \Gamma_{a234} - f_{1a} \eta_{a234}) - (1 - P_{34})(\eta_{a3} \Gamma_{12a4} - f_{a3} \eta_{12a3})] \\ &+ \frac{1}{2} \sum_{ab} (1 - n_a - n_b) (\eta_{12ab} \Gamma_{ab34} - \Gamma_{12ab} \eta_{ab34}) \\ &- \sum_{ab} (n_a - n_b) (1 - P_{12})(1 - P_{34}) \eta_{b2a4} \Gamma_{a1b3} \end{aligned}$$

$n_a$  : occupation number

$$\bar{n}_a = 1 - n_a$$

## IM-SRG (Quick reminder)

- In-medium similarity renormalization group:

Generator is chosen to suppress the off diagonal component:

$$\eta_{12} = \frac{f_{12}}{f_{11} - f_{22} + \Gamma_{1212}}$$

$$\eta_{1234} = \frac{\Gamma_{1234}}{f_{11} + f_{22} - f_{33} - f_{44} + A_{1234}}$$

$$A_{1234} = \Gamma_{1212} + \Gamma_{3434} - \Gamma_{1313} - \Gamma_{2424} - \Gamma_{1414} - \Gamma_{2323}$$

$f_{12}, \Gamma_{1234}$  : matrix element we want to suppress

$$\frac{dE}{ds} = \sum_{ab} (n_a - n_b) \eta_{ab} f_{ba} + \frac{1}{2} \sum_{abcd} \eta_{abcd} \Gamma_{cdab} n_a n_b \bar{n}_c \bar{n}_d$$

$$\begin{aligned} \frac{df_{12}}{ds} = & \sum_a (1 + P_{12}) \eta_{1a} f_{a2} + \frac{1}{2} \sum_{ab} (n_a - n_b) \eta_{ab} \Gamma_{b1a2} - f_{a2} \eta_{b1a2} \\ & + \frac{1}{2} \sum_{abc} (n_a n_b \bar{n}_c + \bar{n}_a \bar{n}_b n_c) (1 + P_{12}) \eta_{c1ab} \Gamma_{abc2} \end{aligned}$$

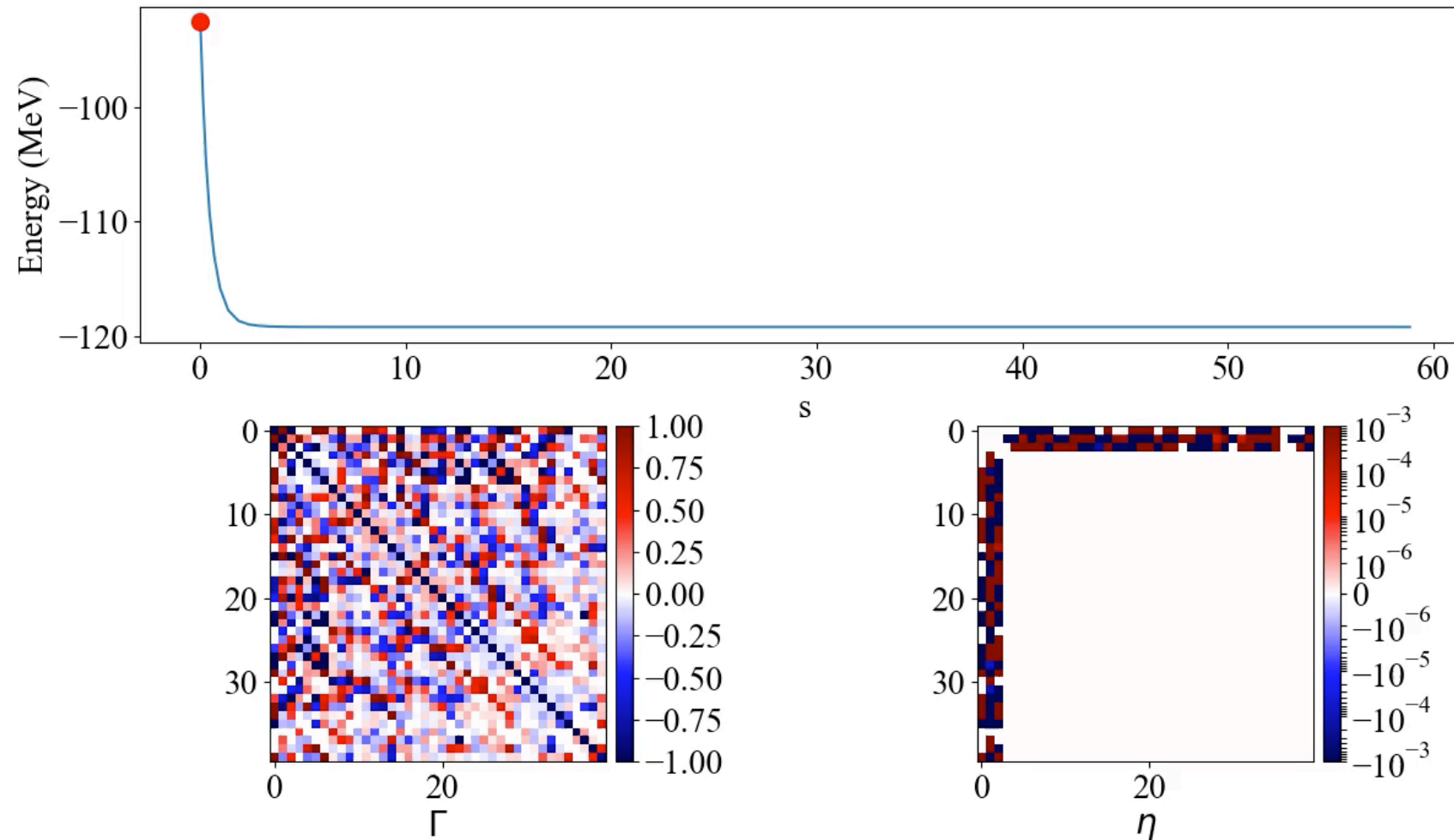
$$\begin{aligned} \frac{d\Gamma_{1234}}{ds} = & \sum_a [(1 - P_{12}) \eta_{1a} \Gamma_{a234} - f_{1a} \eta_{a234}] - (1 - P_{34}) \eta_{a31} \Gamma_{12a4} - f_{a3} \eta_{12a3} \\ & + \frac{1}{2} \sum_{ab} (1 - n_a - n_b) \eta_{12ab} \Gamma_{ab34} - \Gamma_{12a} \eta_{ab34} \\ & - \sum_{ab} (n_a - n_b) (1 - P_{12}) (1 - P_{34}) \eta_{b2a4} \Gamma_{a1b3} \end{aligned}$$

