

TRIUMF Theory Workshop (Mar. 3 – 6, 2020)
Progress in Ab Initio Techniques in Nuclear Physics

Intrinsic structure of light nuclei from no-core Monte Carlo shell model

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Supported by MEXT and JICFuS

Priority Issue 9 to be Tackled by Using Post K Computer “Elucidation of the
Fundamental Laws and Evolution of the Universe”

TRIUMF
March 3, 2020

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- Yusuke Tsunoda (Tokyo)
- James P Vary (Iowa State U)
- Pieter Maris (Iowa State U)

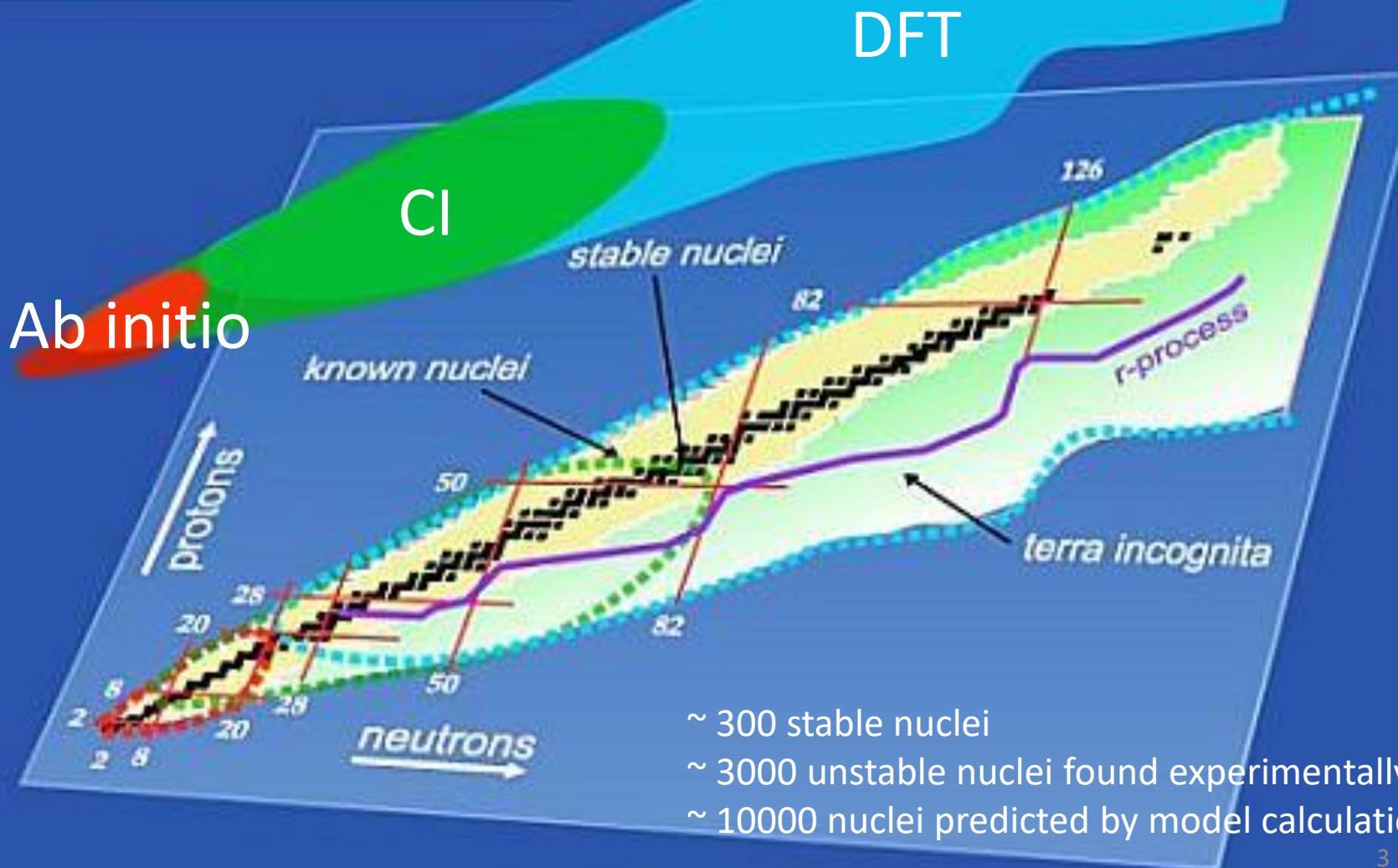
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Priority Issue 9 to be Tackled by Using Post K Computer “Elucidation of the Fundamental Laws and Evolution of the Universe”

