

TRIUMF Theory Workshop (Feb. 26 – Mar. 1, 2019)
Progress in Ab Initio Techniques in Nuclear Physics

Alpha-cluster structure from no-core Monte Carlo shell model

Takashi Abe (U of Tokyo)

Supported by MEXT and JICFuS

Priority Issue 9 to be Tackled by Using Post K Computer “Elucidation of the
Fundamental Laws and Evolution of the Universe”

TRIUMF
March 1, 2019

“Ab initio” in low-energy nuclear structure physics

- Major challenge in nuclear physics
 - Nuclear structure & reactions directly from *ab-initio* calc. w/ nuclear forces
 - *ab-initio* approaches in nuclear structure calculations ($A > 4$):
Light mass: Green's Function Monte Carlo, No-Core Shell Model ($A \sim 12$),
Medium/heavy mass: Coupled Cluster, IM-SRG,
Self-consistent Green's Function theory, Lattice EFT, UMOA, ...
- Solve the non-relativistic many-body Schroedinger eq.
and obtain the eigenvalues and eigenvectors.

$$H|\Psi\rangle = E|\Psi\rangle$$

$$H = T + V_{\text{NN}} + V_{\text{3N}} + \dots + V_{\text{Coulomb}}$$

- **Ab initio**: All nucleons are active, and Hamiltonian consists of realistic NN (+ 3N + ...) potentials.

→ Computationally demanding → Monte Carlo shell model (MCSM)

Monte Carlo shell model (MCSM)

Standard shell model

$$H = \begin{pmatrix} * & * & * & * & * & \cdots \\ * & * & * & * & & \\ * & * & * & & & \\ * & * & & \ddots & & \\ * & & & & & \\ \vdots & & & & & \end{pmatrix} \xrightarrow{\text{Diagonalization}} \begin{pmatrix} E_0 & & & & & 0 \\ & E_1 & & & & \\ & & E_2 & & & \\ & & & \ddots & & \\ 0 & & & & & \end{pmatrix}$$

Large sparse matrix
 $\sim \mathcal{O}(10^{10})$ # non-zero MEs
 $\sim \mathcal{O}(10^{13-14})$

- Importance truncation

Monte Carlo shell model

$$H \sim \begin{pmatrix} * & * & \cdots \\ * & \ddots & \\ \vdots & & \end{pmatrix} \xrightarrow{\text{Diagonalization}} \begin{pmatrix} E'_0 & & 0 \\ & E'_1 & \\ 0 & & \ddots \end{pmatrix}$$

Important bases stochastically selected $\sim \mathcal{O}(100)$

$$|\Psi(J, M, \pi)\rangle = \sum_i^{N_{basis}} f_i |\Phi_i(J, M, \pi)\rangle$$

$$|\Phi(J, M, \pi)\rangle = \sum_K g_K P_{MK}^J P^\pi |\phi\rangle$$

diagonalization

Deformed

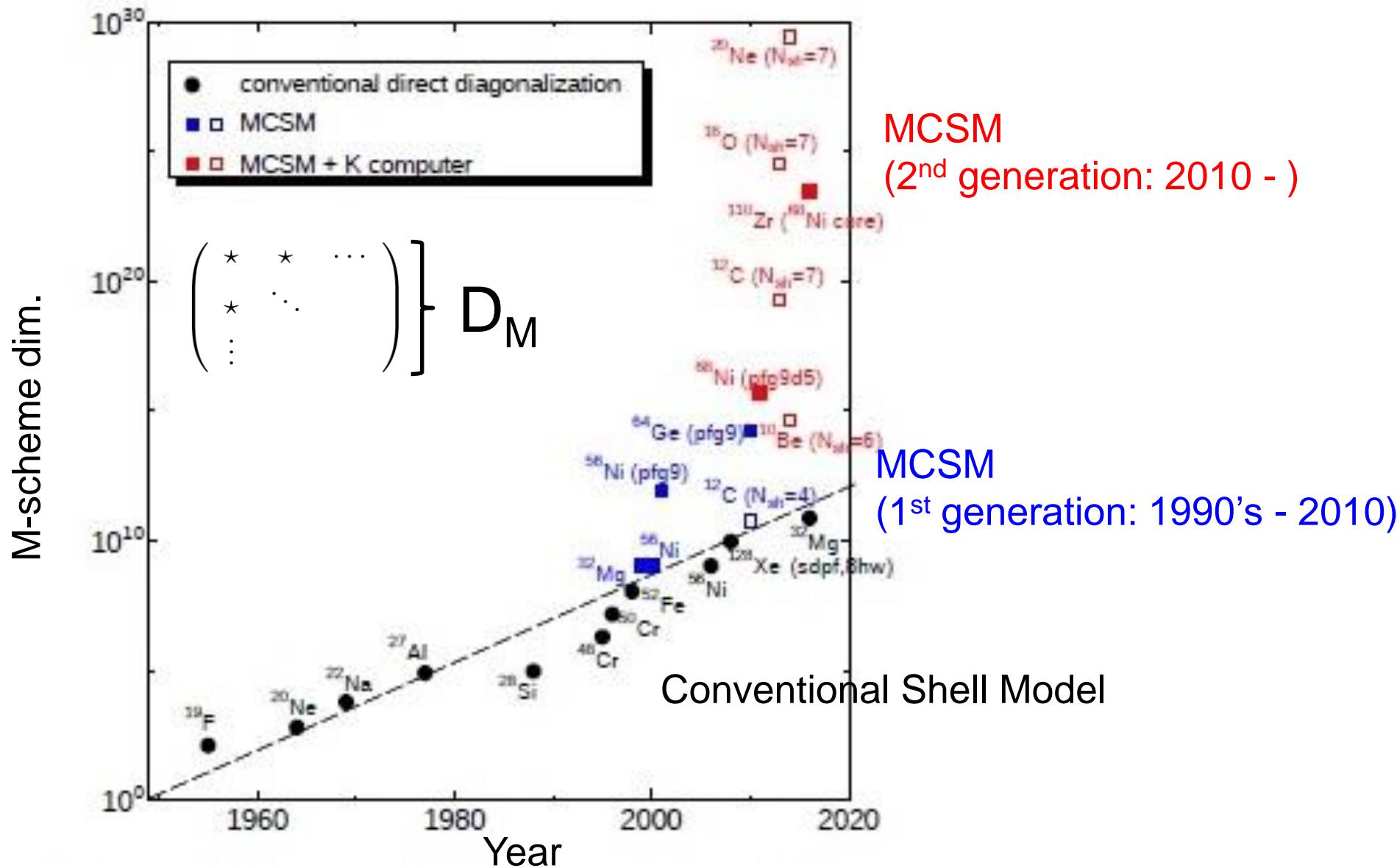
$$|\phi\rangle = \prod_i^A a_i^\dagger |-\rangle$$

Spherical

$$a_i^\dagger = \sum_\alpha c_\alpha^\dagger D_{\alpha i}$$

stochastic sampling & CG method

Historical evolution/development of the MCSM



How to obtain ab-initio results from no-core MCSM

- Two steps of the extrapolation

↙ Same as in the MCSM w/ an inert core

1. Extrapolation of our MCSM (approx.) results to exact results in the fixed size of model space

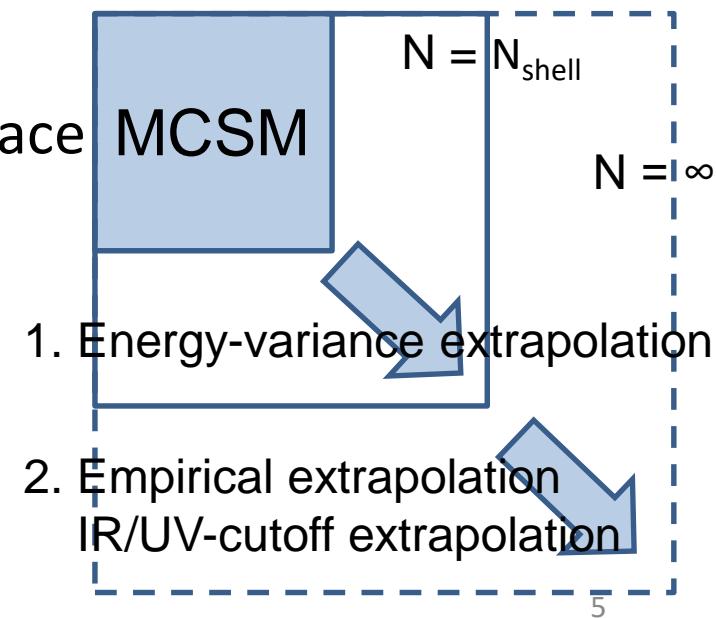
Energy-variance extrapolation

N. Shimizu, Y. Utsuno, T. Mizusaki, T. Otsuka, T. Abe, & M. Honma, Phys. Rev. C82, 061305(R) (2010)

2. Extrapolation into the infinite model space

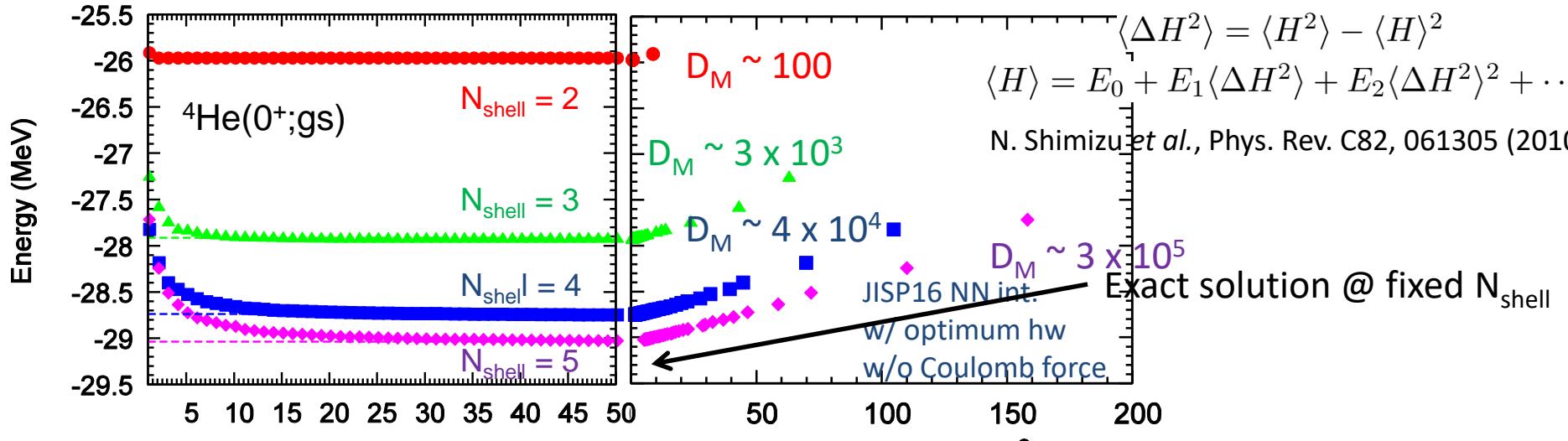
- Empirical extrapolation w.r.t. N_{shell}
- IR- & UV-cutoff extrapolations