$n_a$  : occupation number

$\bar{n}_a = 1 - n_a$

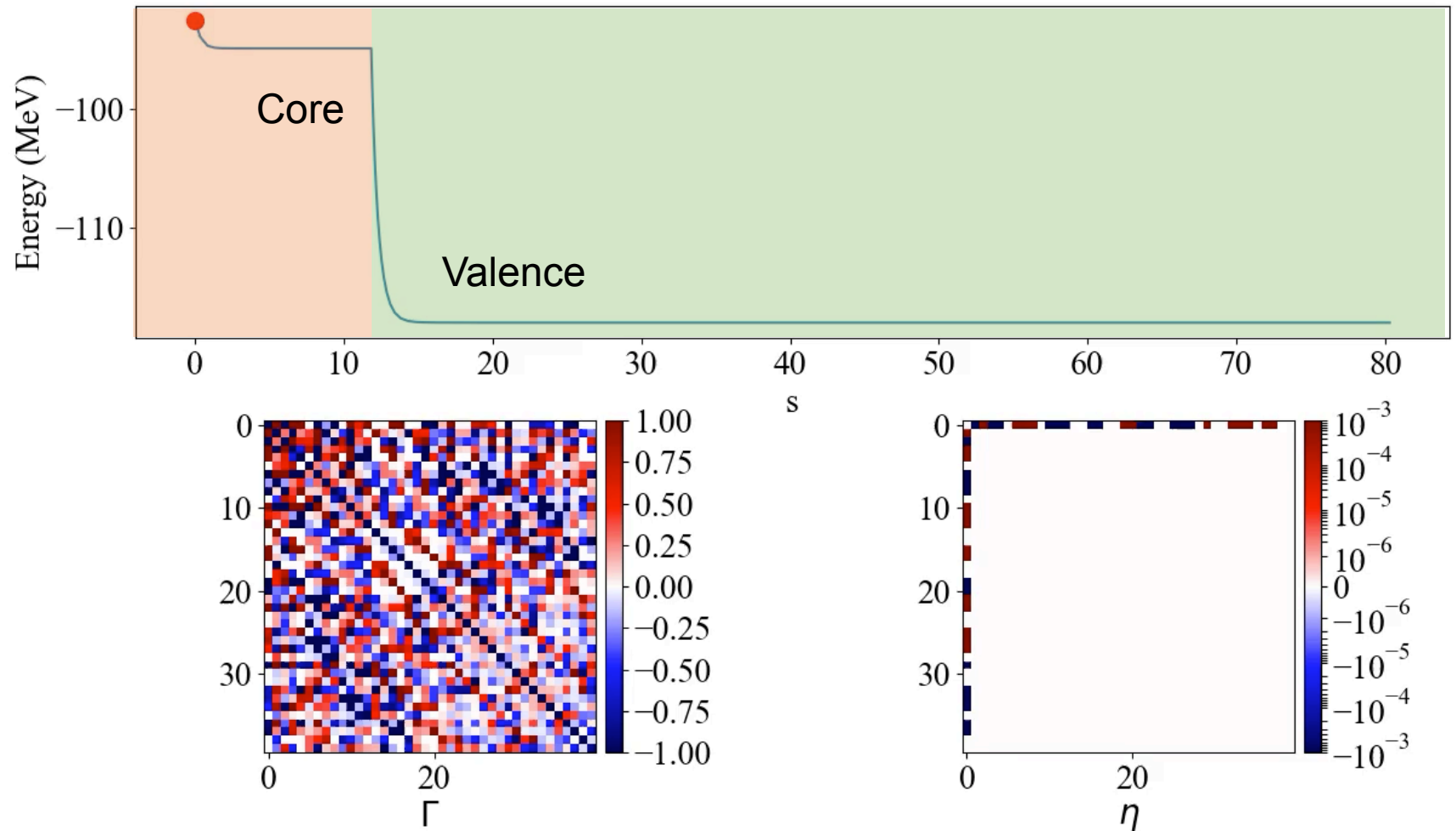
# IM-SRG (Quick reminder)

Evolution for single reference problem ( $^{14}\text{C}$  w/ SRG evolved NN-only)