Nuclear Landscape

UNEDF SciDAC Collaboration: <http://unedf.org/>

-  Ab initio
-  Configuration Interaction
-  Density Functional Theory



Ab initio
Configuration Interaction
Density Functional Theory

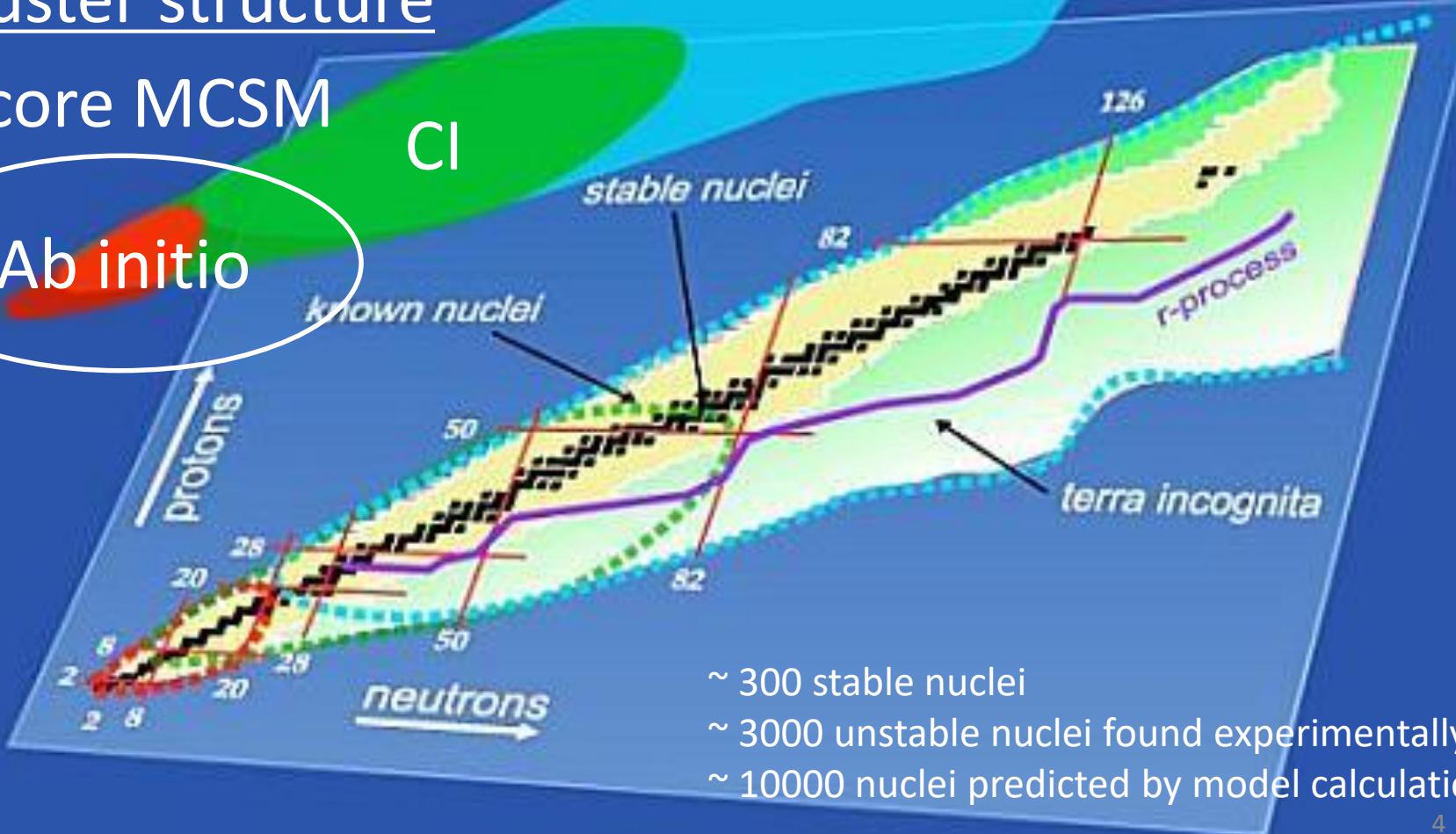
α -cluster structure

No-core MCSM

Ab initio

CI

DFT



“Ab initio” in low-energy nuclear structure physics

- Major challenge in nuclear physics
 - Nuclear structure & reactions directly from *ab-initio* calc. w/ nuclear forces
 - *ab-initio* approaches in nuclear structure calculations ($A > 4$):
Light mass: Green's Function Monte Carlo, No-Core Shell Model ($A \sim 12$),
Medium/heavy mass: Coupled Cluster, IM-SRG,
Self-consistent Green's Function theory, Lattice EFT, UMOA, ...
- Solve the non-relativistic many-body Schroedinger eq.
and obtain the eigenvalues and eigenvectors.

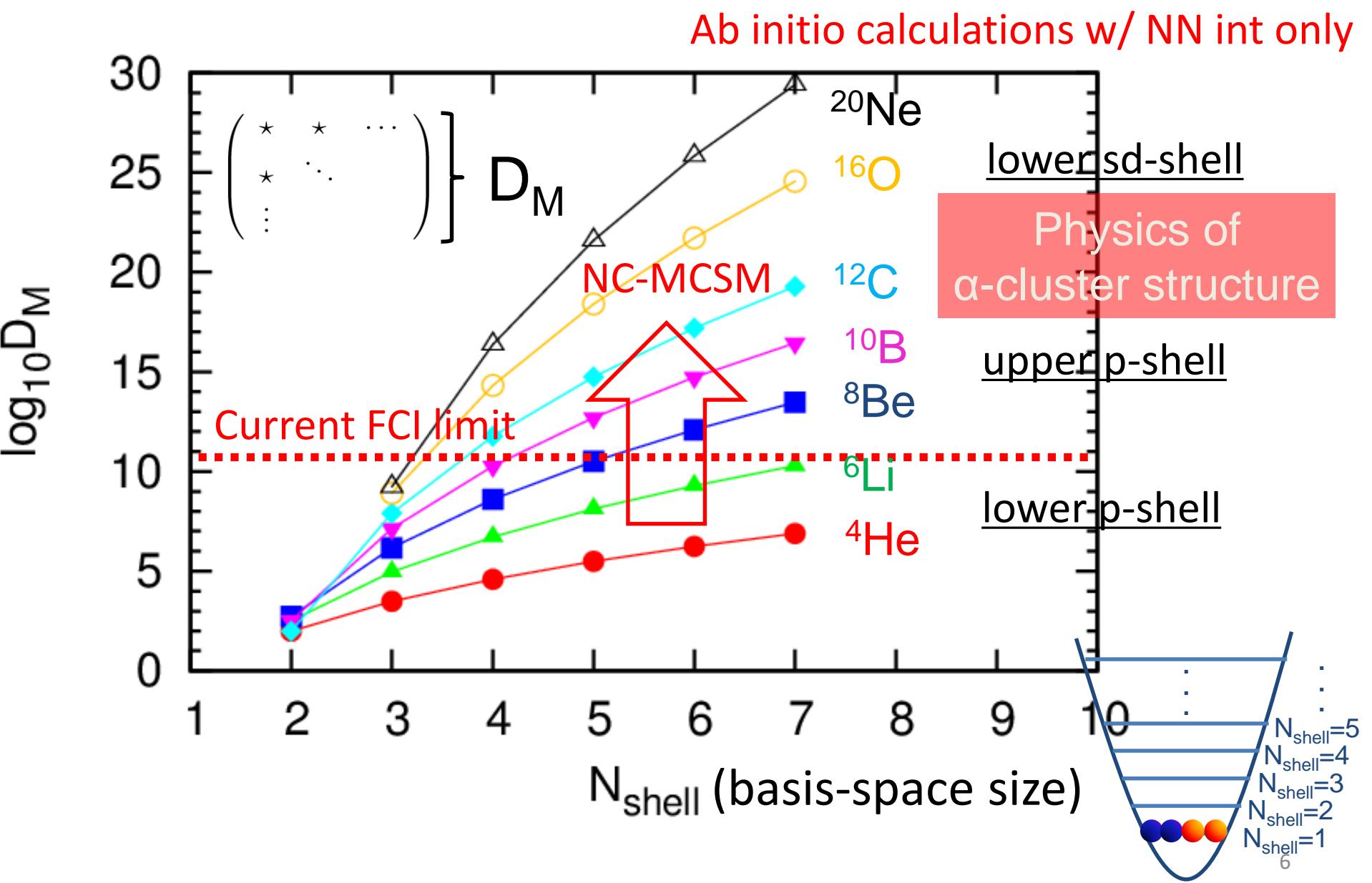
$$H|\Psi\rangle = E|\Psi\rangle$$

$$H = T + V_{\text{NN}} + V_{\text{3N}} + \dots + V_{\text{Coulomb}}$$

- **Ab initio**: All nucleons are active, and Hamiltonian consists of realistic NN (+ 3N + ...) potentials.

→ Computationally demanding → Monte Carlo shell model (MCSM)

M-scheme dimension



Monte Carlo shell model (MCSM)

Standard shell model

$$H = \begin{pmatrix} * & * & * & * & * & \cdots \\ * & * & * & * & & \\ * & * & * & & & \\ * & * & & \ddots & & \\ * & & & & & \\ \vdots & & & & & \end{pmatrix} \xrightarrow{\text{Diagonalization}} \begin{pmatrix} E_0 & & & & & 0 \\ & E_1 & & & & \\ & & E_2 & & & \\ & & & \ddots & & \\ 0 & & & & & \end{pmatrix}$$

Large sparse matrix
 $\sim \mathcal{O}(10^{10})$ # non-zero MEs
 $\sim \mathcal{O}(10^{13-14})$

- Importance truncation

Monte Carlo shell model

$$H \sim \begin{pmatrix} * & * & \cdots \\ * & \ddots & \\ \vdots & & \end{pmatrix} \xrightarrow{\text{Diagonalization}} \begin{pmatrix} E'_0 & & 0 \\ & E'_1 & \\ 0 & & \ddots \end{pmatrix}$$