→ Ab initio solution

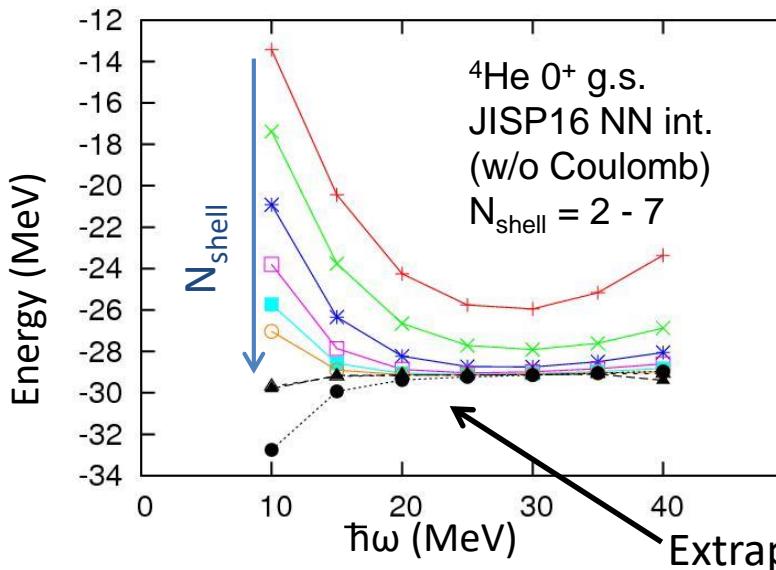


Extrapolations

- Extrapolation to FCI results (@ fixed size of basis space) <- Energy variance



- Extrapolation to full ab initio solution (@ infinite size of basis space)

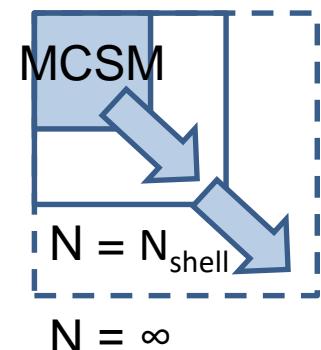


$$(N_{\text{shell}}, \hbar\omega)$$

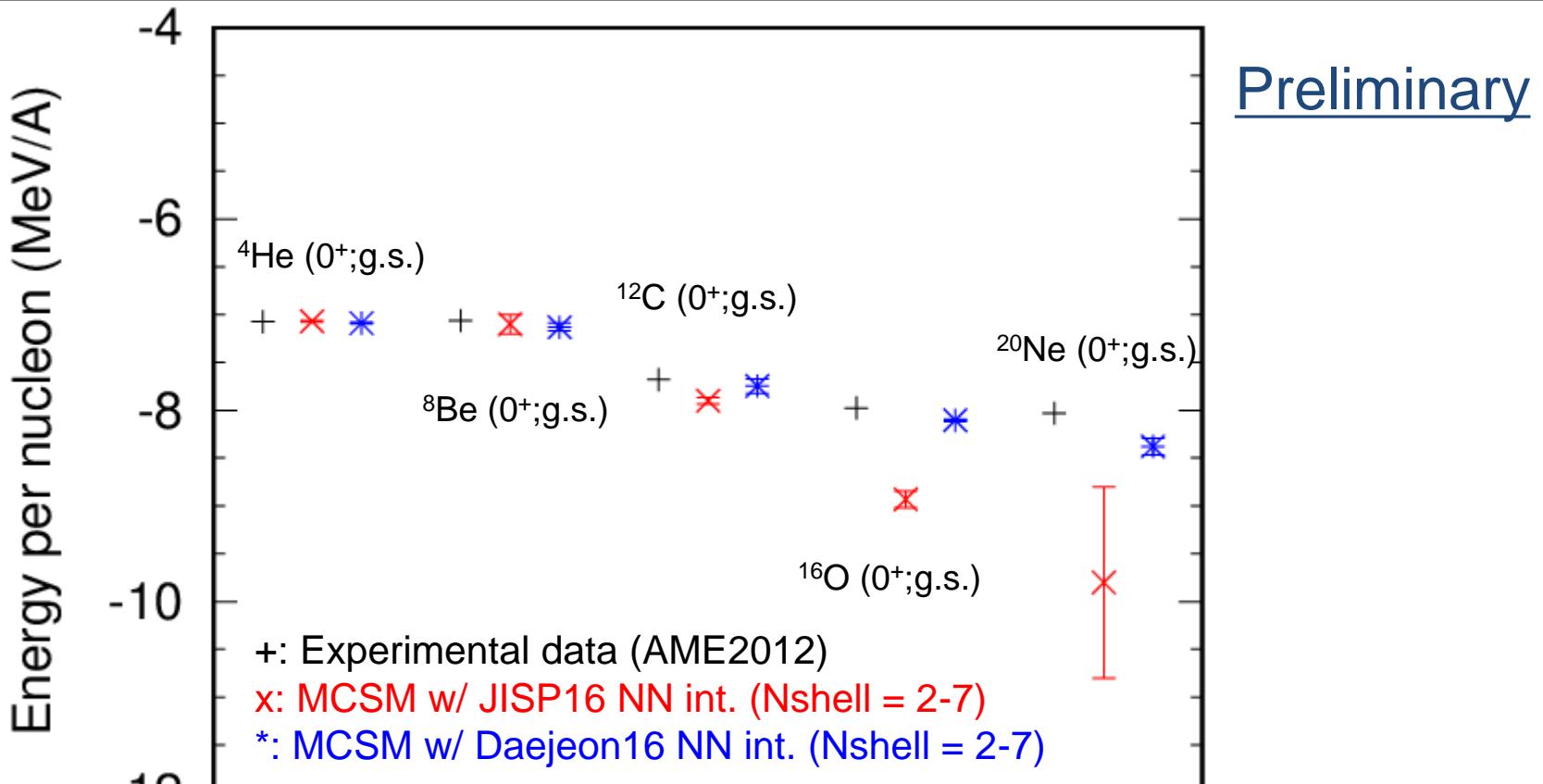
$$E(N) = E(N = \infty) + a \exp(-bN)$$

MCSM(traditional): -29.15(3) MeV
($N_{\text{shell}} = 3 - 7$, $\hbar\omega = 20 - 30$ MeV)