# IM-SRG (Quick reminder)

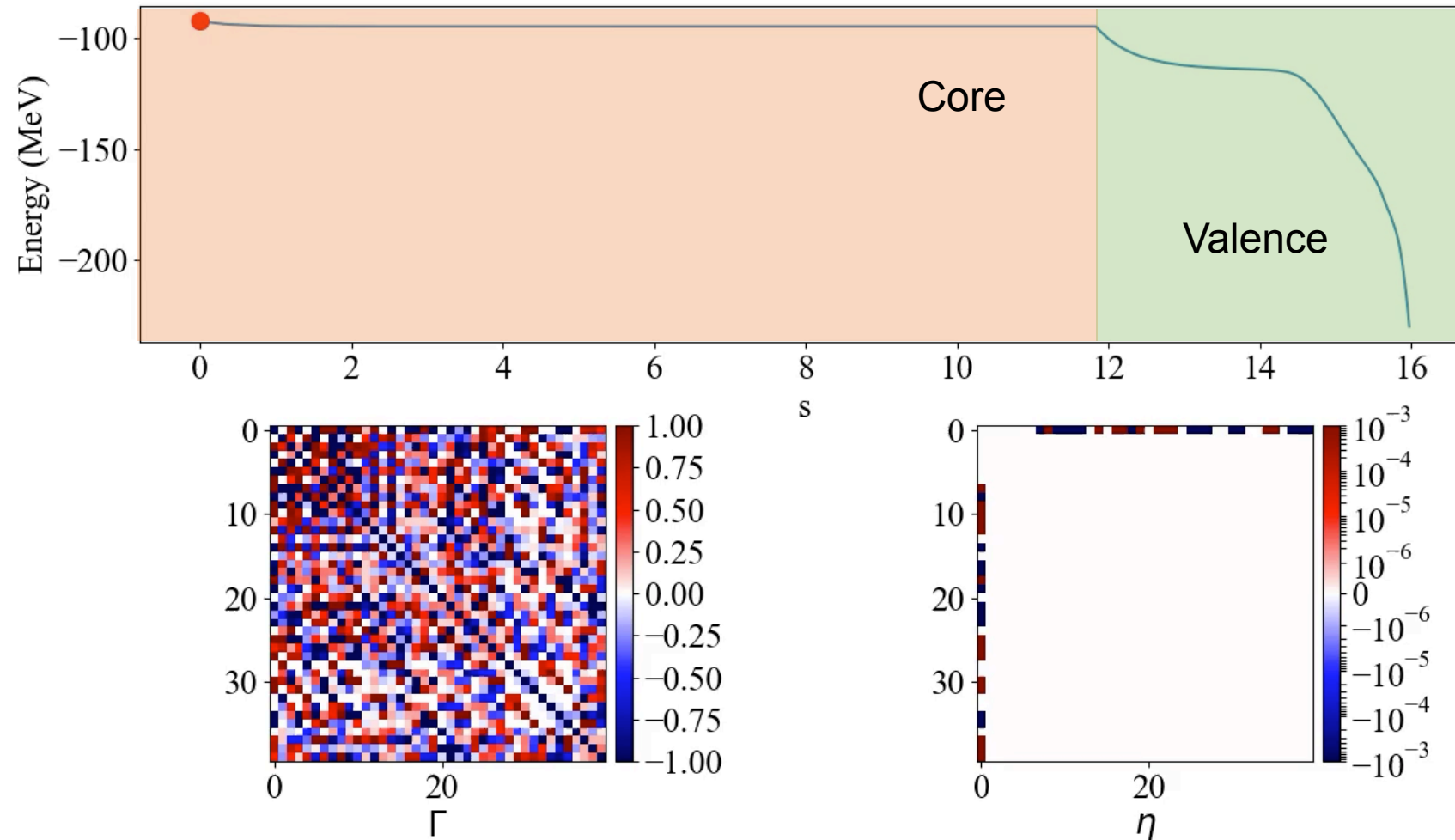
Evolution for p-shell problem ( $^{14}\text{C}$  [ $^4\text{He}$  core] w/ SRG evolved NN-only)





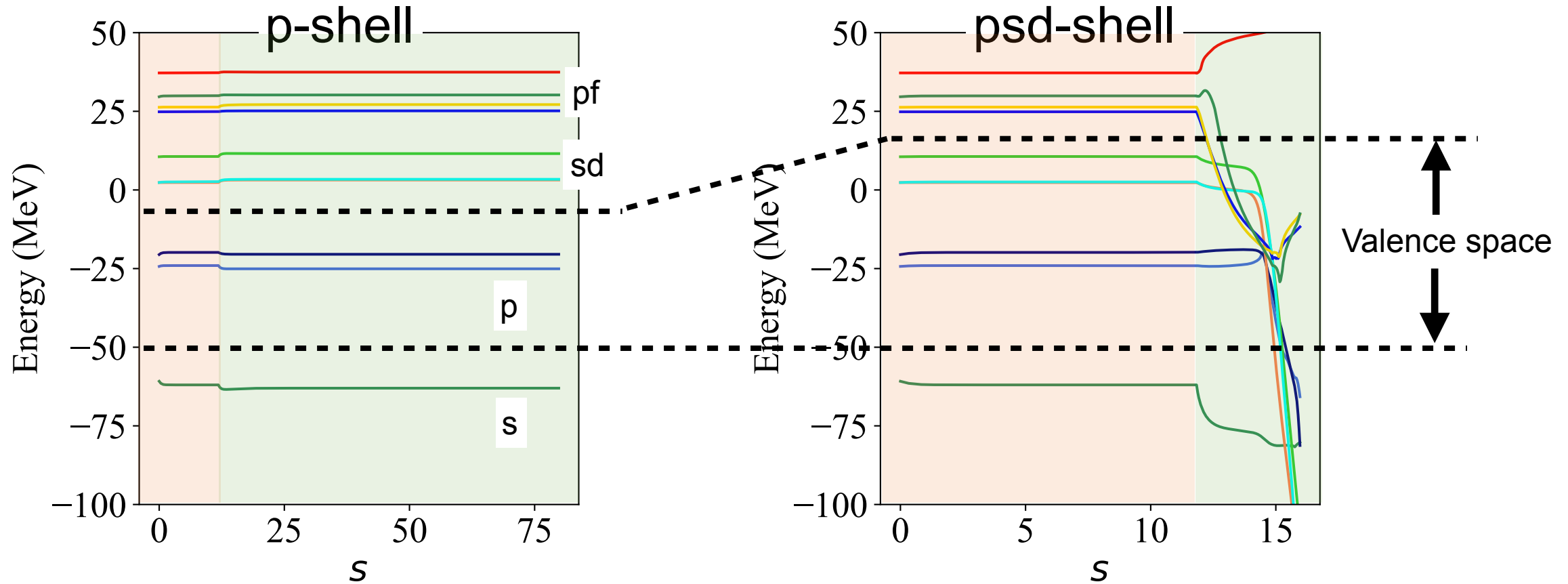
# IM-SRG for two-major shell valence space problem

Evolution for psd-shell problem ( $^{14}\text{C}$  [ $^4\text{He}$  core] w/ SRG evolved NN-only)