Important bases stochastically selected $\sim \mathcal{O}(100)$

$$|\Psi(J, M, \pi)\rangle = \sum_i^{N_{basis}} f_i |\Phi_i(J, M, \pi)\rangle$$

$$|\Phi(J, M, \pi)\rangle = \sum_K g_K P_{MK}^J P^\pi |\phi\rangle$$

diagonalization

Deformed

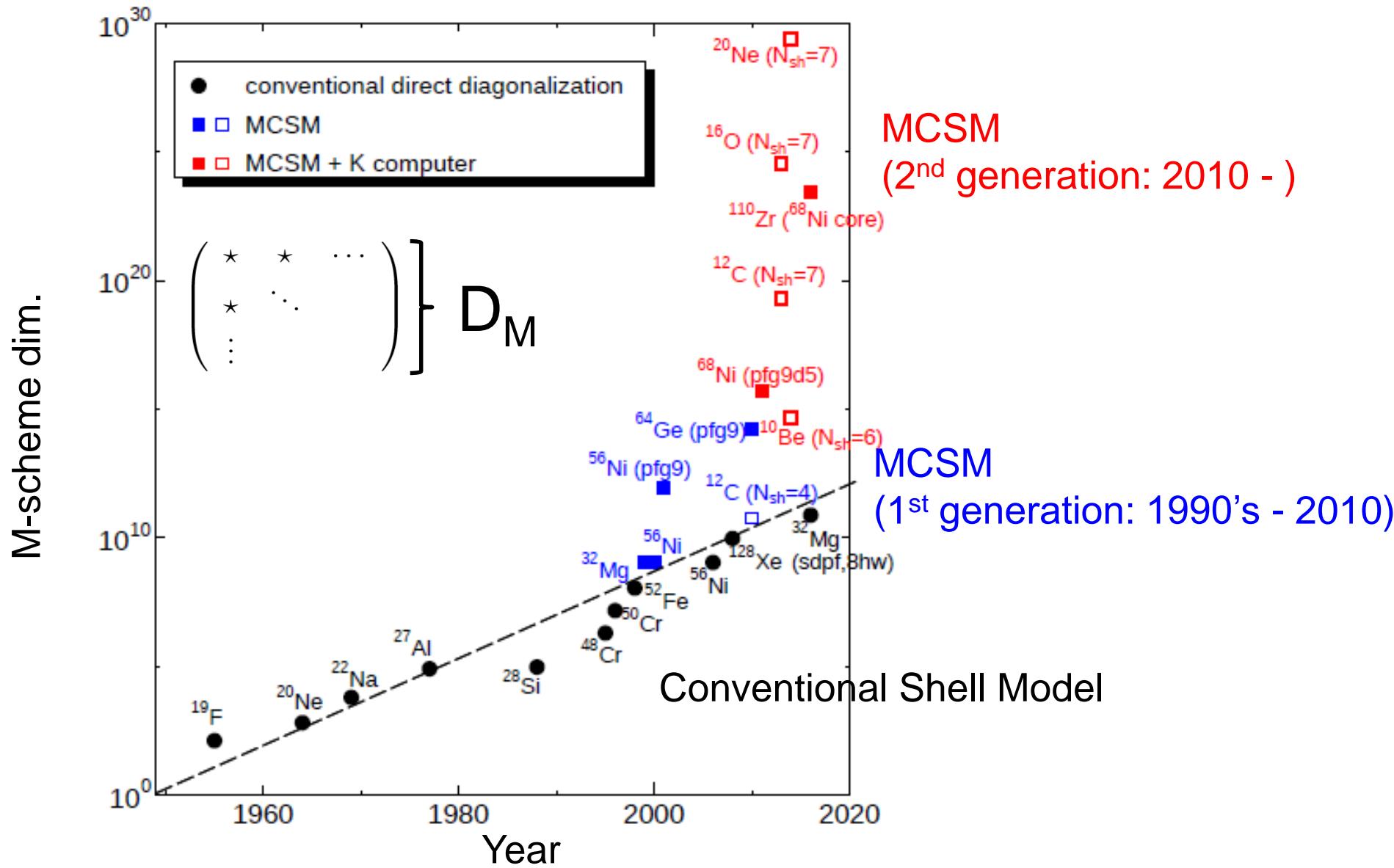
$$|\phi\rangle = \prod_i^A a_i^\dagger |-\rangle$$

Spherical

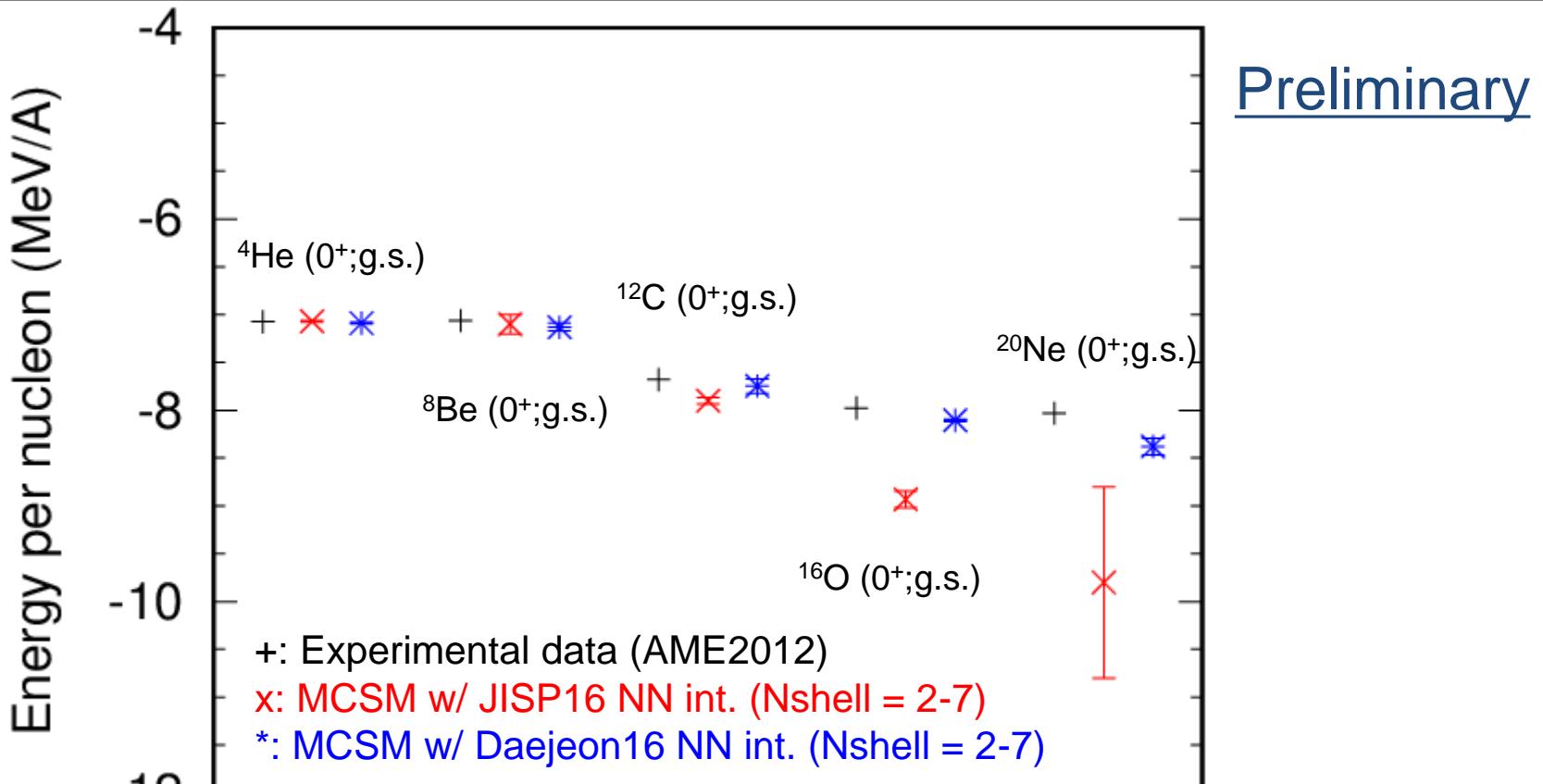
$$a_i^\dagger = \sum_\alpha c_\alpha^\dagger D_{\alpha i}$$

stochastic sampling & CG method

Historical evolution/development of the MCSM



Comparison of MCSM results w/ experiments



MCSM results are obtained using K computer by traditional extrapolation w/ optimum harmonic oscillator energies.