c.f.) NCFC: -29.164(2) MeV
Extrapolated results to infinite N_{max} space



Comparison of MCSM results w/ experiments

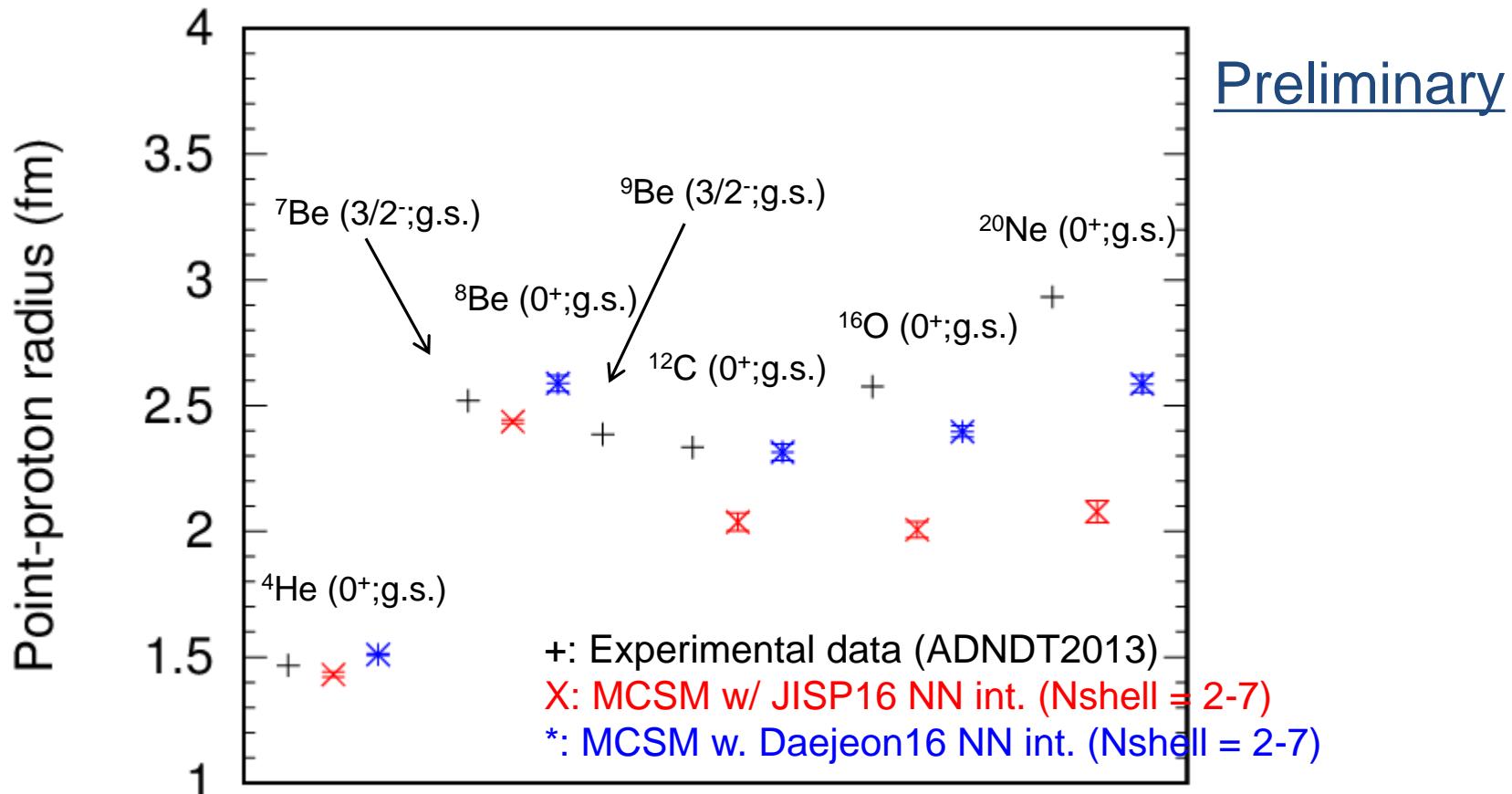


MCSM results are obtained using K computer by traditional extrapolation w/ optimum harmonic oscillator energies.

JISP16 results show good agreements w/ experimental data up to ${}^{12}\text{C}$, slightly overbound for ${}^{16}\text{O}$, and clearly overbound for ${}^{20}\text{Ne}$.

Daejeon16 results show good agreements w/ experimental data up to ${}^{20}\text{Ne}$.

Comparison of MCSM results w/ experiments



MCSM results are obtained using K computer around optimum harmonic oscillator energies for radii.

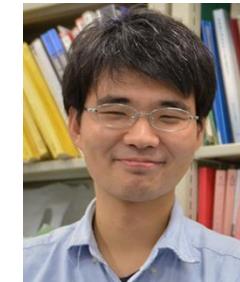
JISP16 results show good agreements w/ experimental data up to ${}^8\text{Be}$, clearly smaller for heavier nuclei beyond ${}^{12}\text{C}$ as A increases.

Daejeon16 results show larger radii than JISP16 ones.

G.S. energies & excitation spectra of Be isotopes

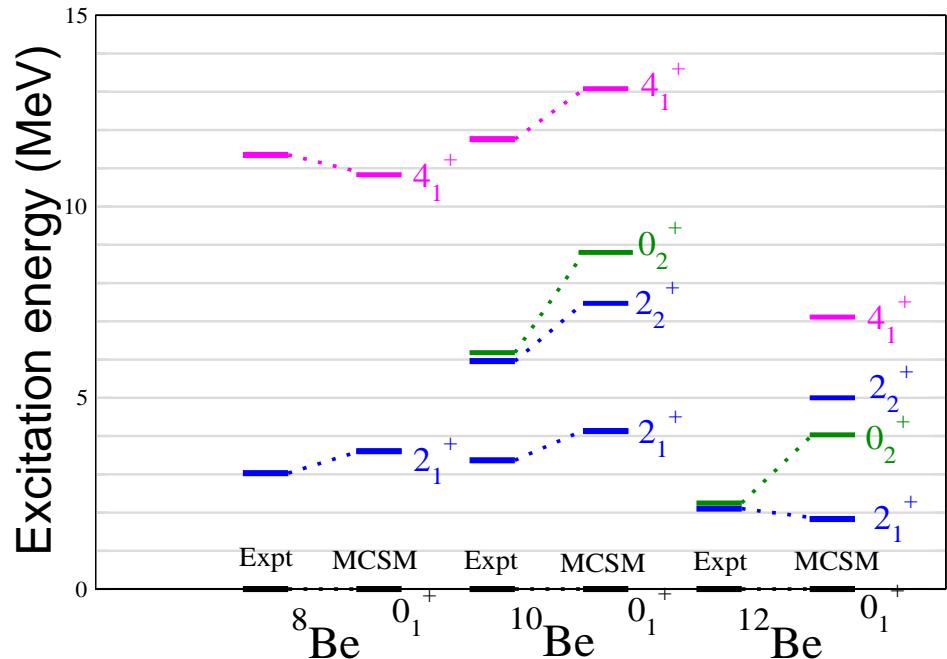
- Ground-state energies

	MCSM (MeV)	Expt (MeV)
${}^8\text{Be} (0^+)$	-49.95	-56.499
${}^{10}\text{Be} (0^+)$	-53.4	-64.98
${}^{12}\text{Be} (0^+)$	-49.93	-68.65



T. Yoshida (RIST)

- Excitation energies



JISP16 NN

$N_{\text{shell}} = 6, \hbar w = 15 \text{ MeV}$

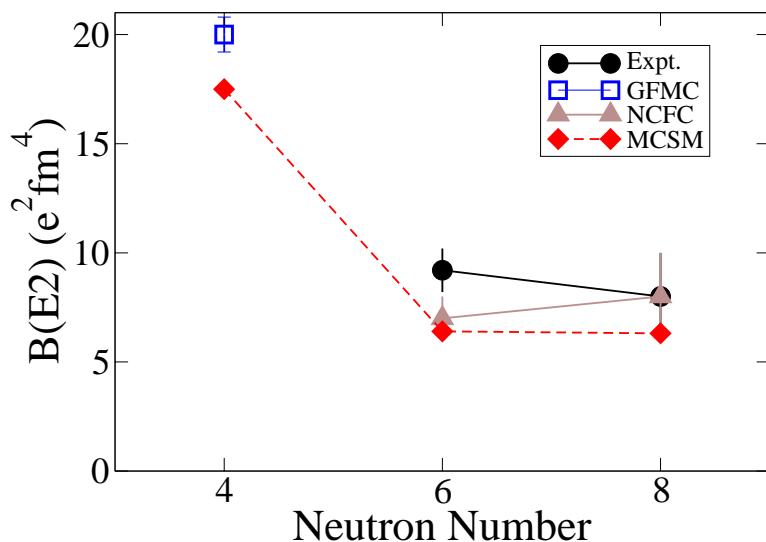
Expt.: ${}^8\text{Be}, {}^{10}\text{Be}$ (Tilley et al., 2004),
 ${}^{12}\text{Be}$ (Shimoura et al., 2003)

MCSM: JISP16 NN int., $N_{\text{shell}} = 6, \hbar w = 15 \text{ MeV}$

E2 & E0 transition strengths of Be isotopes

- E2 transition strengths

$$B(E2; 2^+_1 \rightarrow 0^+_1)$$

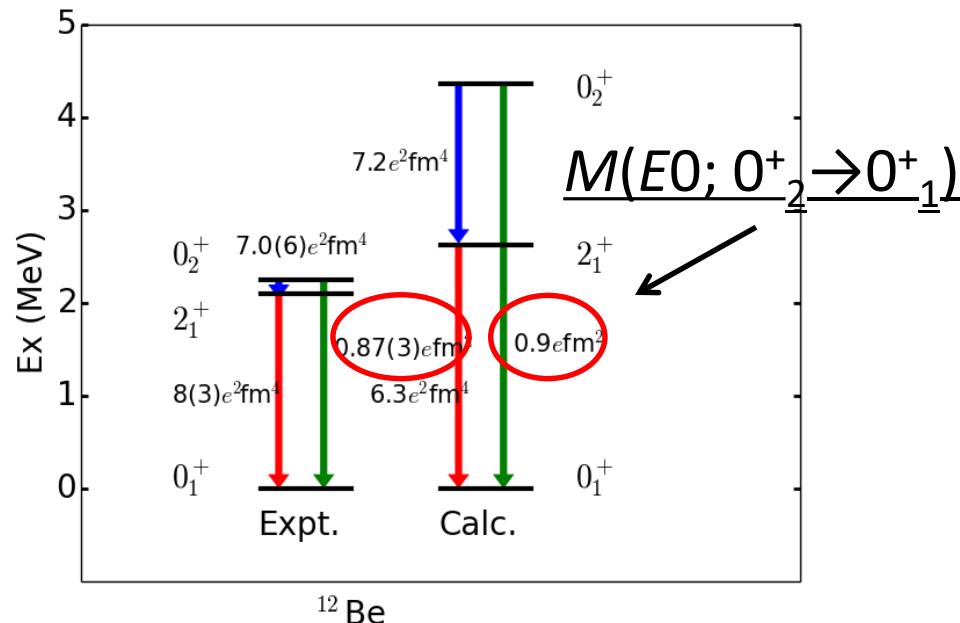


Expt.: ^{8}Be Datar *et al.* 2013

^{10}Be McCutchan *et al.* 2009

^{12}Be Imai *et al.* 2009

- E2 & E0 transition strengths in ^{12}Be



Expt.:

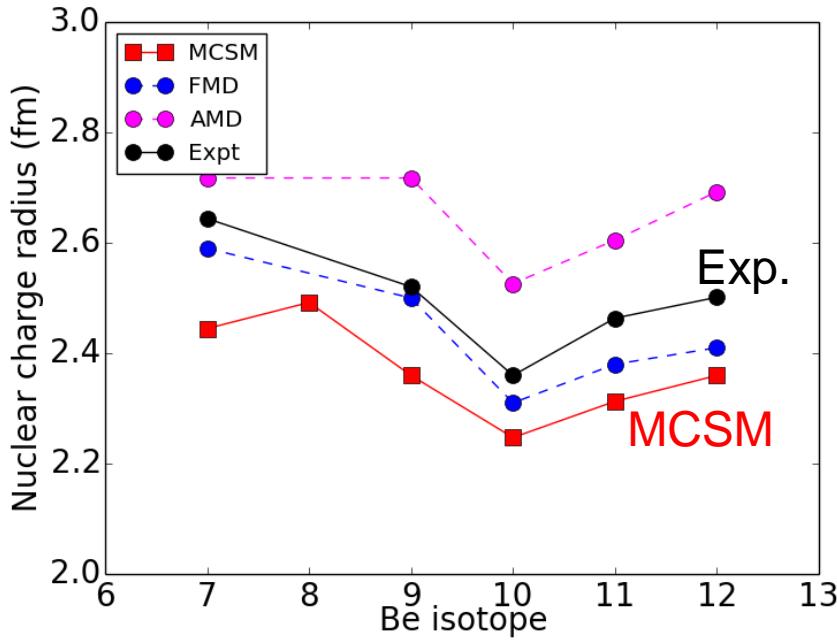
S. Shimoura, et al., Phys. Lett. B 654 87 (2007)