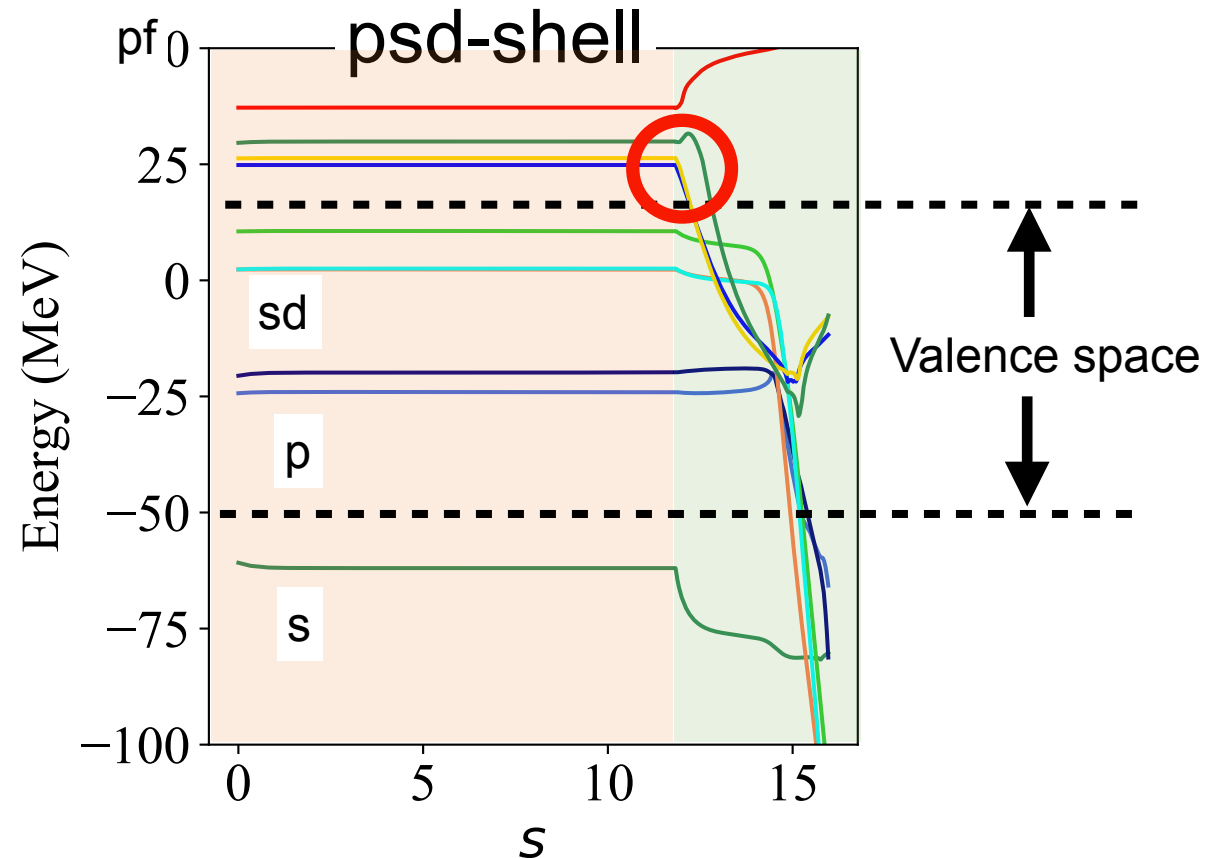
# How it fail

- Flow of single-particle energies



# How it fail

- Flow of single-particle energies
  - ◆ At the very beginning of valence-decoupling flow, some of pf-shell orbits come down.
  - ◆ Intuitively, we expect that P- and Q-space single particle energies do not mix.
  - ◆ At the beginning of the flow, the slope of single-particle energies ( $df/ds$ ) seems to be crucial.



# Flow equation for single-particle energies

- Flow equation for S.P.Es:

$$\begin{aligned} \frac{df_{pp}}{ds} = & 2 \sum_a \eta_{pa} f_{ap} + \frac{1}{2} \sum_{ab} (n_a - n_b) (\eta_{ab} \Gamma_{bpap} - f_{ab} \eta_{bpap}) \\ & + \sum_{abc} (n_a n_b \bar{n}_c + \bar{n}_a \bar{n}_b n_c) \eta_{cpab} \Gamma_{abcp} \end{aligned}$$

- Assuming White generator:

$$\begin{aligned} \eta_{12} &= \frac{f_{12}}{f_{11} - f_{22}} \\ \eta_{1234} &= \frac{\Gamma_{1234}}{f_{11} + f_{22} - f_{33} - f_{44}} \end{aligned}$$

## Flow equation for single-particle energies

- Assuming HF basis:

$$\frac{df_{pp}}{ds} \sim \sum_{p'} \sum_{hh'} \frac{|\Gamma_{hh'pp'}|^2}{f_{pp} + f_{p'p'} - f_{hh} - f_{h'h'}} + \sum_{p'p''} \sum_h \frac{|\Gamma_{php'p''}|^2}{f_{pp} + f_{hh} - f_{p'p'} - f_{p''p''}}$$