JISP16 results show good agreements w/ experimental data up to ${}^{12}\text{C}$, slightly overbound for ${}^{16}\text{O}$, and clearly overbound for ${}^{20}\text{Ne}$.

Daejeon16 results show good agreements w/ experimental data up to ${}^{20}\text{Ne}$.

Density distribution from ab initio calc.

- Green's function Monte Carlo (GFMC)
 - "Intrinsic" density is constructed by aligning the moment of inertia among samples
R. B. Wiringa, S. C. Pieper, J. Carlson, & V. R. Pandharipande, Phys. Rev. C62, 014001 (2000)
- No-core full configuration (NCFC)
 - Translationally-invariant density is obtained by deconvoluting the intrinsic & CM w.f.
C. Cockrell J. P. Vary & P. Maris, Phys. Rev. C86, 034325 (2012)
- Lattice EFT
 - Triangle structure of carbon-12
E. Epelbaum, H. Krebs, T. A. Lahde, D. Lee, & U.-G. Meissner, Phys. Rev. Lett. 109, 252501 (2012), ...
- FMD
 - H. Feldmeier, Nucl. Phys. A515, 147 (1990), ...

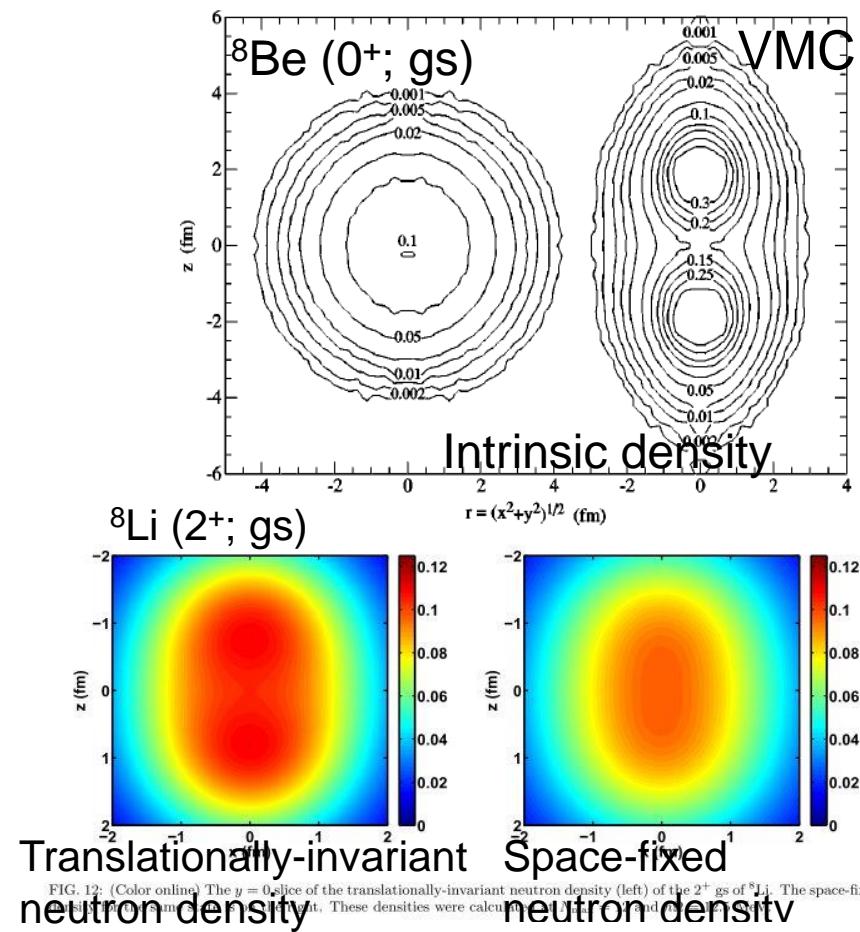
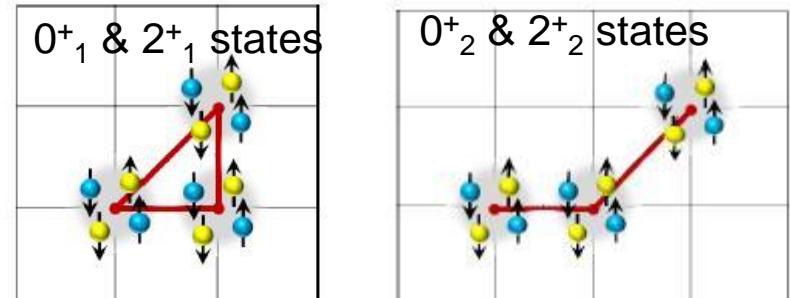


FIG. 12: (Color online) The $y = 0$ slice of the translationally-invariant neutron density (left) of the 2^+ gs of ^{8}Li . The space-fixed neutron density is shown on the right. These densities were calculated with the same input parameters and basis sets as the intrinsic density.



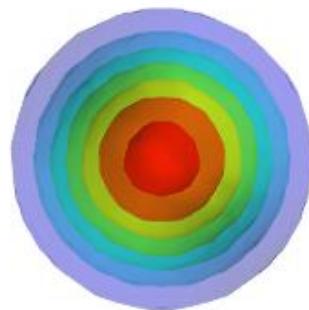
Density distribution in MCSM

$$|\Phi\rangle = \sum_{i=1}^{N_{basis}} c_i |\Phi_i\rangle = c_1 \begin{array}{c} \text{image} \\ \text{of a density} \\ \text{distribution} \end{array} + c_2 \begin{array}{c} \text{image} \\ \text{of a density} \\ \text{distribution} \end{array} + c_3 \begin{array}{c} \text{image} \\ \text{of a density} \\ \text{distribution} \end{array} + c_4 \begin{array}{c} \text{image} \\ \text{of a density} \\ \text{distribution} \end{array} + \dots$$

Angular-momentum projection

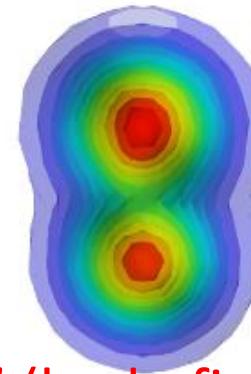
$$|\Psi\rangle = \sum_{i=1}^{N_{basis}} c_i P^J P^\pi |\Phi_i\rangle$$

A way to construct
an “intrinsic” density



${}^8\text{Be}$ 0⁺ ground state

Laboratory frame



“Intrinsic” (body-fixed) frame

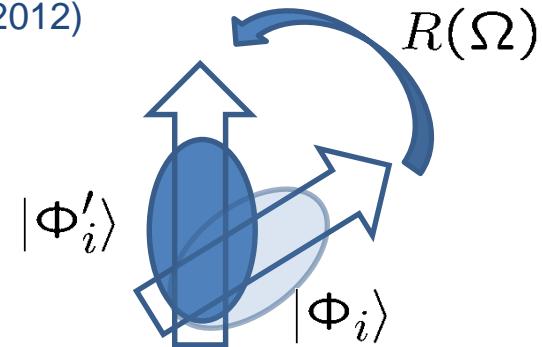
Densities in lab. & body-fixed frames can be constructed by MCSM

How to construct an “intrinsic” density from MCSM w.f.