N. Imai, et al., Phys. Lett. B 673 179 (2009)

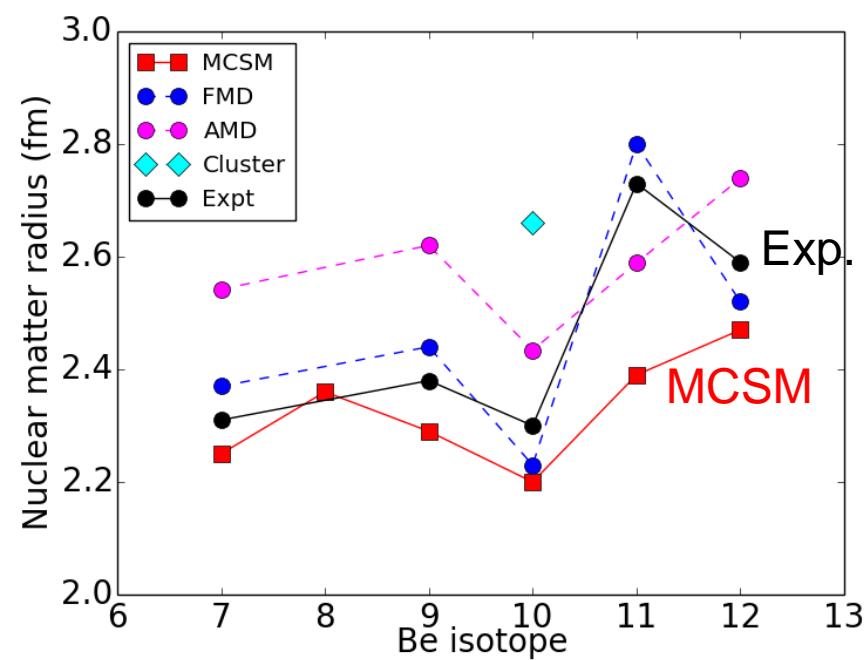
MCSM: JISP16 NN, $N_{\text{shell}} = 6$, $\hbar\omega = 15$ MeV

Radii of Be isotopes

Point-proton radius



Matter radius



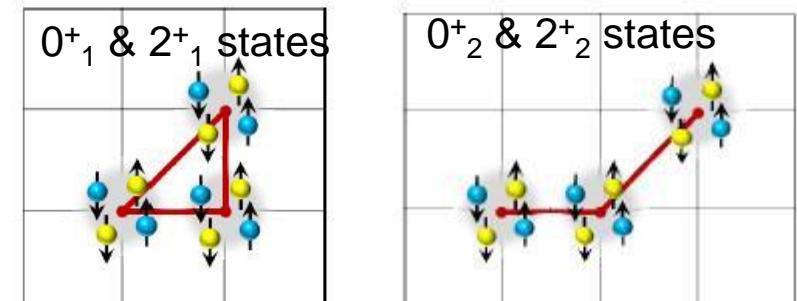
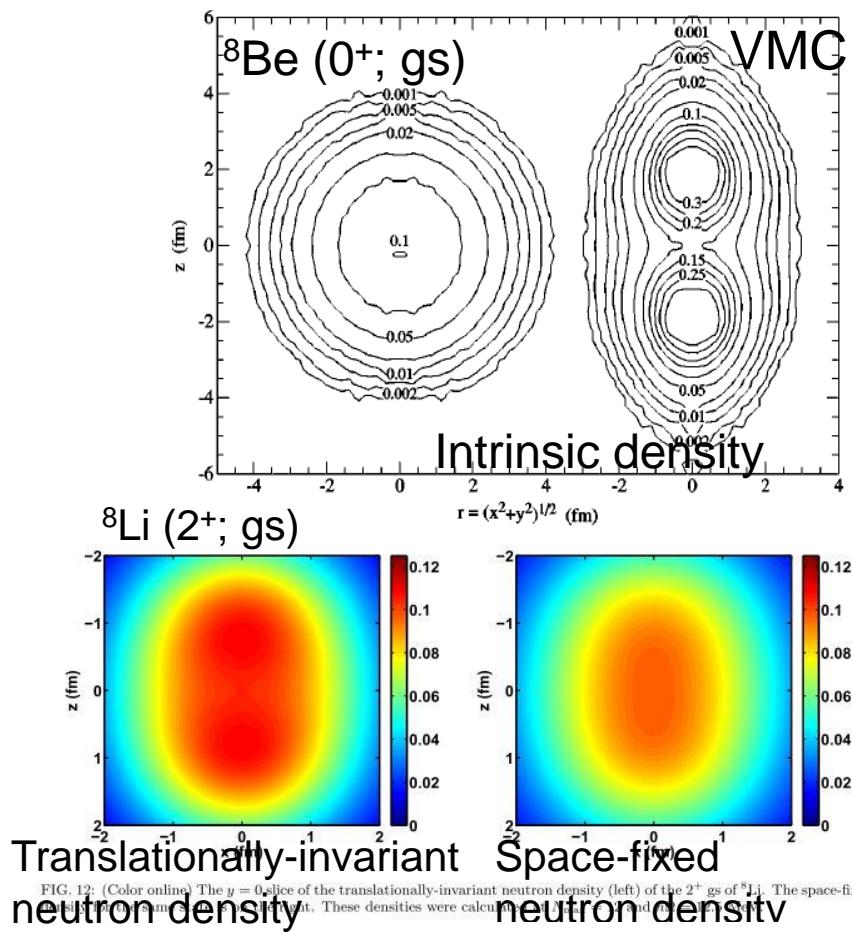
Rather good agreement w/ overall trend, except for ^{11}Be neutron halo

MCSM < Expt., FMD < AMD, Cluster model

- Expt., FMD: F. Ajzenberg-Selove, NPA 506, 1 (1990), A. Krieger et al., PRL 108, 142501 (2012)
AMD : Y. Kanada-En'yo, PRC91, 014315 (2015)
Cluster: M. Ito & K. Ikeda, Rep. Prog. Phys. 77, 096301 (2014)

Density distribution from ab initio calc.

- Green's function Monte Carlo (GFMC)
 - "Intrinsic" density is constructed by aligning the moment of inertia among samples
R. B. Wiringa, S. C. Pieper, J. Carlson, & V. R. Pandharipande, Phys. Rev. C62, 014001 (2000)
- No-core full configuration (NCFC)
 - Translationally-invariant density is obtained by deconvoluting the intrinsic & CM w.f.
C. Cockrell J. P. Vary & P. Maris, Phys. Rev. C86, 034325 (2012)
- Lattice EFT
 - Triangle structure of carbon-12
E. Epelbaum, H. Krebs, T. A. Lahde, D. Lee, & U.-G. Meissner, Phys. Rev. Lett. 109, 252501 (2012), ...
- FMD
 - H. Feldmeier, Nucl. Phys. A515, 147 (1990), ...



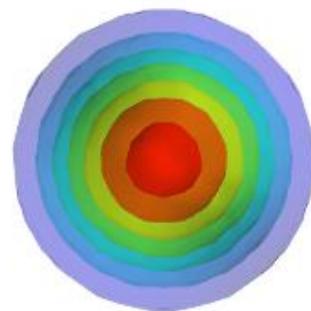
Density distribution in MCSM

$$|\Phi\rangle = \sum_{i=1}^{N_{basis}} c_i |\Phi_i\rangle = c_1 \begin{array}{c} \text{image} \\ \text{of a density} \\ \text{distribution} \end{array} + c_2 \begin{array}{c} \text{image} \\ \text{of a density} \\ \text{distribution} \end{array} + c_3 \begin{array}{c} \text{image} \\ \text{of a density} \\ \text{distribution} \end{array} + c_4 \begin{array}{c} \text{image} \\ \text{of a density} \\ \text{distribution} \end{array} + \dots$$

Angular-momentum projection

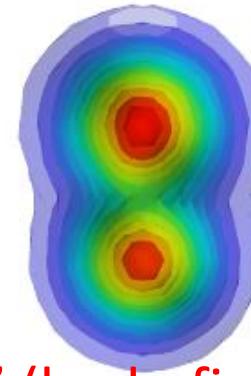
$$|\Psi\rangle = \sum_{i=1}^{N_{basis}} c_i P^J P^\pi |\Phi_i\rangle$$

A way to construct
an “intrinsic” density



${}^8\text{Be}$ 0⁺ ground state

Laboratory frame



“Intrinsic” (body-fixed) frame

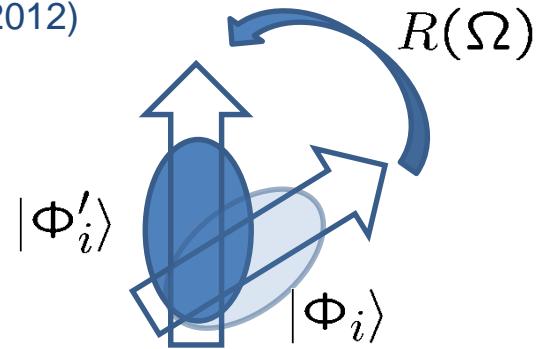
Densities in lab. & body-fixed frames can be constructed by MCSM

How to construct an “intrinsic” density from MCSM w.f.