- We can modify the denominator so that  $\frac{df_{pp}}{ds} \geq 0$

# Flow equation for single-particle energies

- Modifying the generator

- ◆ Simple way is to give the constant shift to energy denominator

K. Suzuki, Prog. Theor. Phys. **58**, 1064 (1977).

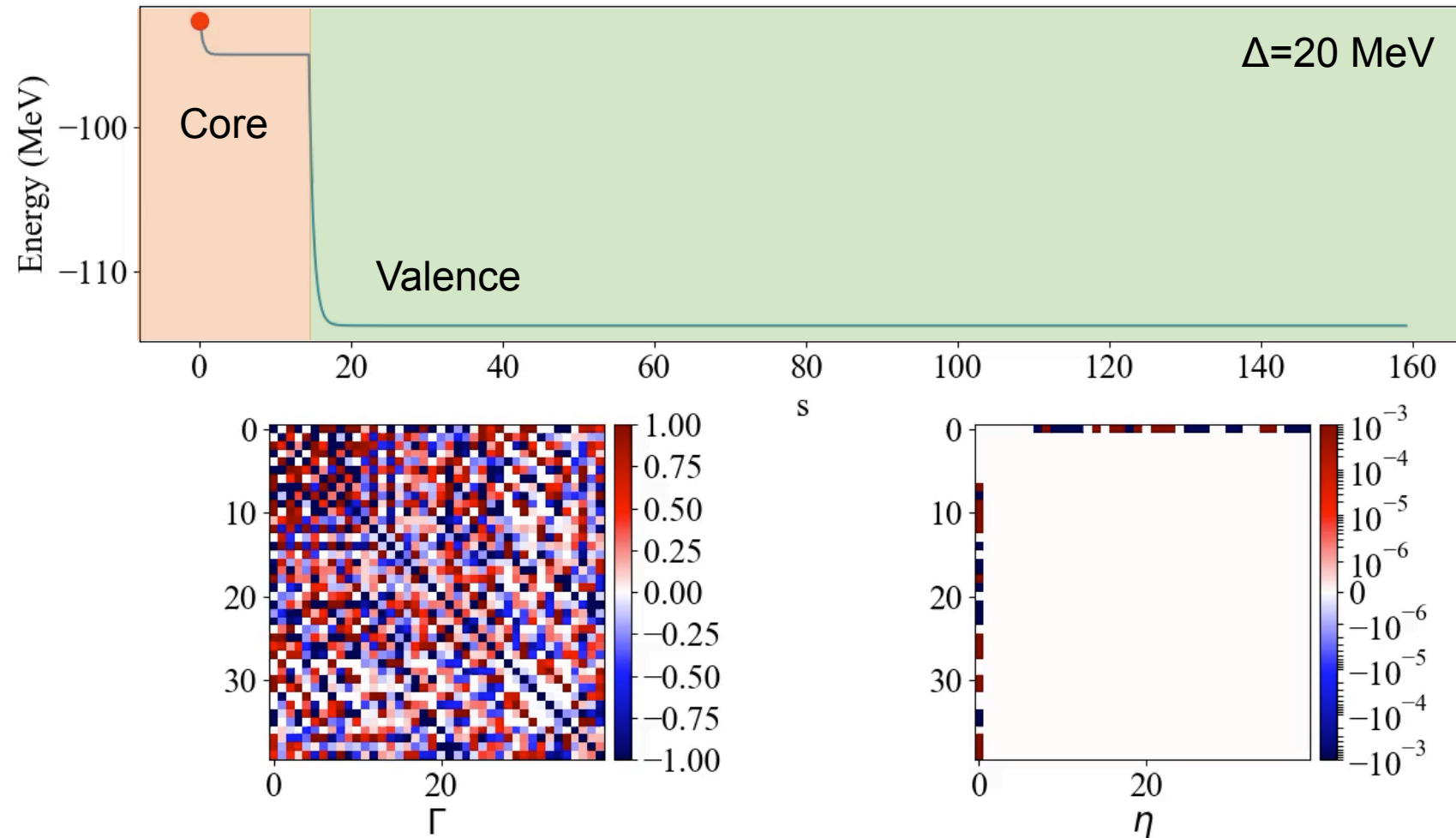
N. Tsunoda, K. Takayanagi, M. Hjorth-Jensen, and T. Otsuka, Phys. Rev. C **89**, 024313 (2014).

$$\begin{array}{ccc}
 \eta_{12} = \frac{f_{12}}{f_{11} - f_{22} + \Gamma_{1212}} & & \eta_{12} = \frac{f_{12}}{f_{11} - f_{22} + \Gamma_{1212} + \Delta} \\
 \eta_{1234} = \frac{\Gamma_{1234}}{f_{11} + f_{22} - f_{33} - f_{44} + A_{1234}} & \longrightarrow & \eta_{1234} = \frac{\Gamma_{1234}}{f_{11} + f_{22} - f_{33} - f_{44} + A_{1234} + \Delta} \\
 A_{1234} = \Gamma_{1212} + \Gamma_{3434} - \Gamma_{1313} - \Gamma_{2424} - \Gamma_{1414} - \Gamma_{2323} & & A_{1234} = \Gamma_{1212} + \Gamma_{3434} - \Gamma_{1313} - \Gamma_{2424} - \Gamma_{1414} - \Gamma_{2323}
 \end{array}$$

- ◆ For our purpose, suitable choice of  $\Delta$  would be comparable to  $\hbar\omega$ .