N. Shimizu, T. Abe, Y. Tsunoda, Y. Utsuno, **T. Yoshida**, T. Mizusaki, M. Honma, T. Otsuka,
Progress in Theoretical and Experimental Physics, 01A205 (2012)

- MCSM wave function

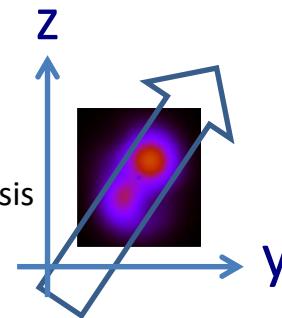
$$|\Psi\rangle = \sum_{i=1}^{N_{basis}} c_i P^J P^\pi |\Phi_i\rangle$$



- Wave function w/o the projections

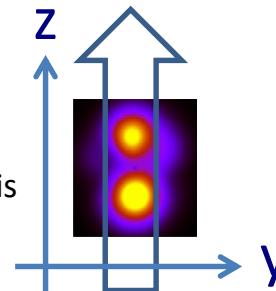
$$\sum_{i=1}^{N_{basis}} c_i |\Phi_i\rangle = c_1 \begin{matrix} \text{[Heatmap]} \\ \text{[Arrows]} \end{matrix} + c_2 \begin{matrix} \text{[Heatmap]} \\ \text{[Arrows]} \end{matrix} + \dots + c_{N_{basis}} \begin{matrix} \text{[Heatmap]} \\ \text{[Arrows]} \end{matrix}$$

Rotation by diagonalizing Q-moment
 $(Q_{zz} > Q_{yy} > Q_{xx})$



- Wave function w/o the projection w/ the alignment of Q-moment

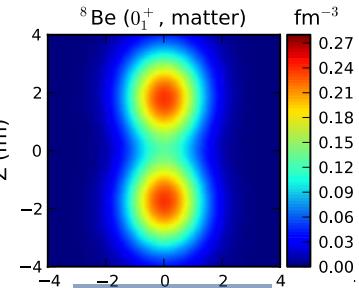
$$\sum_{i=1}^{N_{basis}} c_i |\Phi'_i\rangle = c_1 \begin{matrix} \text{[Heatmap]} \\ \text{[Arrows]} \end{matrix} + c_2 \begin{matrix} \text{[Heatmap]} \\ \text{[Arrows]} \end{matrix} + \dots + c_{N_{basis}} \begin{matrix} \text{[Heatmap]} \\ \text{[Arrows]} \end{matrix}$$



Density distribution of Be isotopes

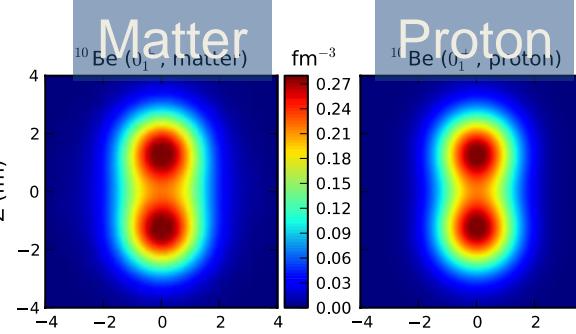
JISP16 NN, $N_{\text{shell}} = 6$, $h\nu = 15 \text{ MeV}$

^8Be
($0^+; \text{gs}$)

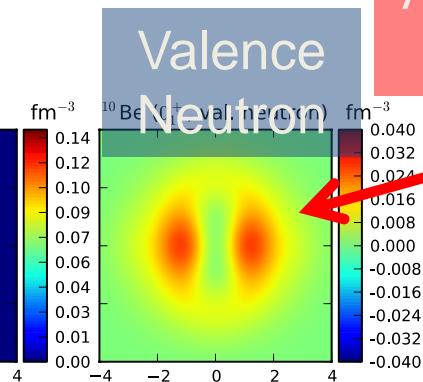


Emergence of
2- α -cluster structure

^{10}Be
($0^+; \text{gs}$)

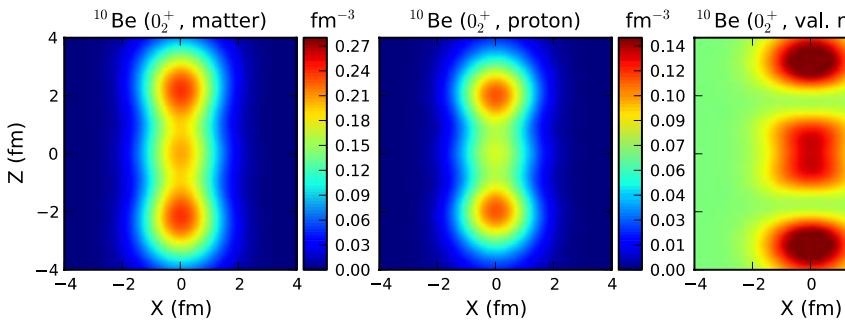


Valence
Neutron

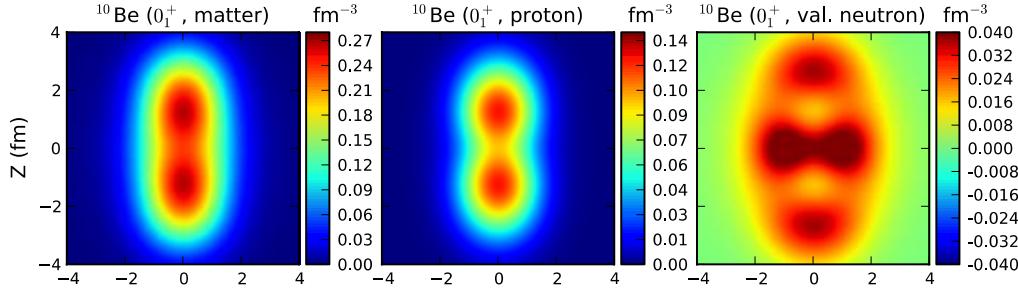


Appearance of molecular-
orbital structure

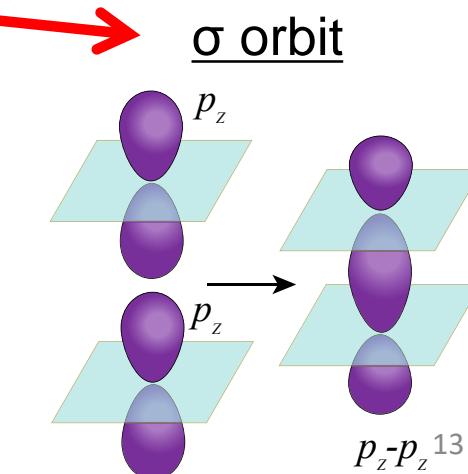
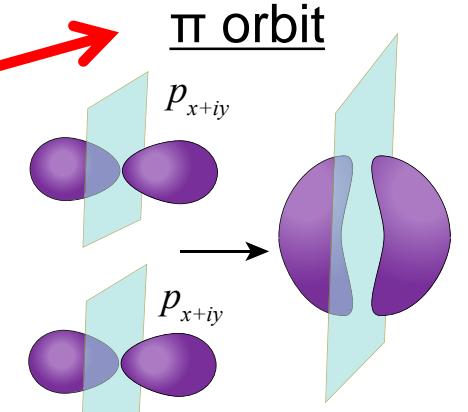
${}^0\text{e}^+$



${}^{12}\text{Be}$
($0^+; \text{gs}$)