N. Shimizu, T. Abe, Y. Tsunoda, Y. Utsuno, **T. Yoshida**, T. Mizusaki, M. Honma, T. Otsuka,
Progress in Theoretical and Experimental Physics, 01A205 (2012)

- MCSM wave function

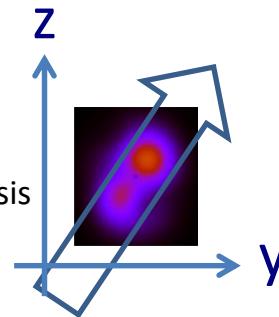
$$|\Psi\rangle = \sum_{i=1}^{N_{basis}} c_i P^J P^\pi |\Phi_i\rangle$$



- Wave function w/o the projections

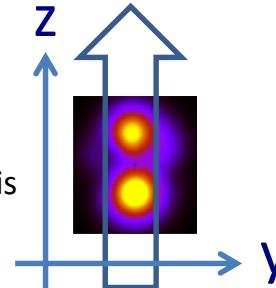
$$\sum_{i=1}^{N_{basis}} c_i |\Phi_i\rangle = c_1 \begin{matrix} \text{image} \\ \downarrow \end{matrix} + c_2 \begin{matrix} \text{image} \\ \downarrow \end{matrix} + \dots + c_{N_{basis}} \begin{matrix} \text{image} \\ \downarrow \end{matrix}$$

Rotation by diagonalizing Q-moment
($Q_{zz} > Q_{yy} > Q_{xx}$)

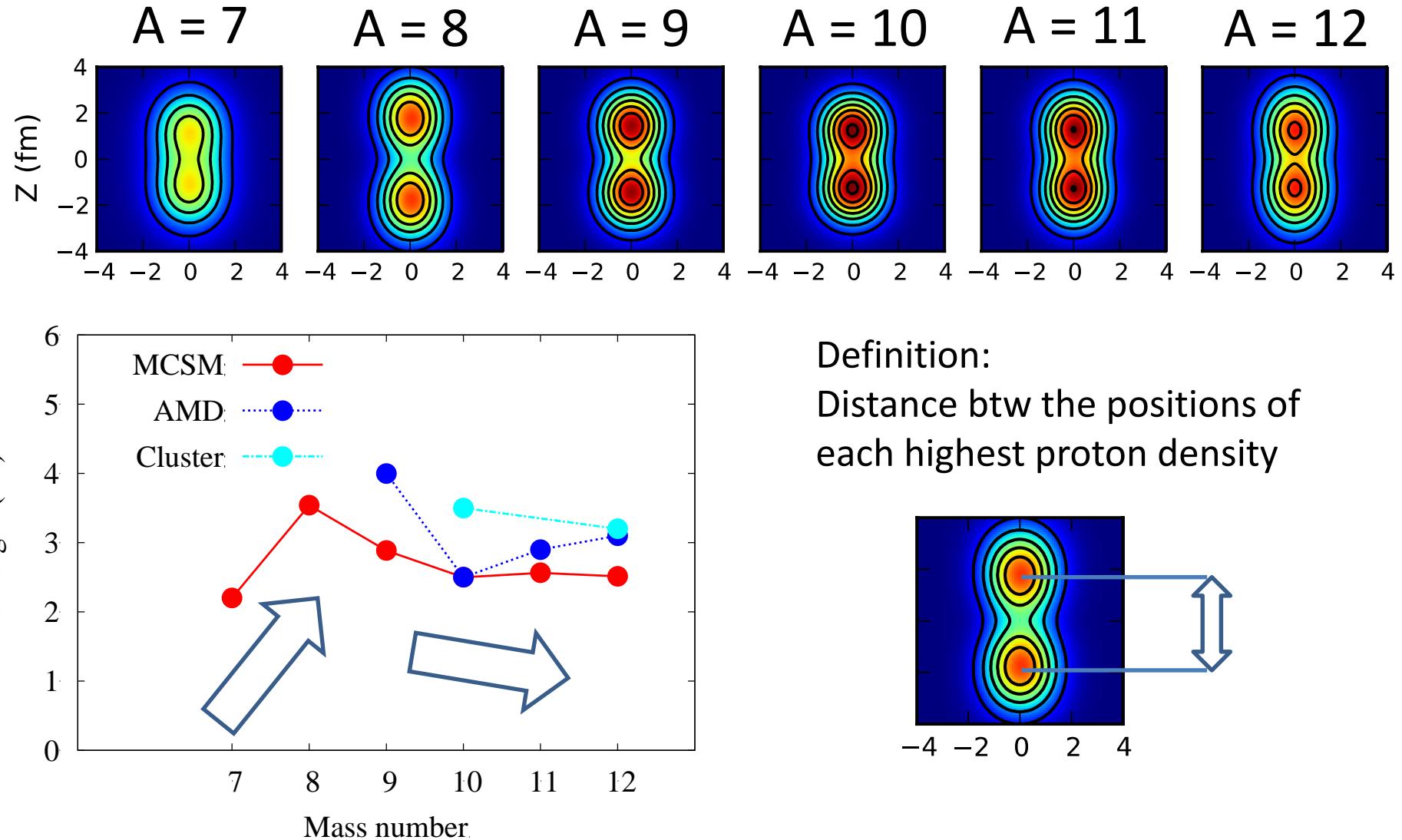


- Wave function w/o the projection w/ the alignment of Q-moment

$$\sum_{i=1}^{N_{basis}} c_i |\Phi'_i\rangle = c_1 \begin{matrix} \text{image} \\ \downarrow \end{matrix} + c_2 \begin{matrix} \text{image} \\ \downarrow \end{matrix} + \dots + c_{N_{basis}} \begin{matrix} \text{image} \\ \downarrow \end{matrix}$$



inter α -cluster distance

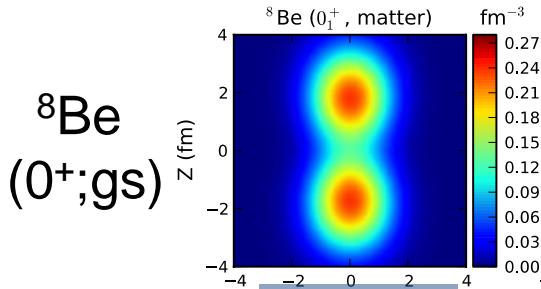


AMD: Y. Kanada-En'yo, Phys. Rev C68, 014319 (2003)

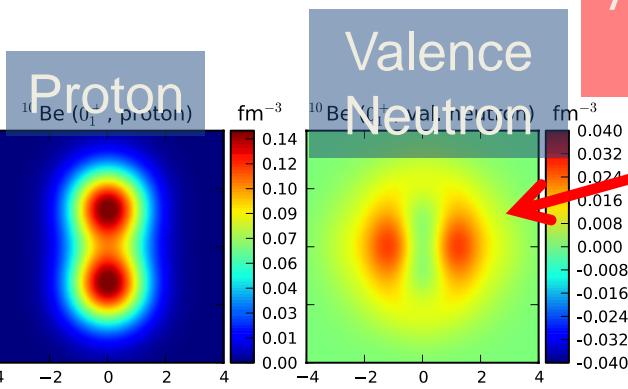
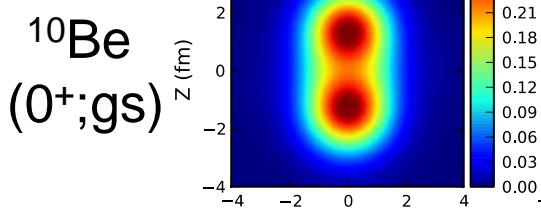
Cluster: M. Ito & K. Ikeda, Rep. Prog. Phys. 77, 096301 (2014)

Density distribution of Be isotopes

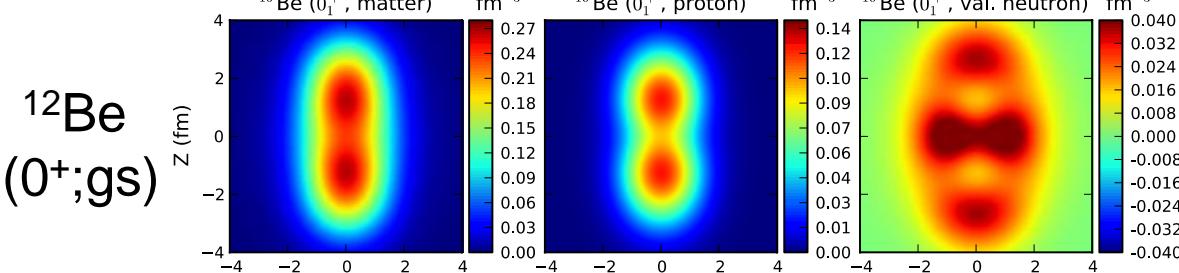
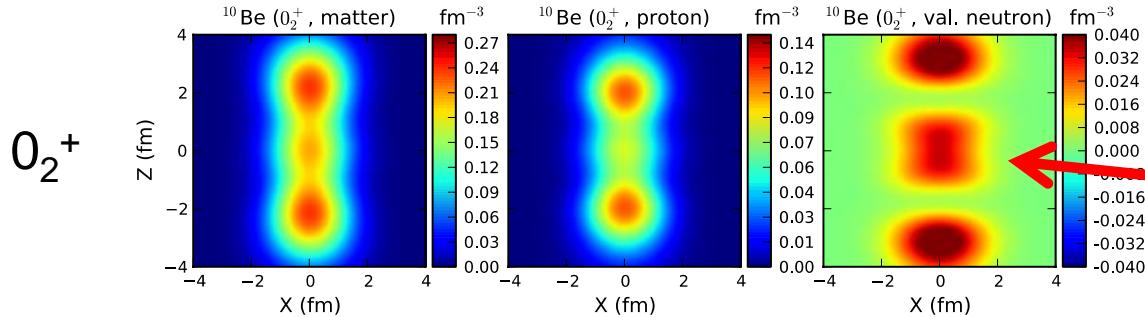
Fading 2- α structure as N increases



