# Eigenvector continuation in nuclear physics 

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## Motivation

## Many physics problems are tremendously difficult...

- huge matrices, possibly too large to store
- ever more so given the evolution of typical HPC clusters
- most exact methods suffer from exponential scaling
- interest only in a few (lowest) eigenvalues



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## Introducing eigenvector continuation

D. Lee, TRIUMF Ab Initio Workshop 2018; Frame et al., PRL 121032501 (2018)


- novel numerical technique
- can solve otherwise untractable problems
- amazingly simple in practice
- broadly applicable
- this talk: look for nuclear nails


## Hubbard model

- three-dimensional Bose-Hubbard model (4 bosons on $4 \times 4 \times 4$ lattice)
- hopping parameter $t$, on-site interaction $U \leadsto H=H(c=U / t)$

- Bose gas for $c>0$, weak binding for $-3.8<c<0$, tight cluster for $c<-3.8$
- eigenvector continuation can extrapolate across regimes




## General idea

## Scenario

- consider physical state (eigenvector) in a large space
- parametric dependence of Hamiltonian $H(c)$ traces only small subspace


## Procedure

- calculate $\left|\psi\left(c_{i}\right)\right\rangle, i=1, \ldots N_{\mathrm{EC}}$ in "easy" regime
- solve generalized eigenvalue problem $H|\psi\rangle=\lambda N|\psi\rangle$ with
- $H_{i j}=\left\langle\psi_{i}\right| H\left(c_{\text {target }}\right)\left|\psi_{j}\right\rangle$
- $N_{i j}=\left\langle\psi_{i} \mid \psi_{j}\right\rangle$


## Prerequisite

- smooth dependence of $H(c)$ on $c$
- enables analytic continuation of $|\psi(c)\rangle$ from $c_{\text {easy }}$ to $c_{\text {target }}$


## SRG evolution

- unitary transformation of Hamiltonian: $H \rightarrow H_{\lambda}=U_{\lambda} H U_{\lambda}^{\dagger} \rightsquigarrow V_{\lambda}$
- decouple low and high momenta at scale $\lambda$

R. Furnstahl, HUGS 2014 lecture slides
- interaction becomes more amenable to numerical methods...
- ...at the cost of induced many-body forces!


Bogner et al., PPNP 6594 (2010)


Hebeler+Furnstahl, RPP 76126301 (2013)

## SRG evolution = ODE solving

$$
\frac{\mathrm{d} H_{s}}{\mathrm{~d} s}=\frac{\mathrm{d} V_{s}}{\mathrm{~d} s}=\left[\left[G, H_{s}\right], H_{s}\right], \lambda=1 / s^{1 / 4}
$$

ordinary differential equation ensures smooth parametric dependence
$\hookrightarrow$ SRG evolution satisfies EC prerequisites!

## Reverse SRG

Consider $A=3,4$ test cases

- EMN N3LO(500) interaction, Jacobi NCSM calculation

Entem et al., PRC 96024004 (2017); A. Ekström implementation of Navratil et al., PRC 61044001 (2000)

## Reverse SRG

## Consider $\mathrm{A}=3,4$ test cases

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Not even induced 3N forces kept here!


- possible to extrapolate back from small $\lambda$ to bare interaction
- information about missing many-body forces in wavefunctions
- not in any single wavefunction, but in how they change


## Mind the gap

## Still no free lunch, however...

- EC is a variational method
- cannot go beyond what bare interaction gives in same model space!



## So now what?

## Perturbation theory

- consider a Hamiltonian diagonalized in a (small) subspace

$$
\begin{gathered}
H=\left(\begin{array}{ll}
H_{\phi \phi} & H_{\phi \psi} \\
H_{\psi \phi} & H_{\psi \psi}
\end{array}\right) \\
N_{0}=\operatorname{dim} H_{\phi \phi} \ll \operatorname{dim} H=N_{1} \\
H_{\phi \phi}=\operatorname{diag}\left(\left\{\lambda_{i}\right\}_{i=1, \cdot \cdot N_{0}}\right)
\end{gathered}
$$



- factor out large number $X$ from diagonal entries of $H_{\psi \psi}$
- perturbative expansion for lowest eigenvalue and vector

$$
\left|\psi_{1}\right\rangle=\sum_{n=0}^{\infty} X^{-n}\left(\sum_{i=1}^{N_{0}} x_{i}^{(n)}\left|\phi_{i}\right\rangle+\sum_{j=N_{0}+1}^{N_{1}} x_{j}^{(n)}\left|\psi_{j}\right\rangle\right), \lambda_{1}^{\text {full }}=\sum_{n=0}^{\infty} X^{-n} \lambda_{1}^{(n)}
$$

- matching powers gives coupled recursive expressions for $x_{j}^{(n)}$ and $\lambda_{1}^{(n)}$


## Perturbation theory (continued)

Diagonalizing a small space can still be too expensive...

## Perturbation theory (continued)

Diagonalizing a small space can still be too expensive...

- actually, a partial diagonalization per se is ok ( $\rightarrow$ Lanczos)
- but transforming the Hamiltonian is problematic...

- cost for adjusting off-diagonal elements is prohibitive
- scales with size of the full (large) space


## Way out

Start from one-dimensional space $\left(N_{\max }=0\right) \ldots$

...i.e., directly use the given Hamiltonian

## Failure

## ${ }^{3} \mathrm{H}$ NCSM calculation, $N_{\max }=12$ model space

- EMN N3LO 500 interaction

- perturbation theory does not converge!
- however, interaction clearly "more perturbative" for small SRG $\lambda$
- convergence perhaps for very small $\lambda$


## Saved by EC

- span space by the wavefunction corrections $\left|\psi_{1}^{(n)}\right\rangle \rightarrow x_{j}^{(n)}, n=0, \cdots$ order
- evaluate Hamiltonian between these states
- interpretation: $H=H_{\text {diag }}+c H_{\text {off-diag }}$, EC-extrapolate to $c=1$

- same input as PT, but now things converge (to the correct result!)


## Note

expensive part of is setting up the $x_{j}^{(n)}, j=1, \cdots N_{1}$
essentially an $N_{1}$-dim. matrix-vector multiplication...
$\hookrightarrow$ compare PT-EC to Lanczos!

## EC vs. Lanczos

- for EC: effective $N_{m v}=2 \times($ order -1$)$
- comparison: vanilla Lanczos in GNU Octave (i.e., ARPACK)

- EC looks quite competitive in this benchmark!
- but note: only calculating a single eigenvalue here


## Summary and outlook

## This talk

- eigenvector continuation can be used to reverse SRG
- conceptually interesting: implicit information about induced forces
- convergent perturbative model-space extension
- effectively tame divergent expansion coefficients
- interesting as computational method


## Future directions

- larger systems, other methods
- in particular: m-scheme NCSM
- combined model-space and SRG EC
- other applications



## Thanks...

...to my collaborators:

- A. Schwenk, K. Hebeler (TU Darmstadt)
- D. Lee, A. Sarkar (Michigan State U.)
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...and to you, for your attention!