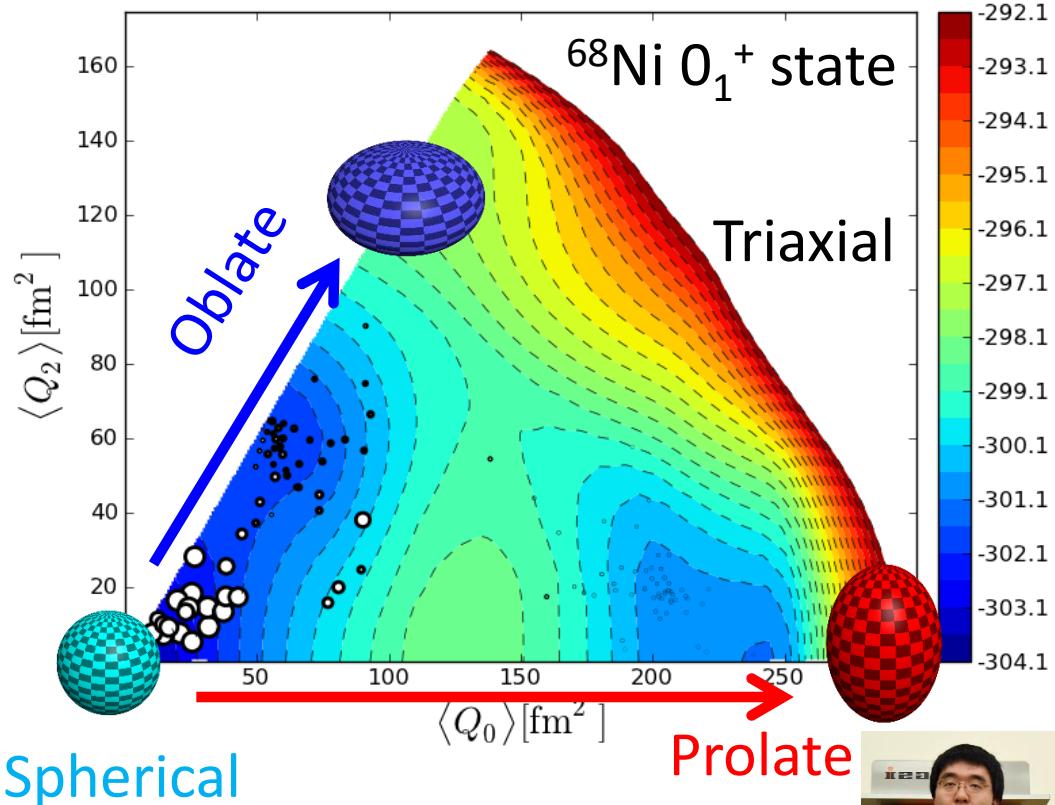
Fading 2- α structure as N increases



Preliminary

T-plot in MCSM

- MCSM wf: $|\Psi\rangle = \sum_{i=1}^{N_{basis}} c_i P^J P^\pi |\Phi_i\rangle$
 - Deformed SD
-> info of intrinsic shape
 - Coefficient
- Potential Energy Surface (PES):
Calculated by Q-constrained HF
- Location of circles:
Quadrupole deformation of unprojected deformed SDs
- Area of circles:
Overlap probability
btw deformed SD & MCSM wf

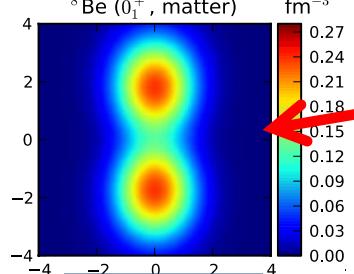


Y. Tsunoda, et al., Phys Rev C89, 031301 (R) (2014).

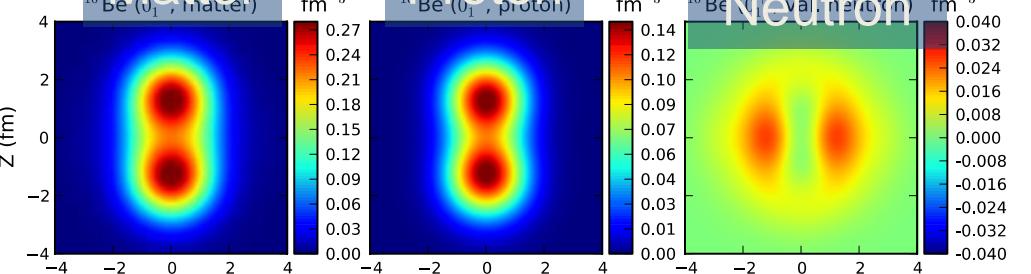


^8Be
($0^+;\text{gs}$)

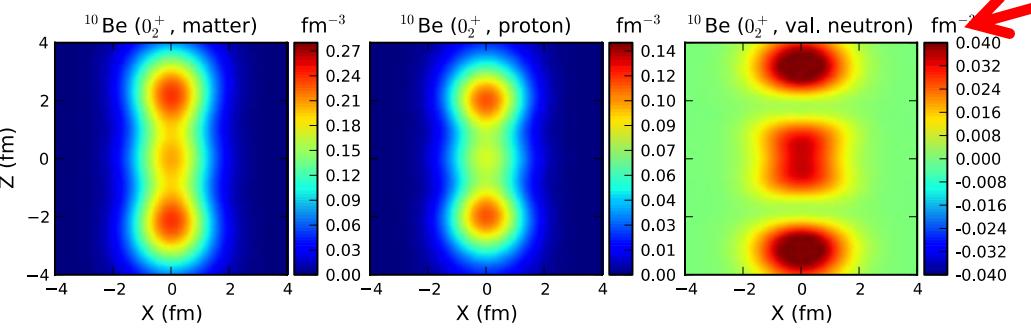
Densi



^{10}Be
($0^+;\text{gs}$)

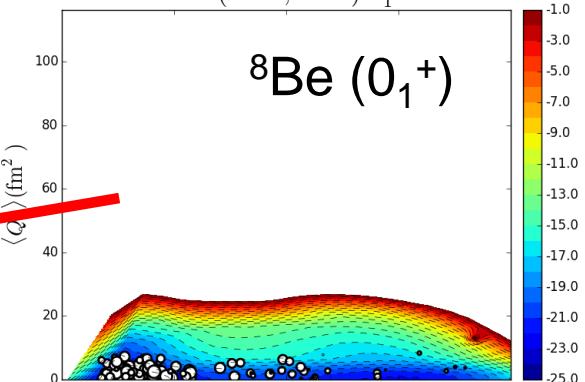


0_2^+



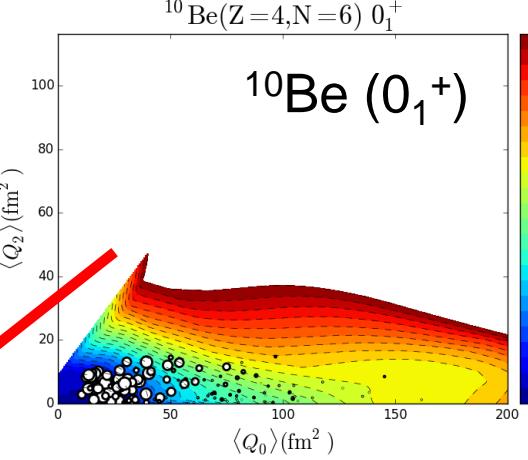
^{12}Be
($0^+;\text{gs}$)