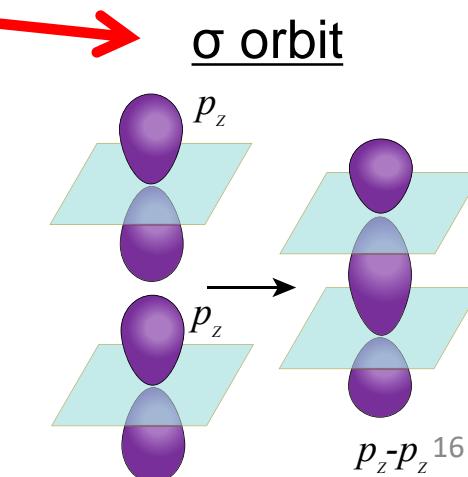
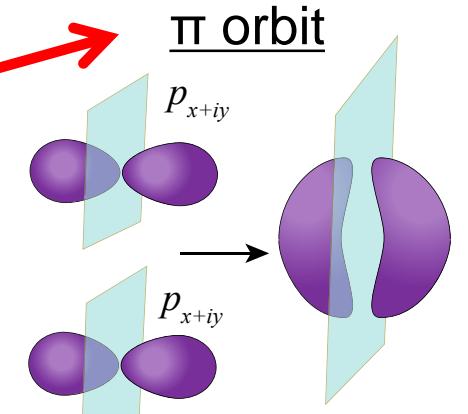
Emergence of
2- α -cluster structure



Appearance of molecular-
orbital structure

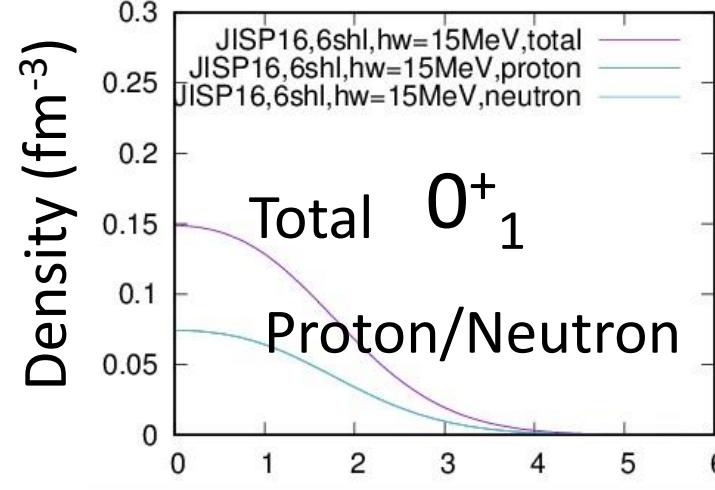


Preliminary

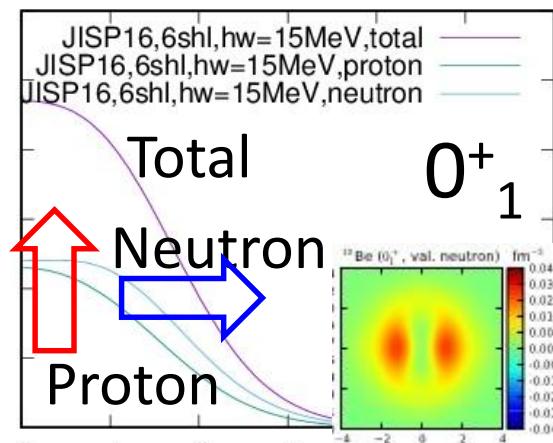


Radial distributions of Be isotopes

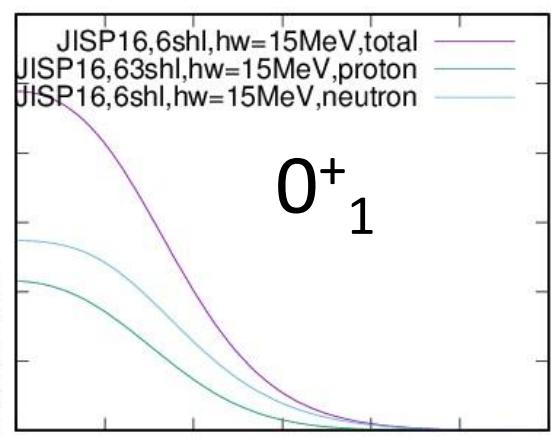
^8Be



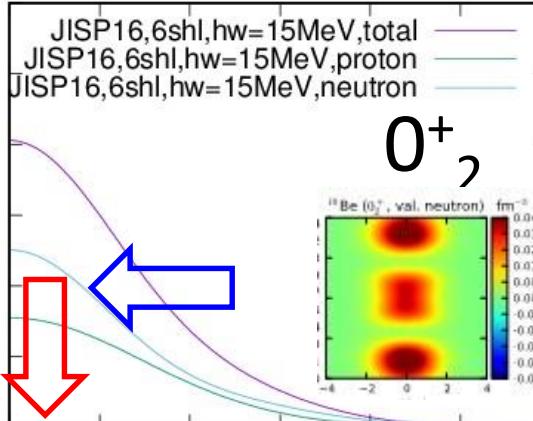
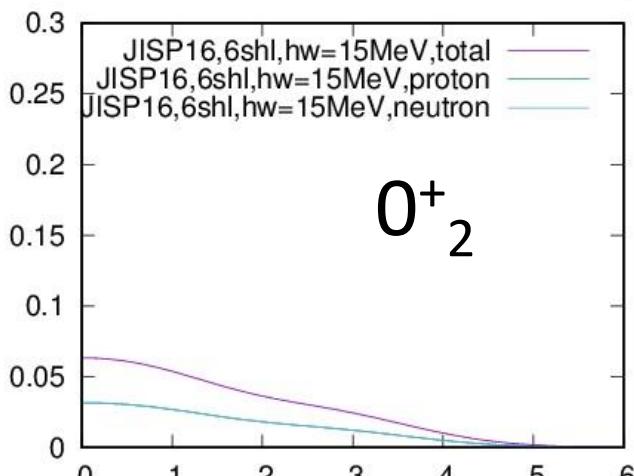
^{10}Be



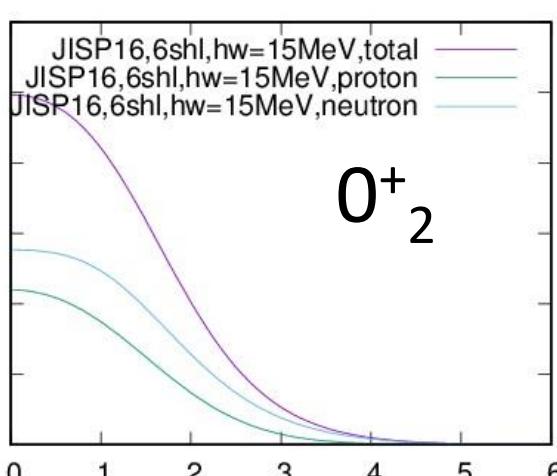
^{12}Be



0^+_2

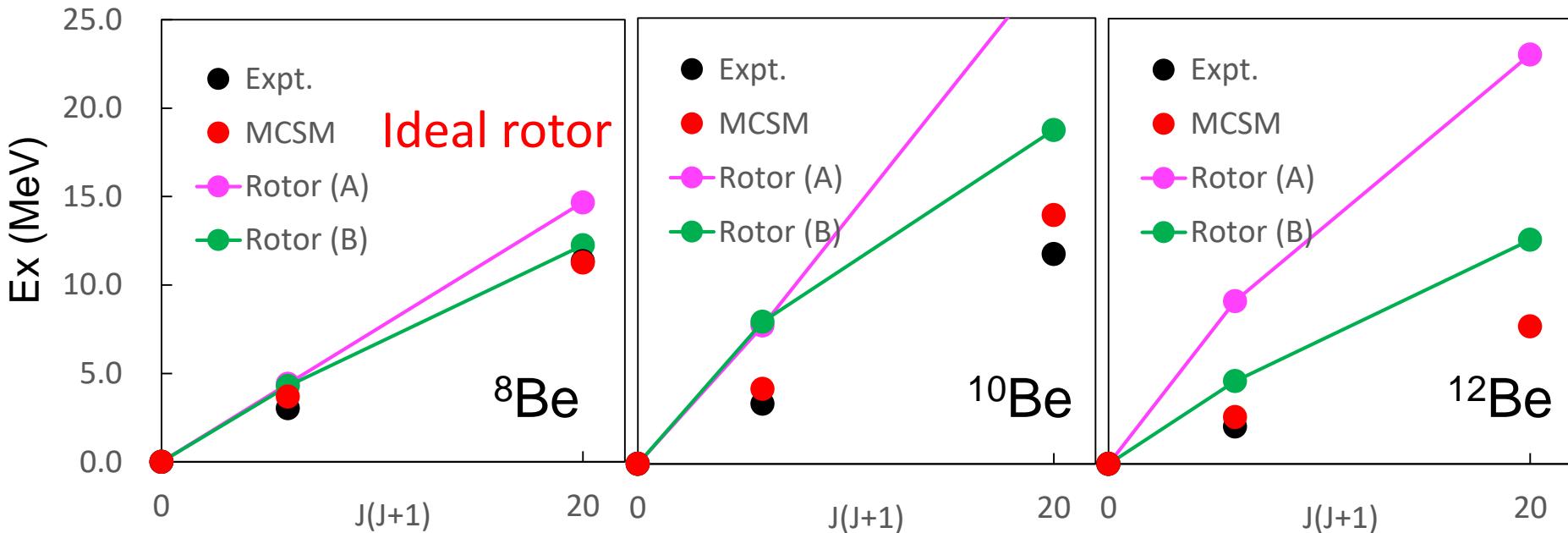


0^+_2



Radius (fm)

Rotational bands of Be isotopes



● Projection from MCSM basis vectors for $J=0$ state

● Projection from MCSM basis vectors for $J=0$ state
with amplitudes re-optimized for each J