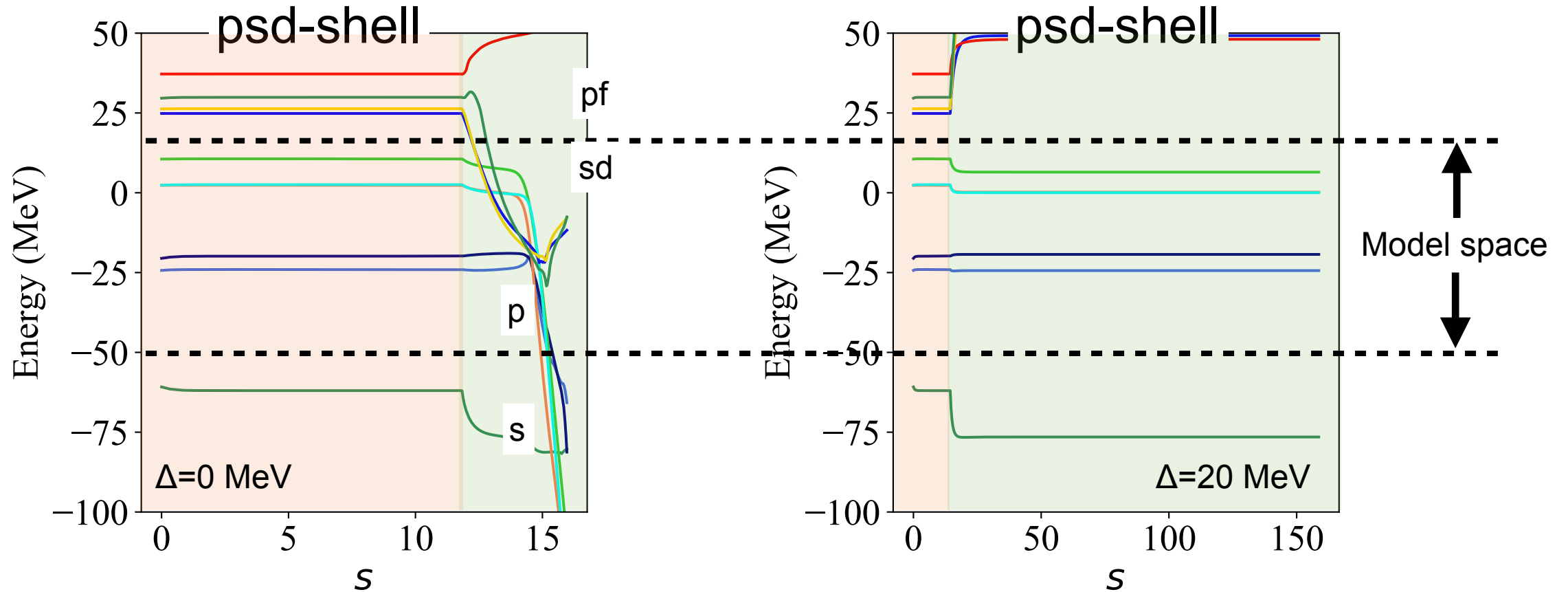
# Does it really work?

Evolution for psd-shell problem ( $^{14}\text{C}$  [ $^4\text{He}$  core] w/ SRG evolved NN-only)



# Does it really work?

- Flow of single-particle energies



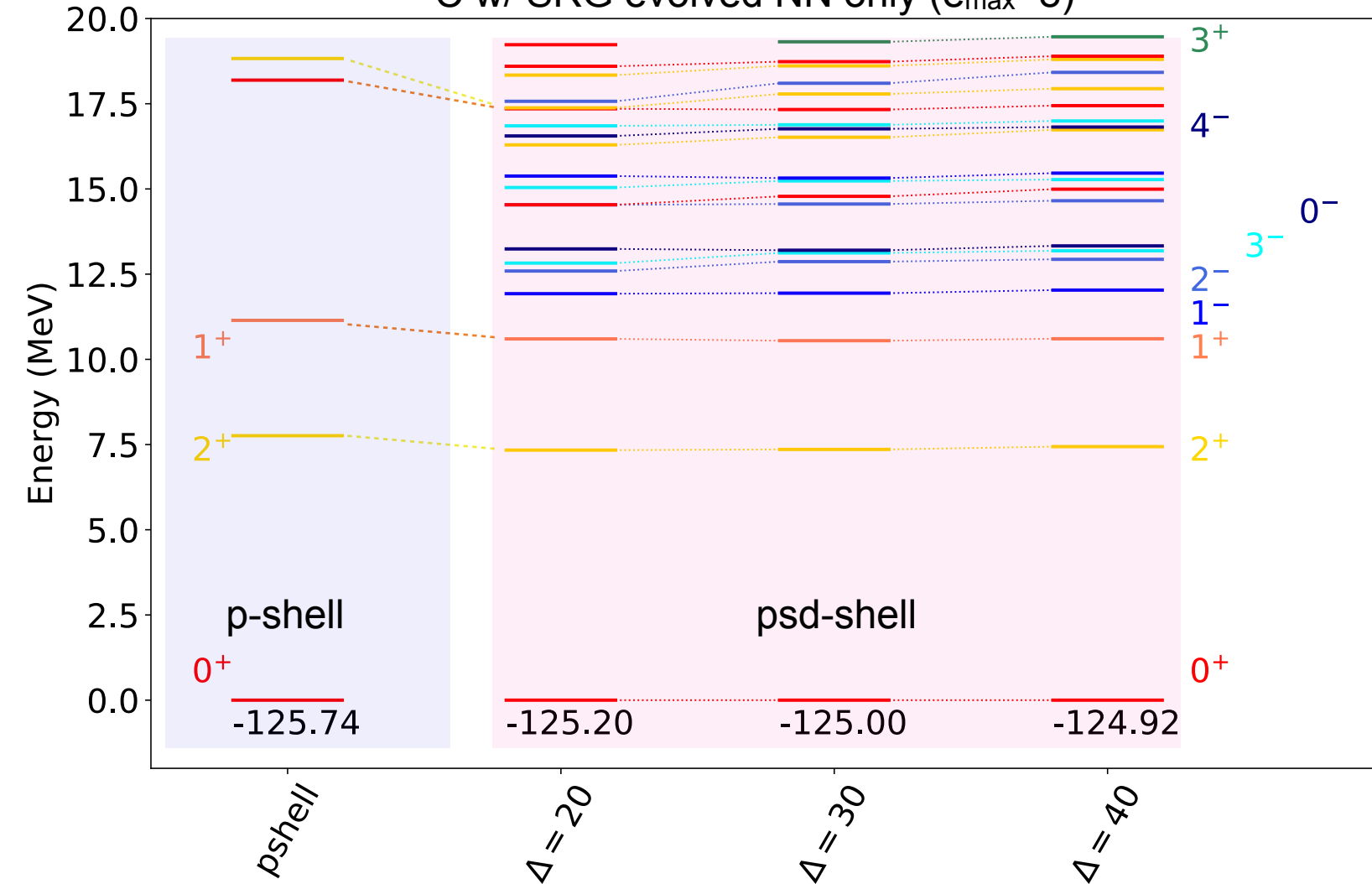


## What is $\Delta$ ?

- Introduction of  $\Delta$  can be regarded as the choice of the generator.
  - ◆ There are some variations:
    - ❖ Wegner
    - ❖ White
    - ❖ Imaginary time
    - ❖ ...
  - ◆ It can affect to the strength of induced many-body terms