$^8\text{Be}(Z=4,N=4) \, 0_1^+$

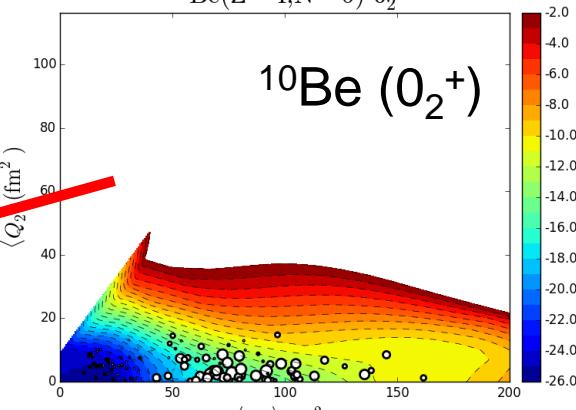


e is

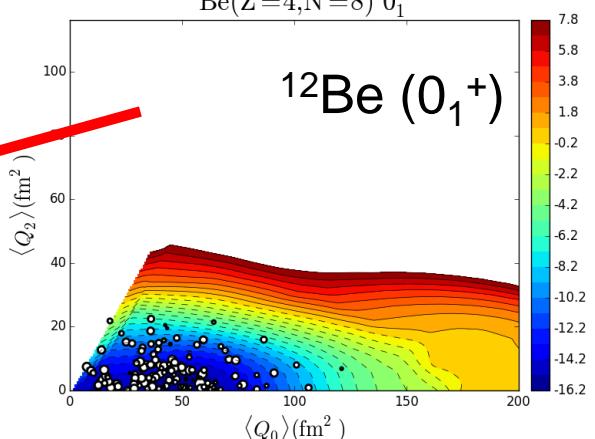
$^{10}\text{Be}(Z=4,N=6) \, 0_1^+$



$^{10}\text{Be}(Z=4,N=6) \, 0_2^+$

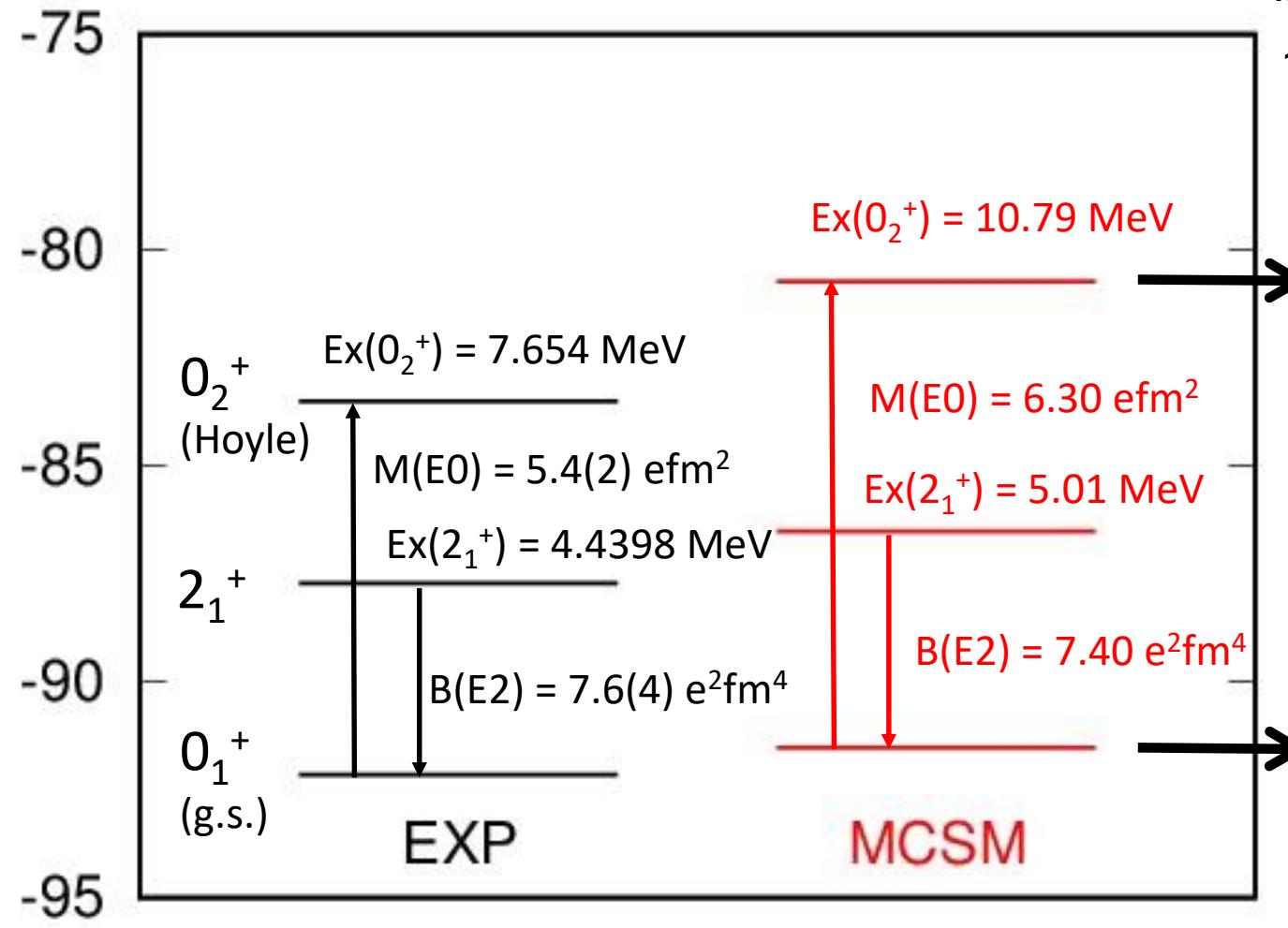


$^{12}\text{Be}(Z=4,N=8) \, 0_1^+$



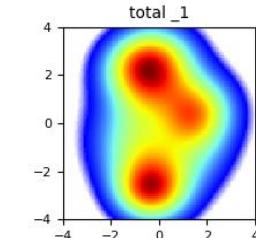
Excitation spectrum of ^{12}C

Preliminary



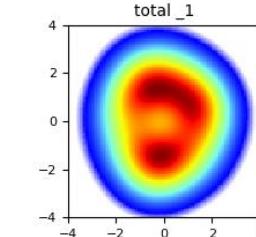
Intrinsic density

1st excited 0⁺ state
(Hoyle state?)