MCSM basis vector

$$|\Psi(J, M, \pi)\rangle = \sum_i^{N_{basis}} f_i |\Phi_i(J, M, \pi)\rangle$$

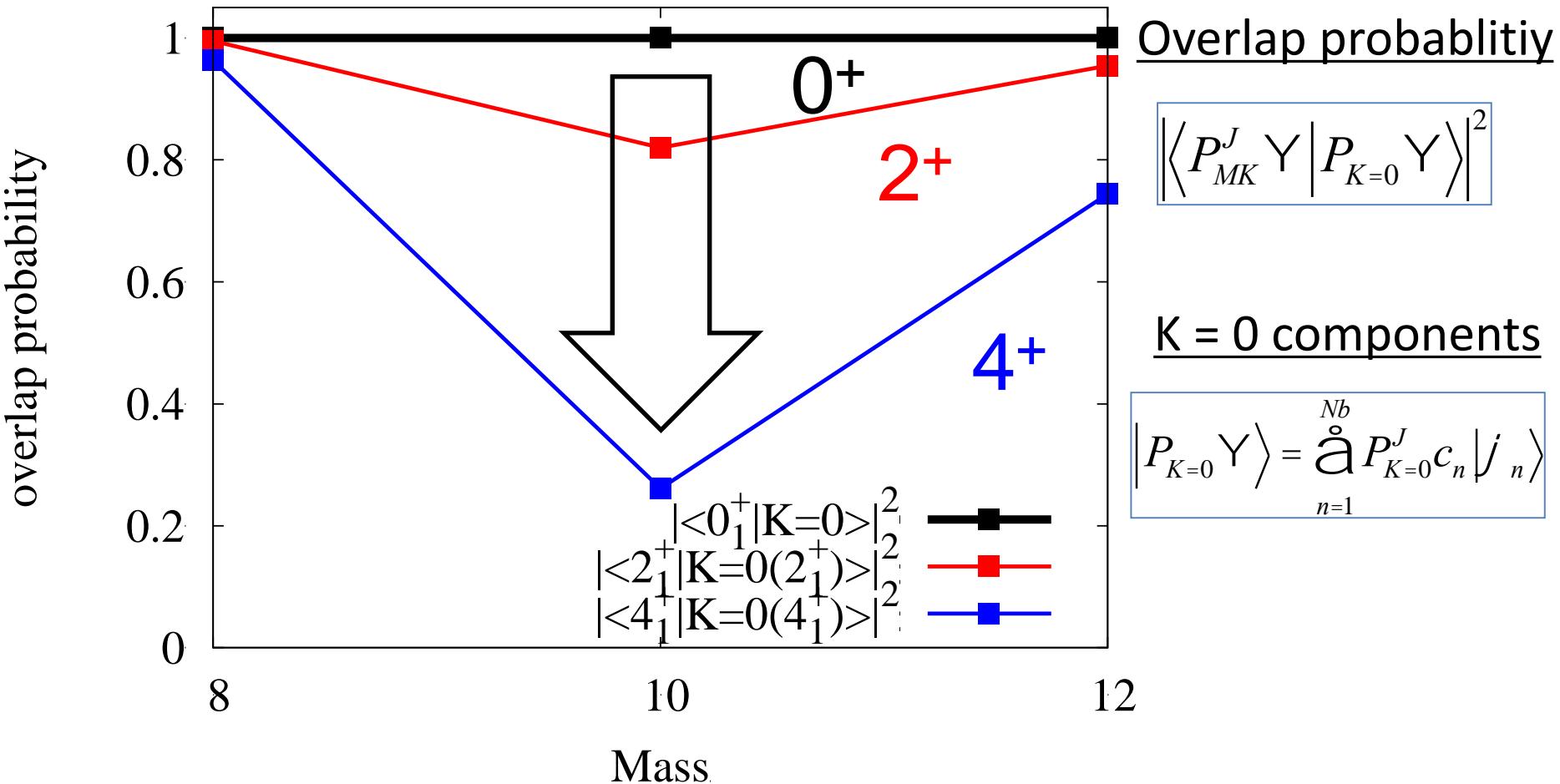
$$|\Phi(J, M, \pi)\rangle = \sum_K g_K P_{MK}^J P^\pi |\phi\rangle$$