# Comparison with single major shell calculation

$^{14}\text{C}$  w/ SRG evolved NN only ( $e_{\text{max}}=8$ )

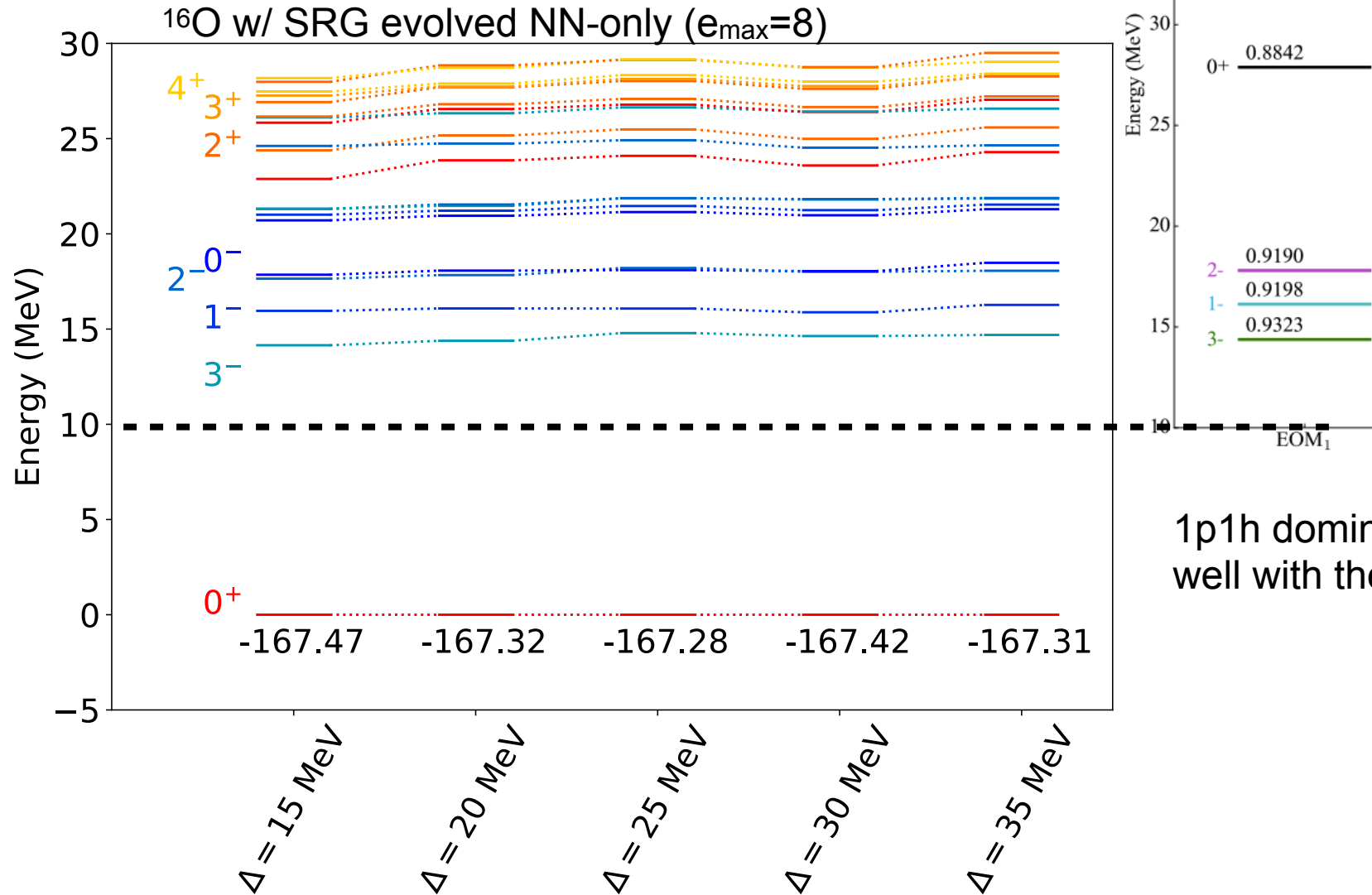


Including negative parity states, many states can appear from psd calculations.

For  $0^+_{1, 2}$ , and  $1^+_{1, 2}$ , p-shell configurations are more than 90%. Also the energies are reasonably close to p-shell results.

$\Delta$  dependence is not unique

# Comparison with EOM method



1p1h dominant lowest 3-, 1-, 2- agree well with the EOM-IMSRG results.