$$r_{pp} = 2.60 \text{ fm}$$

Ground state



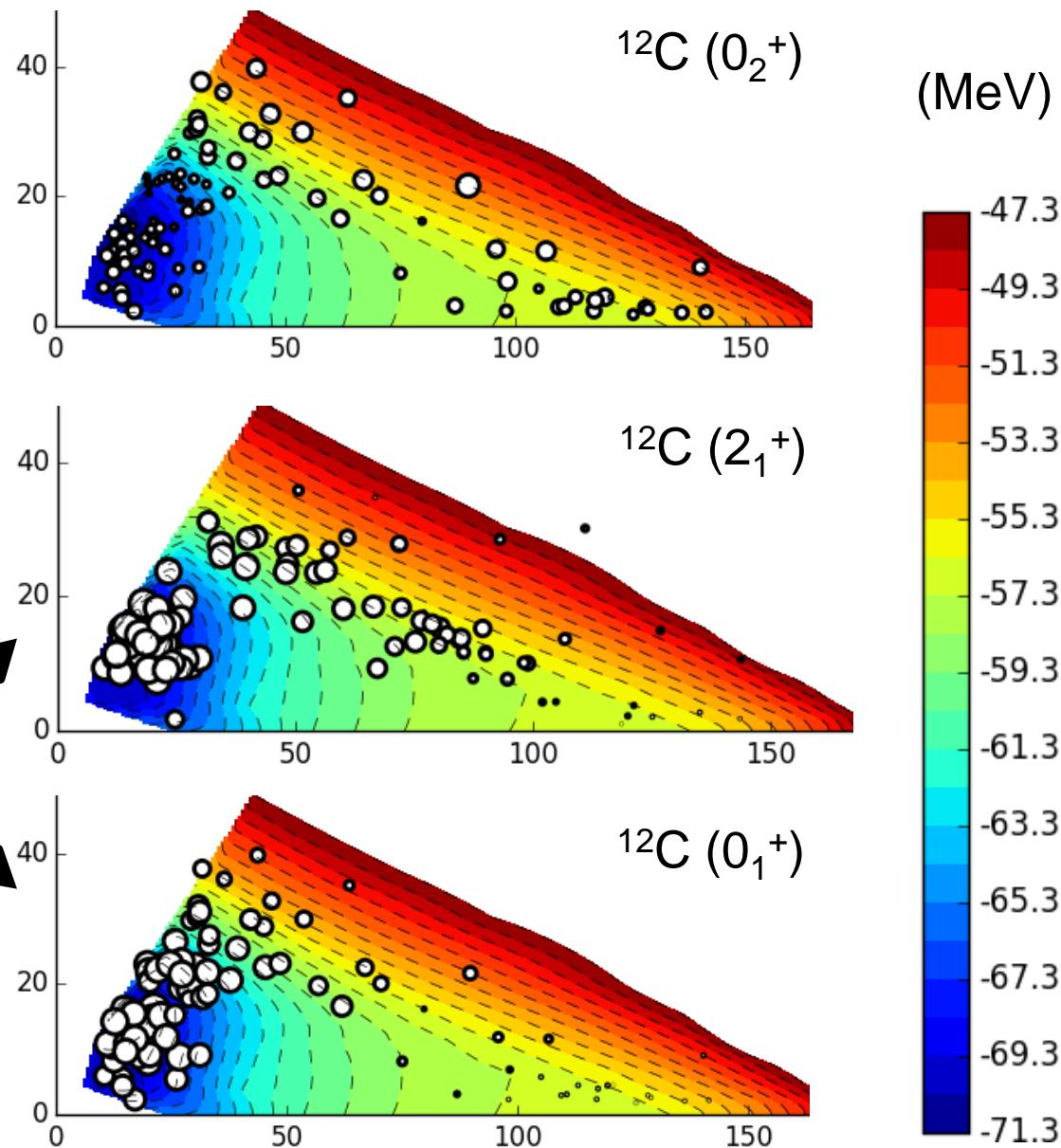
$$r_{pp} = 2.28 \text{ fm}$$

$$r_{pp} = 2.32(2) \text{ fm (exp)}$$

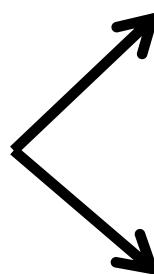
Daejeon16 NN, $N_{\text{shell}} = 7$, $hw = 20 \text{ MeV}$
w/o energy-variance extrapolation

T-plot of ^{12}C low-lying states

Scattered more than
those of ground & 2^+ states



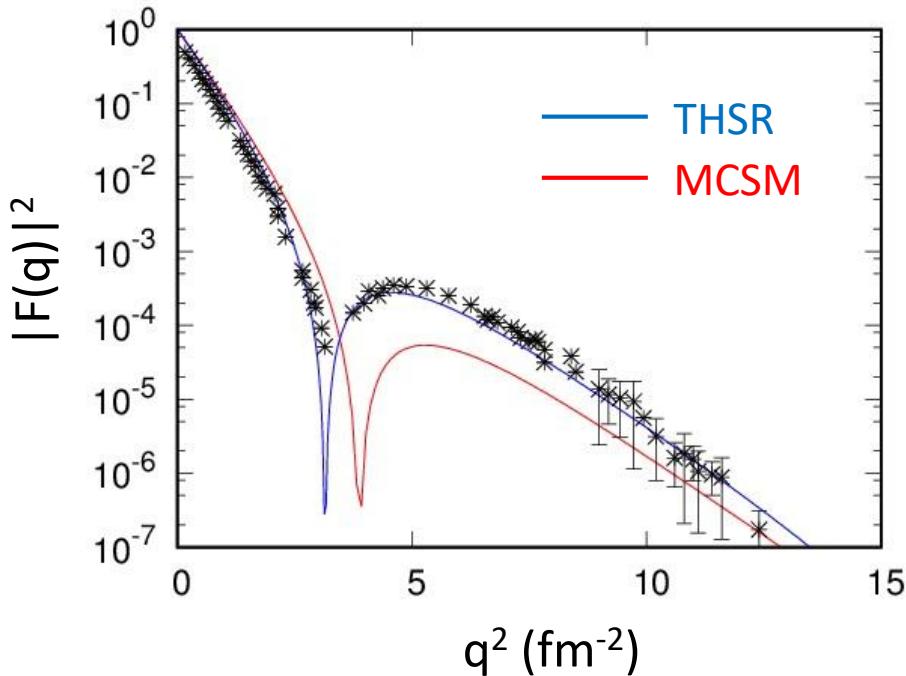
Similar distributions



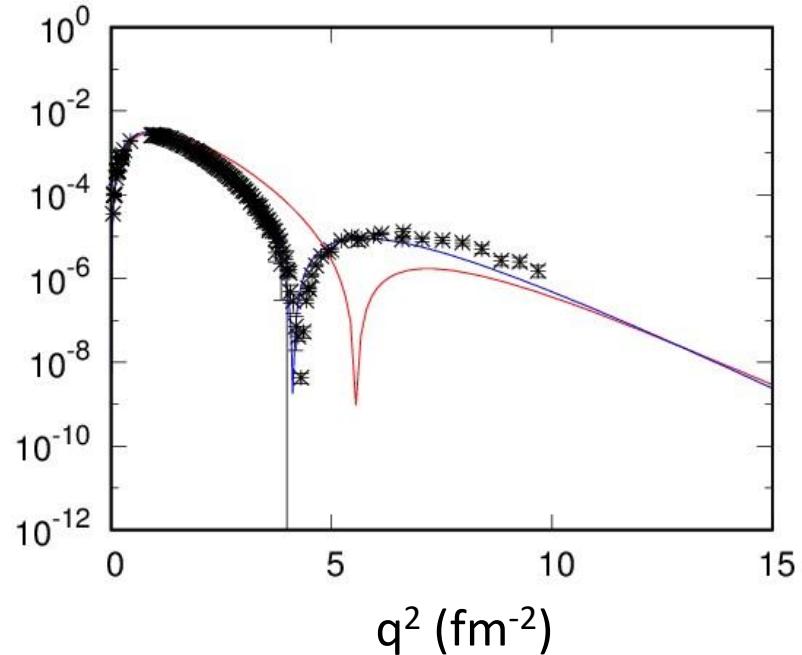
Elastic/Inelastic form factors of ^{12}C

Preliminary

Elastic form factor ($0^+_1 \rightarrow 0^+_1$)



Inelastic form factor ($0^+_1 \rightarrow 0^+_2$)



Daejeon16 NN, $N_{\text{shell}} = 7$, $h\omega = 20$ MeV

$$|F(\mathbf{q})|^2 = \frac{4\pi}{12^2} \left| \int_0^\infty \rho_{J,0_1}^{(J)}(r) j_J(qr) r^2 dr \right|^2 \exp\left(-\frac{1}{2} a_p^2 q^2\right)$$

$$\rho_{J,0_1}^{(J)}(r) = \langle \Psi_{\lambda=k}^{JM} | \sum_{i=1}^{12} \delta(\mathbf{r} - \mathbf{r}_i) | \Psi_{\lambda=1}^{J=0} \rangle / \mathbf{Y}_{JM}^*(\hat{\mathbf{r}})$$

THSR & experimental data:
Provided by Y. Funaki et al.,
Eur. Phys. J. A28 (2006) 259-263.

Summary

- MCSM results for light nuclei ($A \leq 20$) w/ a NN potential can be extrapolated to the infinite basis space to obtain ab initio solution.
 - Daejoen16 NN interaction gives better agreement w. experimental data than those by JISP16.
- Intrinsic structure of Be & C isotopes can be investigated using MCSM wave functions. In the Be isotopes, we can observe two-alpha cluster structure of nucleons and molecular-orbital structure of valence neutrons. T-plot analysis of low-lying states of Be & C isotopes was also shown.

Future perspective

- Heavier (sd-shell) nuclei beyond ^{20}Ne enable to provide a comparison with recent experiments
- Quantitative analysis on alpha-cluster structure of Be & C isotopes