$K = 0$ components

$$\left| P_{K=0} Y \right\rangle = \sum_{n=1}^{Nb} P_{K=0}^J c_n |j_n\rangle$$

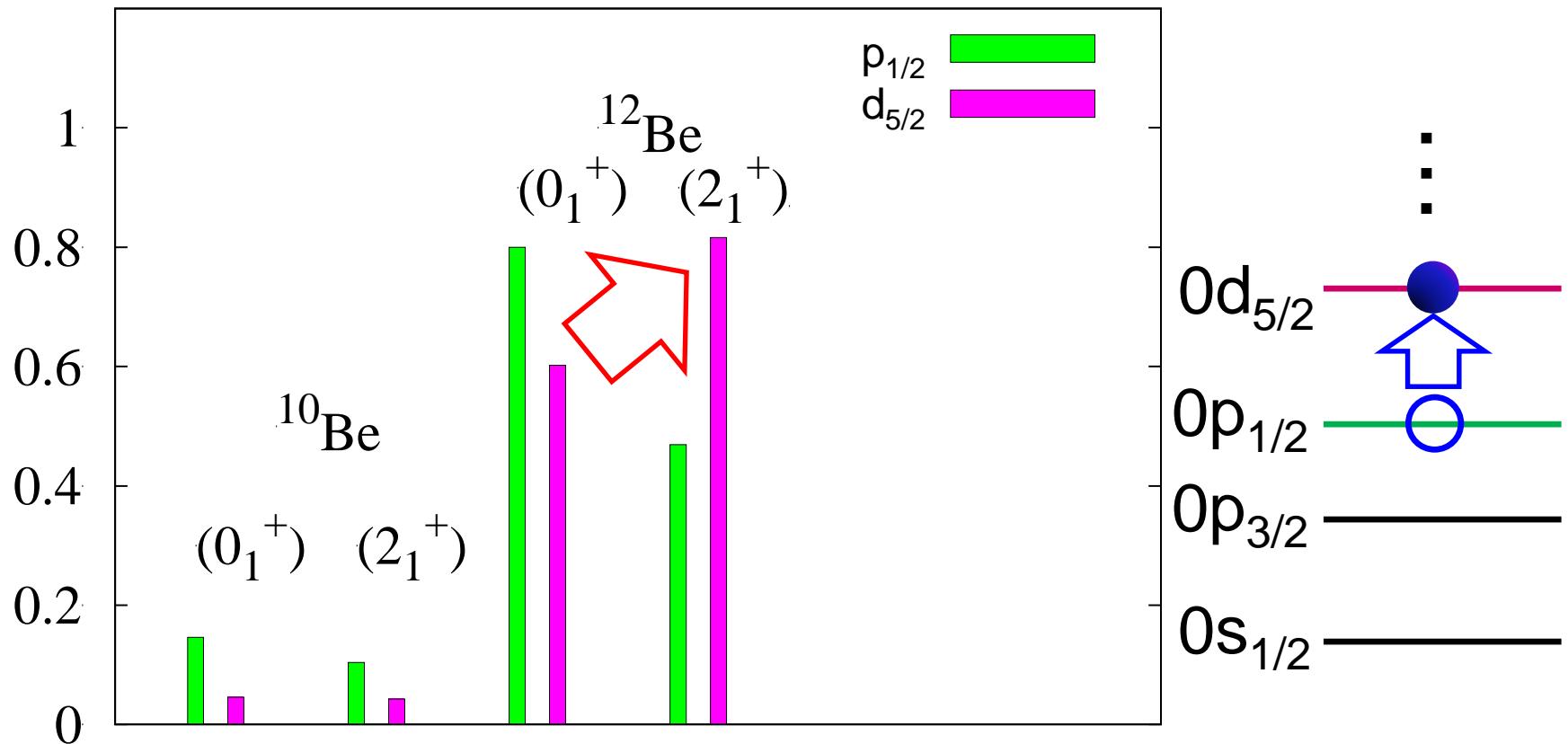
$K \neq 0$ components in basis vectors

Large K-mixing in ^{10}Be



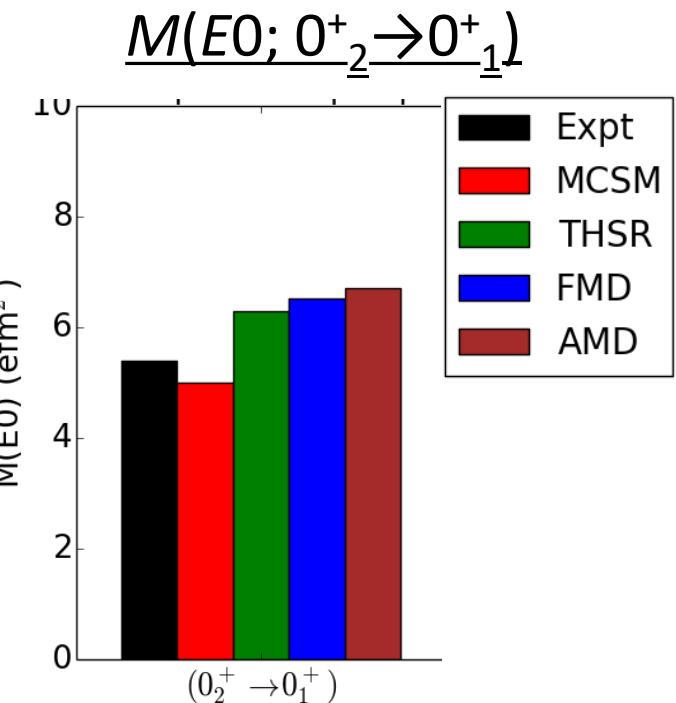
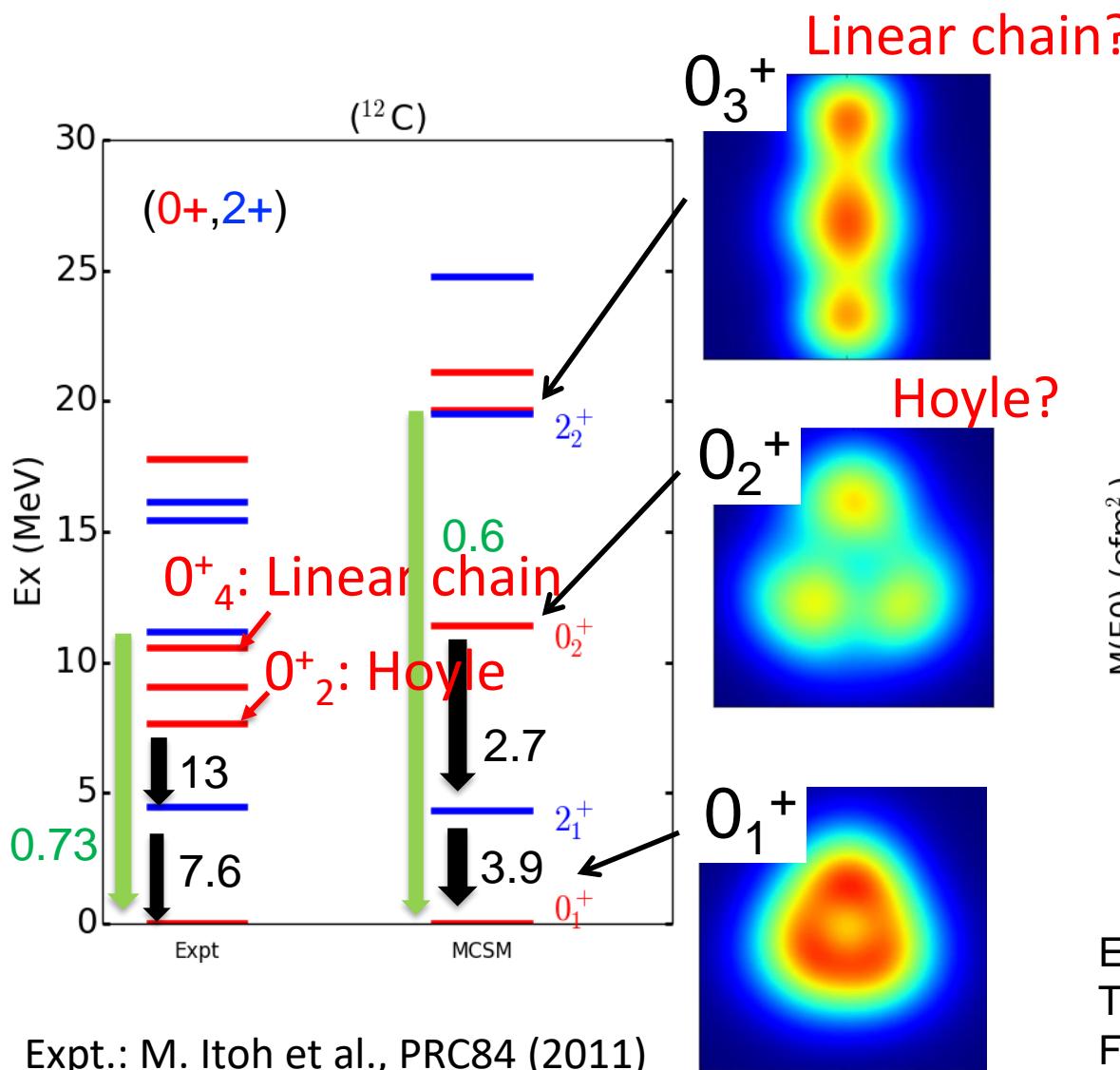
Occupation numbers of valence neutrons

valence neutron $p_{1/2}$ and $d_{5/2}$ occupation number



Excitation of valence neutrons from $0p_{1/2}$ to $0d_{5/2}$ orbital in ^{12}Be

Energy level & transition strength of ^{12}C

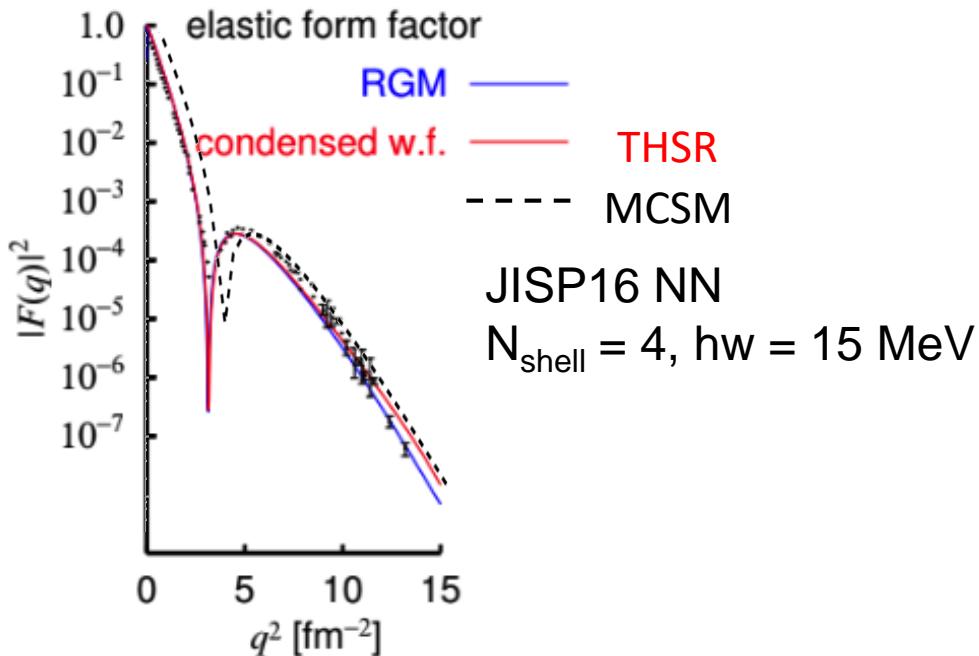


Expt.: P. Strehl 1970
 THSR: Y. Funaki 2015
 FMD: M. Chernykth 2007
 AMD: Y. Kanada-En'yo 2007

$$E_{\text{gs}} = -76.64 \text{ MeV} (\text{MCSM, JISP16, } N_{\text{shell}} = 6, h\nu = 15 \text{ MeV})$$

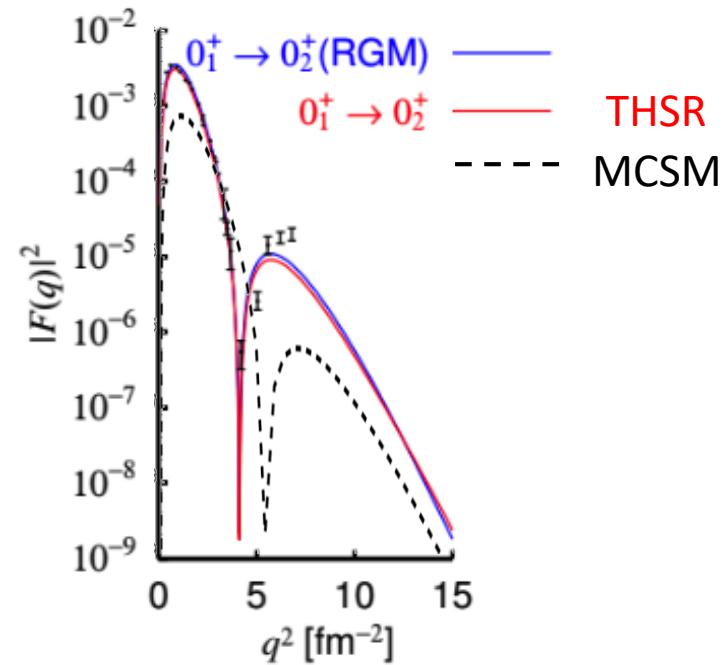
Elastic and inelastic form factors

Elastic form factor ($0^+_1 \rightarrow 0^+_1$)



(a)

Inelastic form factor ($0^+_1 \rightarrow 0^+_2$)



(b)

$$|F(q)|^2 = \frac{4\pi}{12^2} \left| \int_0^\infty \rho_{J,0_1}^{(J)}(r) j_J(qr) r^2 dr \right|^2 \exp \left(-\frac{1}{2} a_p^2 q^2 \right). \quad (1)$$

$$\rho_{J,0_1}^{(J)}(r) = \langle \Psi_{\lambda=k}^{JM} | \sum_{i=1}^{12} \delta(\mathbf{r} - \mathbf{r}_i) | \Psi_{\lambda=1}^{J=0} \rangle / Y_{JM}^*(\hat{\mathbf{r}}), \quad (2)$$

Figs, & Eqs. taken from Y. Funaki et al.,
Eur.Phys.J. A28 (2006) 259-263

Preliminary

Summary

- MCSM results for light nuclei ($A \leq 20$) w/ a NN potential can be extrapolated to the infinite basis space to obtain ab initio solution.
 - Daejoen16 NN interaction gives better agreement w. experimental data than those by JISP16.
- Cluster structure of Be & C isotopes can be investigated using MCSM wave functions and we can observe two-alpha cluster of nucleons and molecular-orbital structure of valence neutrons in ^{10}Be .

Future perspective

- Heavier nuclei beyond ^{20}Ne
- Quantitative analysis on cluster structure of Be & C isotopes

Collaborators

- Takaharu Otsuka (RIKEN, Tokyo, Leuven, MSU)
- Yutaka Utsuno (JAEA)
- Noritaka Shimizu (Tokyo)
- Tooru Yoshida (RIST)
- James P Vary (Iowa State U)
- Pieter Maris (Iowa State U)
- Petr Navratil (TRIUMF)
- Takayuki Miyagi (TRIUMF)

Supported by MEXT and JICFuS

Priority Issue 9 to be Tackled by Using Post K Computer “Elucidation of the Fundamental Laws and Evolution of the Universe”