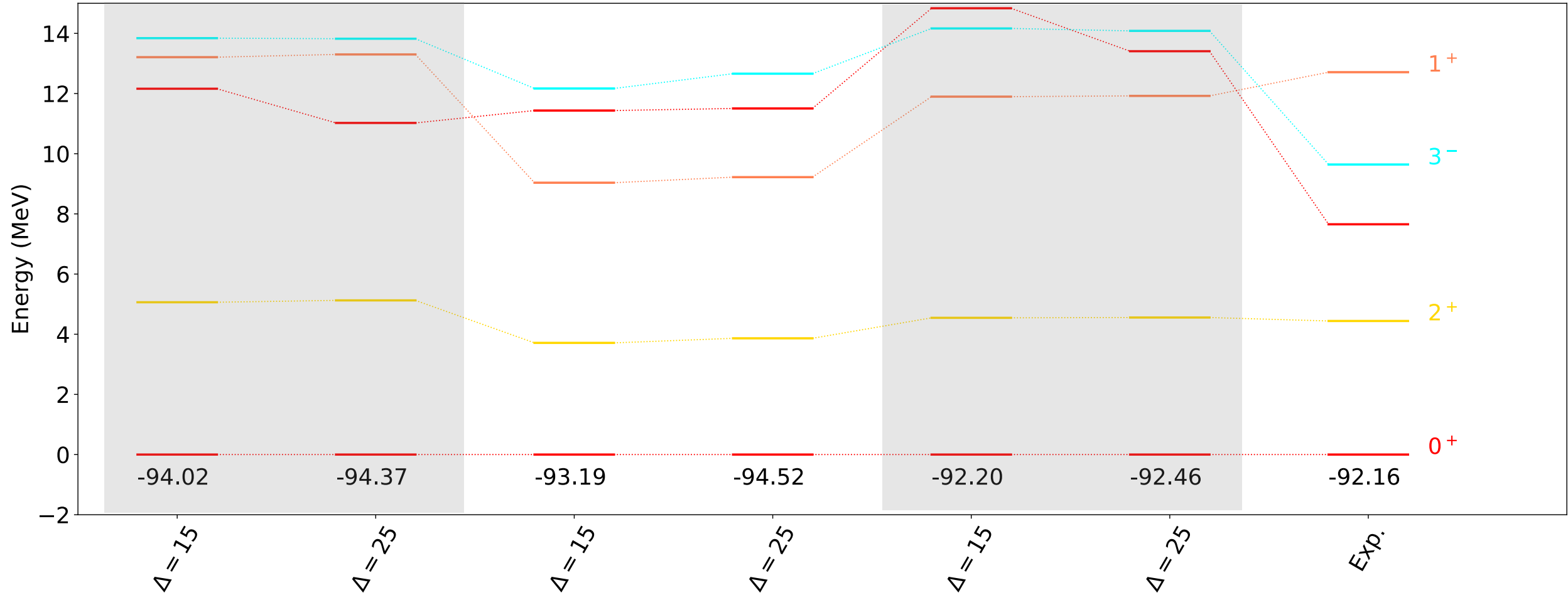
# $^{12}\text{C}$ energies from *psd* calculations

All calculation results are obtained @  $e_{\text{max}}=10$

EM1.8/2.0

$\text{N}^2\text{LO}_{\text{sat}}$

$\text{N}^4\text{LO} + \text{Inl 3NF}$



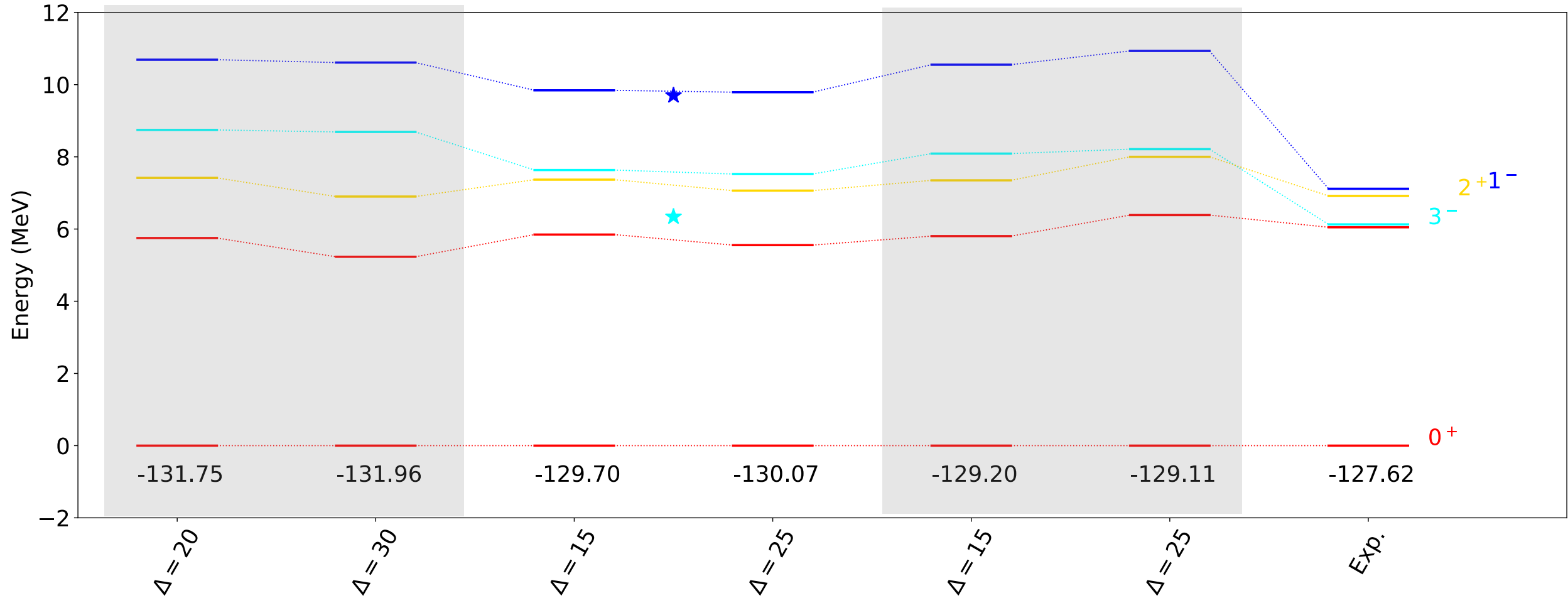
# $^{16}\text{O}$ energies from *psd* calculations

All calculation results are obtained @  $e_{\text{max}}=10$

EM1.8/2.0

$\text{N}^2\text{LO}_{\text{sat}}$

$\text{N}^4\text{LO} + \text{Inl 3NF}$



## Center-of-mass issue

- So far, we added the center-of-mass Hamiltonian at the shell-model calculation stage:

$$H \longrightarrow H_{VS} + \beta H_{cm} \longrightarrow \text{energies}$$

- But,  $H_{VS}$  is no longer represented in HO basis. We should add  $H_{cm}$  from the beginning:

$$H + \beta H_{cm} \longrightarrow H_{VS} \longrightarrow \text{energies}$$

## Summary & Future work

- By adding the shift in energy denominator, the two-major shell-model effective Hamiltonian can be obtained.
- Obtained results are promising
- Treatment of center-of-mass motion
- Comparison with NCSM
- Application to island of inversion